

# **The Phonon Background from Gamma-Rays in Sub-GeV Dark Matter Detectors**

Kim V. Berghaus

Yang Institute for Theoretical Physics

Stony Brook University

Based on: arXiv2112.09702

Collaborators: Rouven Essig, Yonit Hochberg, Yutaro Shoji, Mukul Sholapurkar

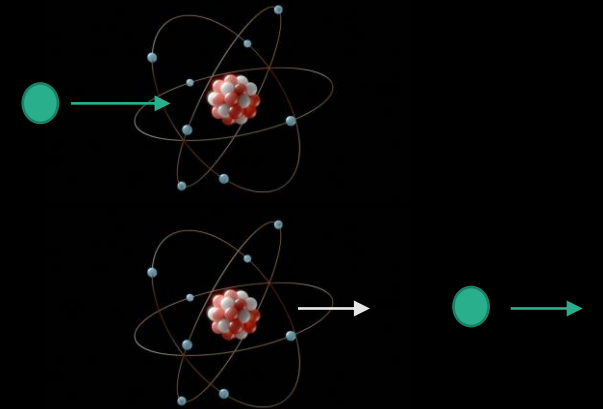
# Sub-GeV Dark Matter Direct Detection

Energy Deposition  $\omega$ :

Elastic  
 $O(\text{keV})$

Signals:

Nuclear recoils



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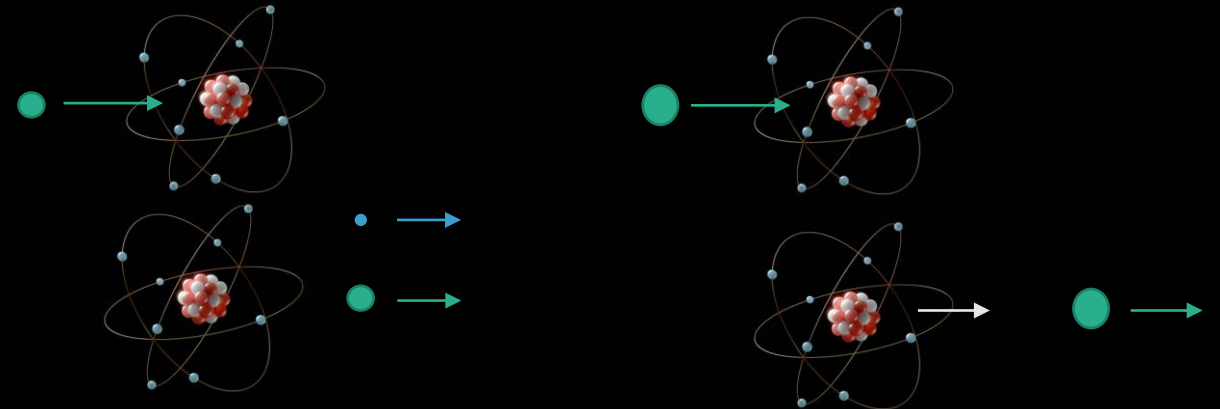
Inelastic  
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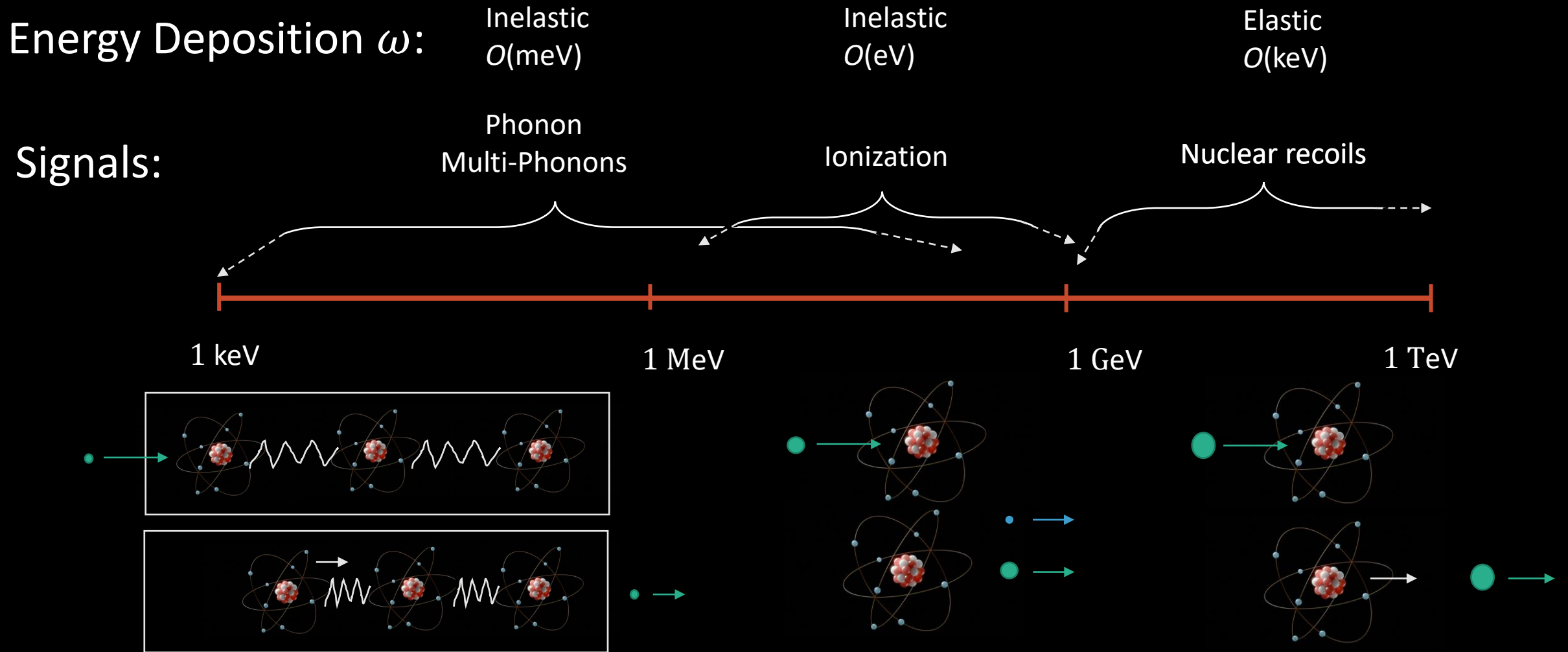
Signals:

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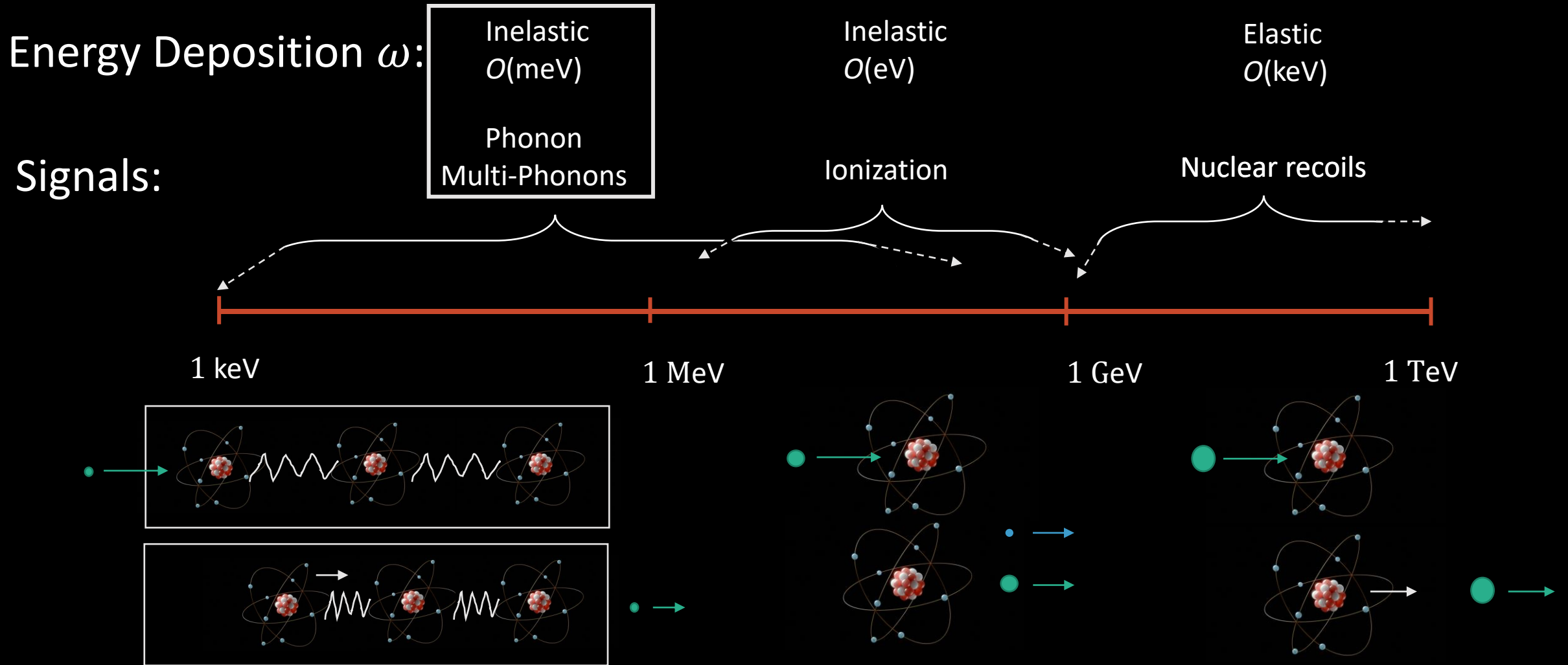
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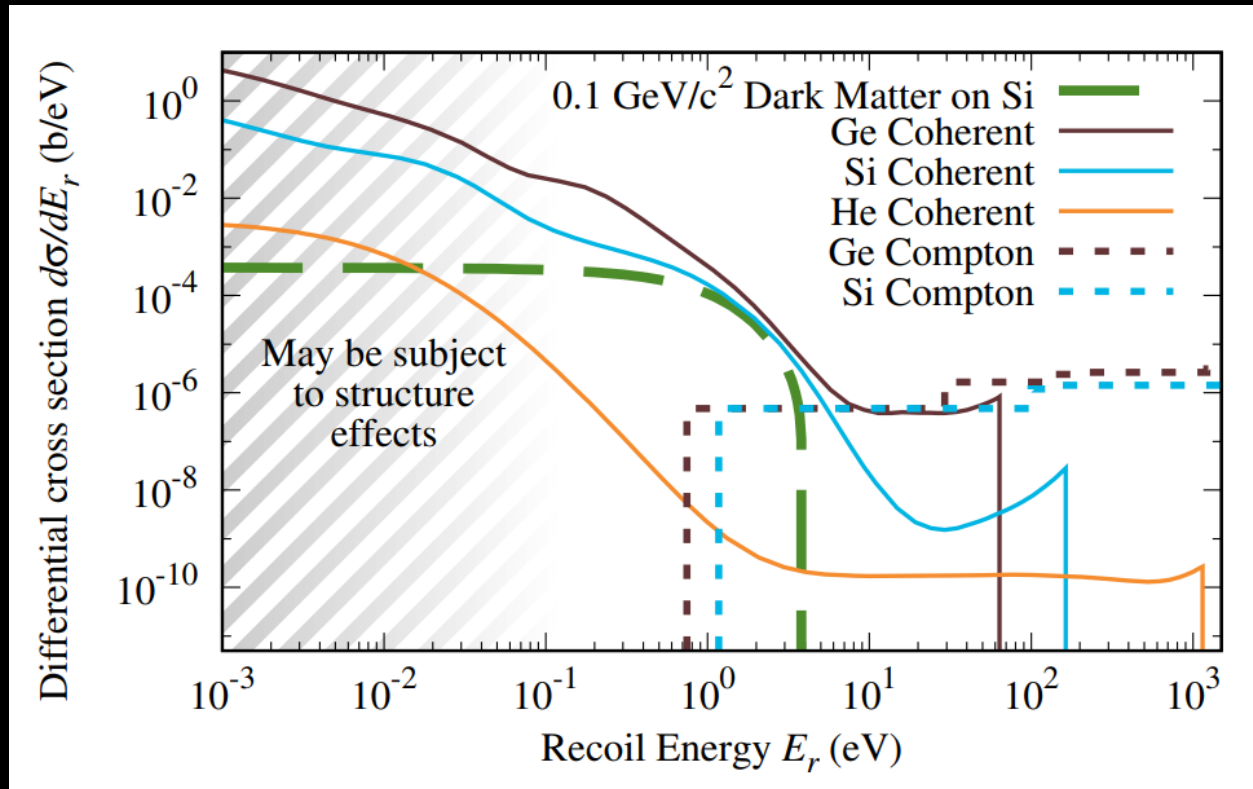
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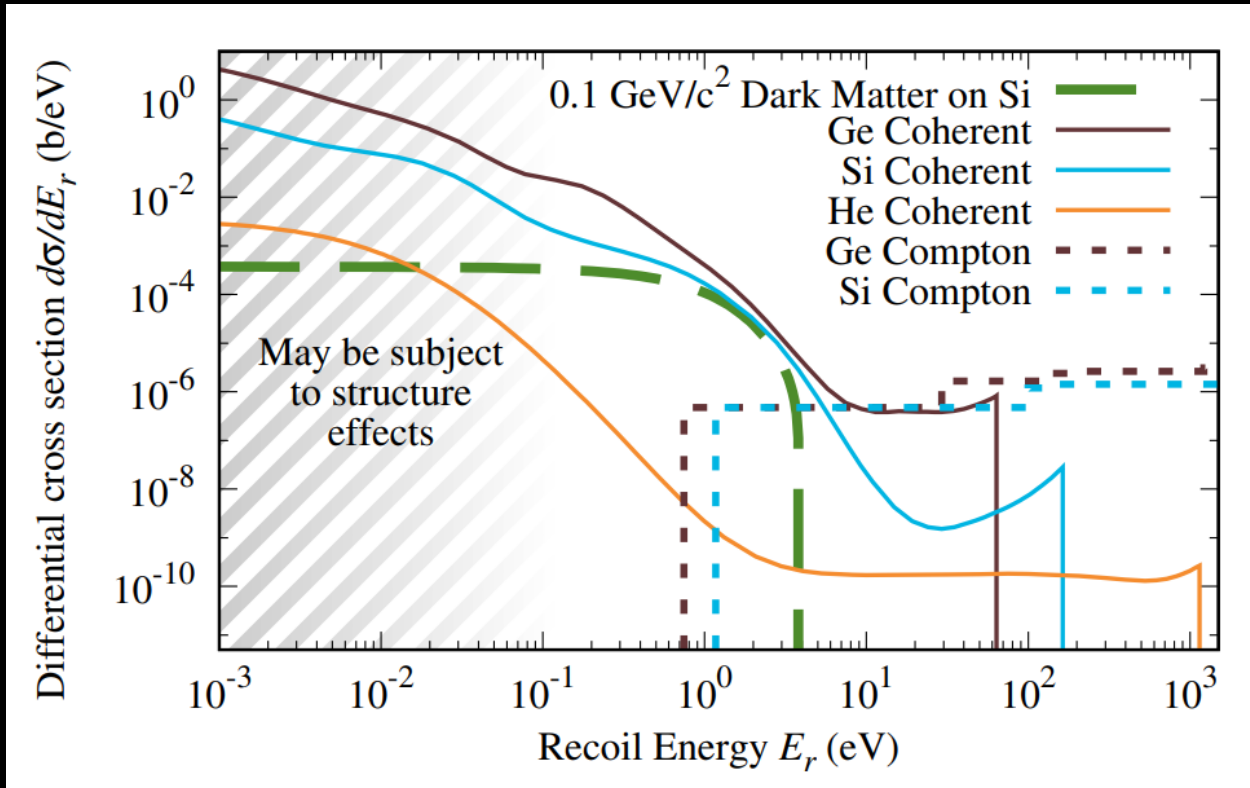


# The Gamma-Ray Background in DM Direct Detection



Alan Robinson, Phs.Rev. D 95,021301(R), 2017

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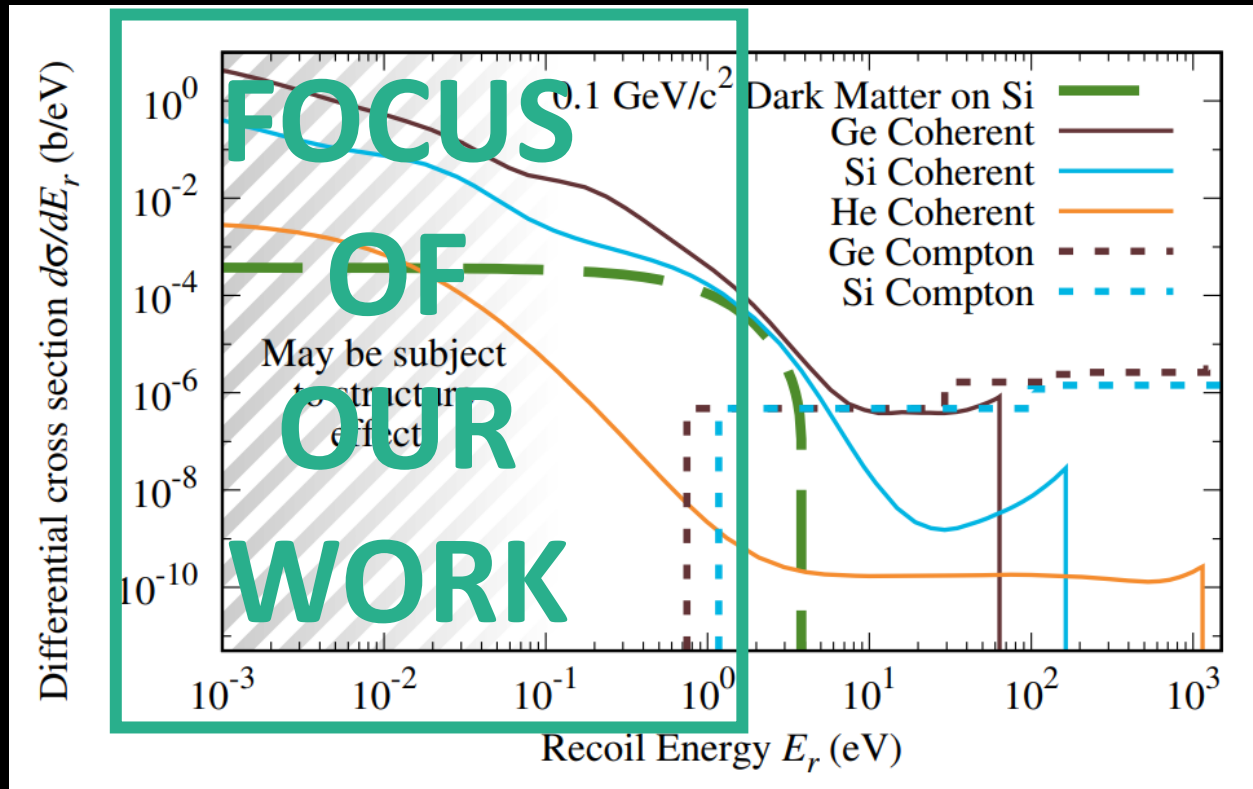
Predicts O(100) events in Ge, Si

assuming 0.04 counts/kg/day/keV

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# The Gamma-Ray Flux in DM Direct Detection

Assumed Gamma-rays based on  
EDELWEISS shielding capabilities

AIP Conference  
Proceedings 1672,  
100002 (2015);

EDELWEISS  
Collaboration 10.1103/PhysRevD.98.082004

$E_\gamma$ [MeV]	Source	$n_\gamma [\times 10^{-18} \text{ cm}^{-3}]$
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$\sim 0.04$  events/kg/day/keV  
of flat Compton background

Table credit Mukul Sholapurkar

# Gamma-Ray Phonon Scattering

The process factors into Photon-Ion scattering  $\times$  material specific response

$$\frac{d\sigma}{d\Omega d\omega} = \frac{d\sigma}{d\Omega}(\mathbf{q}, E_\gamma) \times S(\mathbf{q}, \omega)$$

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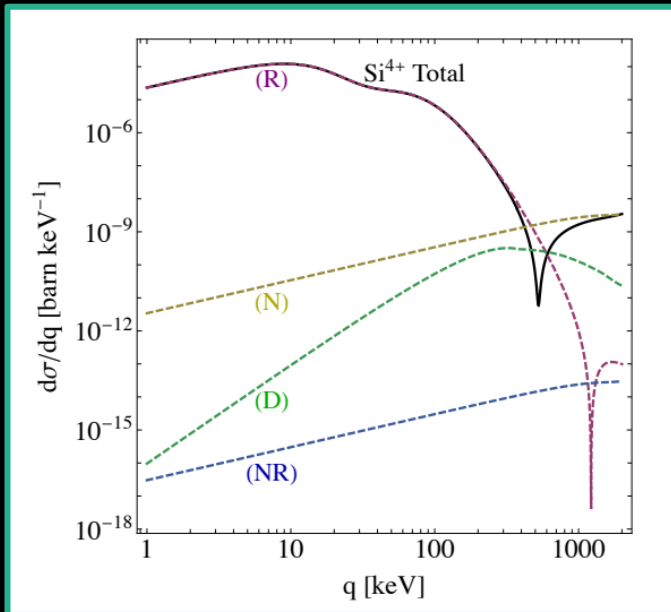
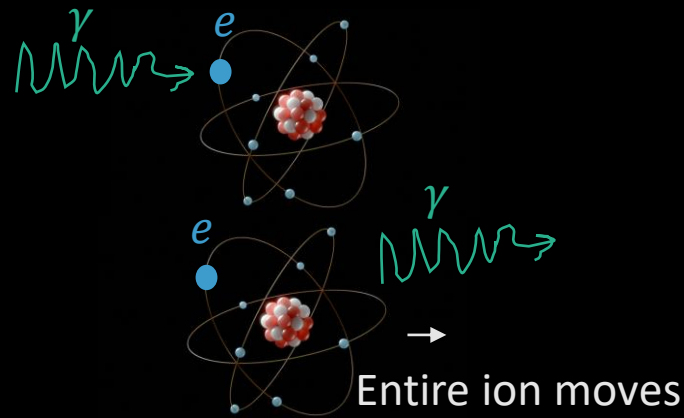


Figure credit Yutaro Shoji

Photon-electron Rayleigh scattering



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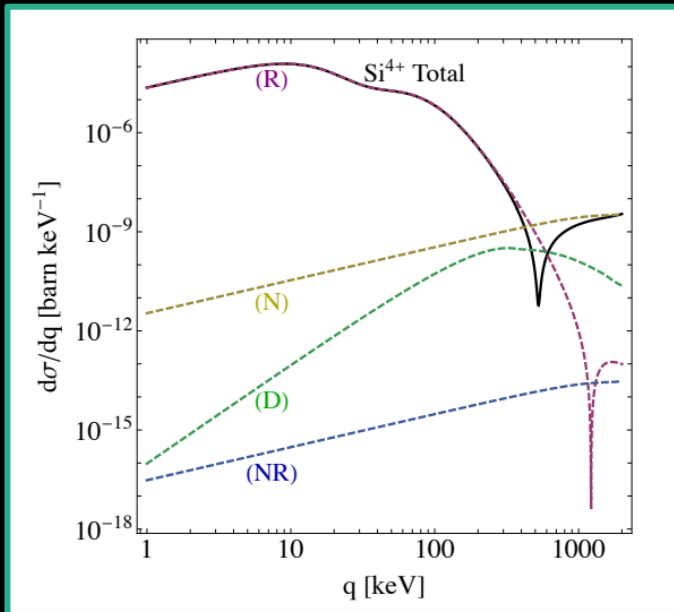
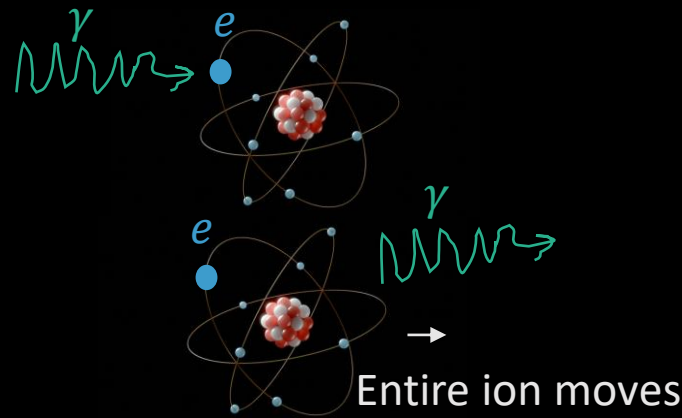


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Photon-electron Rayleigh scattering



$$\frac{d\sigma}{d\Omega}(\mathbf{q}, E_\gamma) = \frac{q^2}{E_\gamma^2} \frac{\alpha^2 \pi}{m_e^2} \times \left( 1 + \left( 1 - \frac{q^2}{2E_\gamma^2} \right)^2 \right) |g(q)|^2$$

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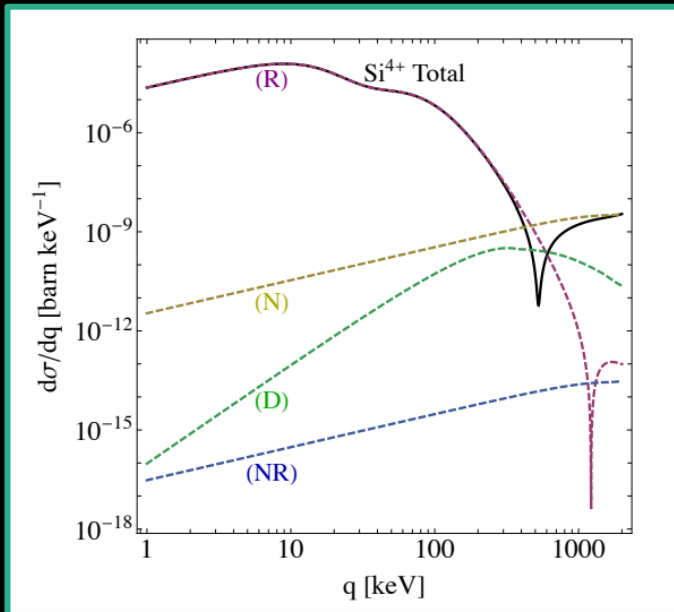
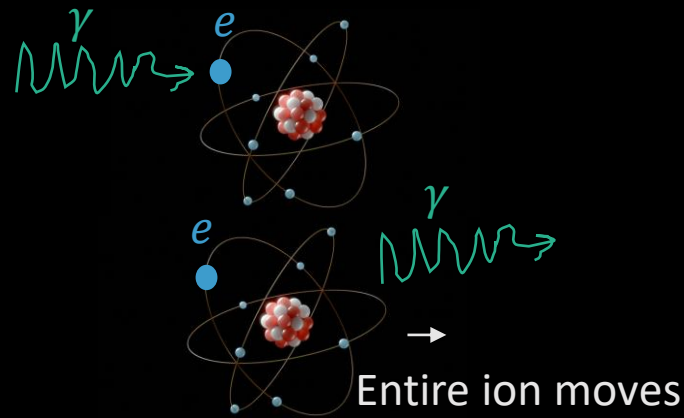


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Kim V. Berghaus, YITP

Photon-electron Rayleigh scattering



EXCESS 2022 Workshop

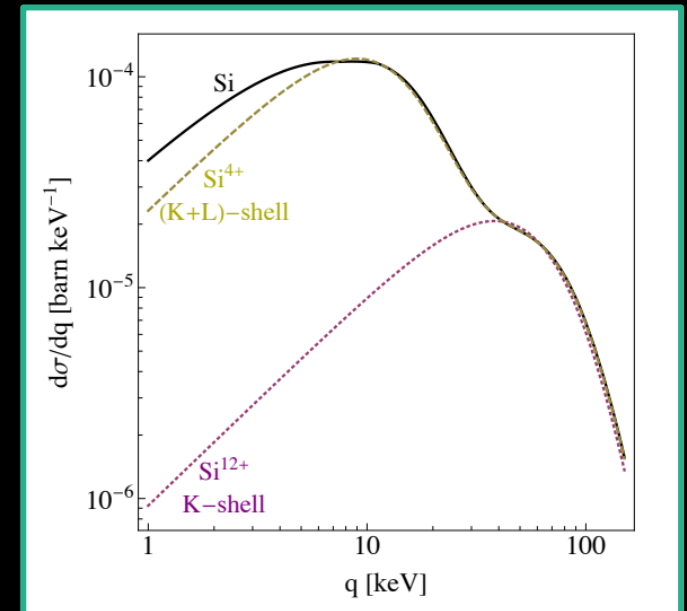


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arXiv 2112.09702

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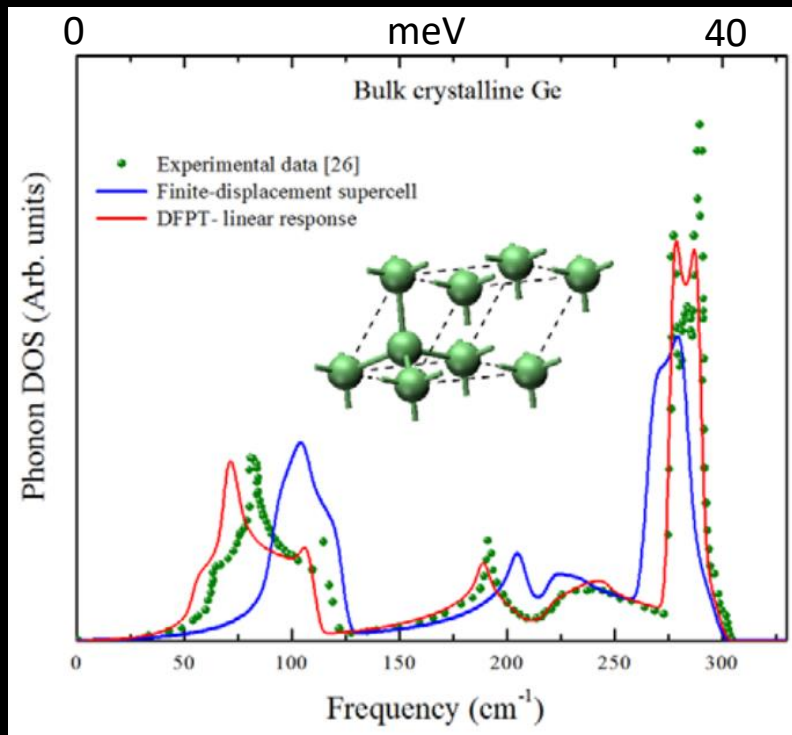
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Same operator appears in  
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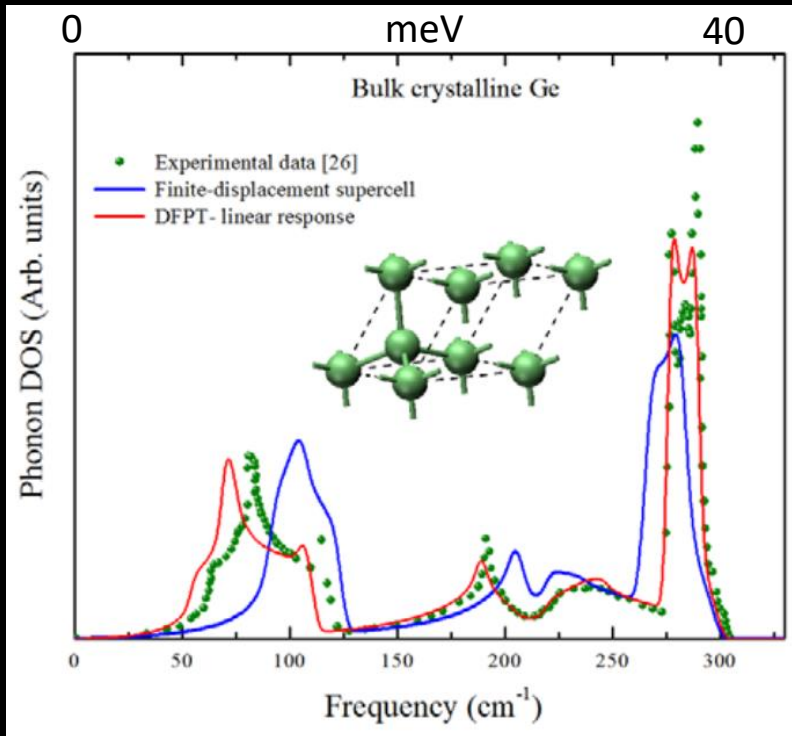
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2013 Apr 22;18(4):4776-85. doi:  
10.3390/molecules18044776. PMID:  
23609626; PMCID: PMC6269924

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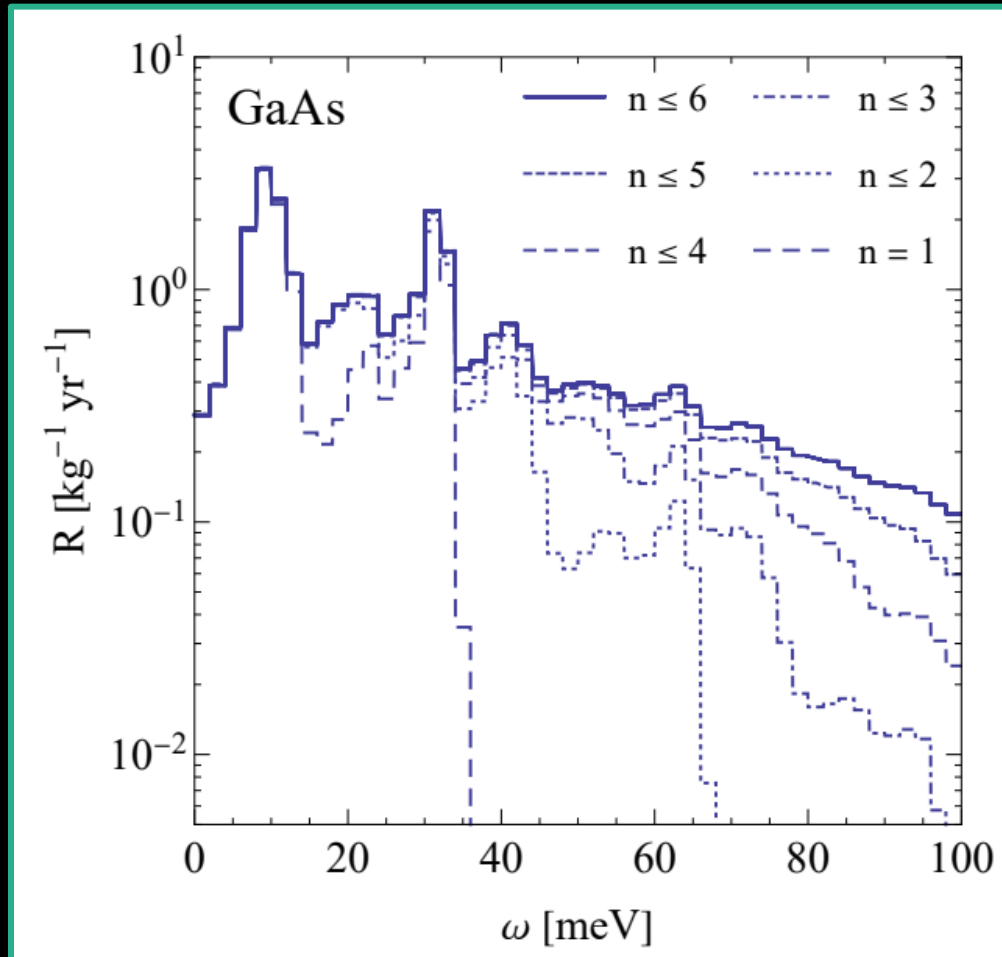
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$$S^1(\mathbf{q}, \omega) \sim e^{-\frac{q^2}{M \bar{\omega}_{ph}}} \frac{q^2}{2M} \frac{F(\omega)}{\omega}$$

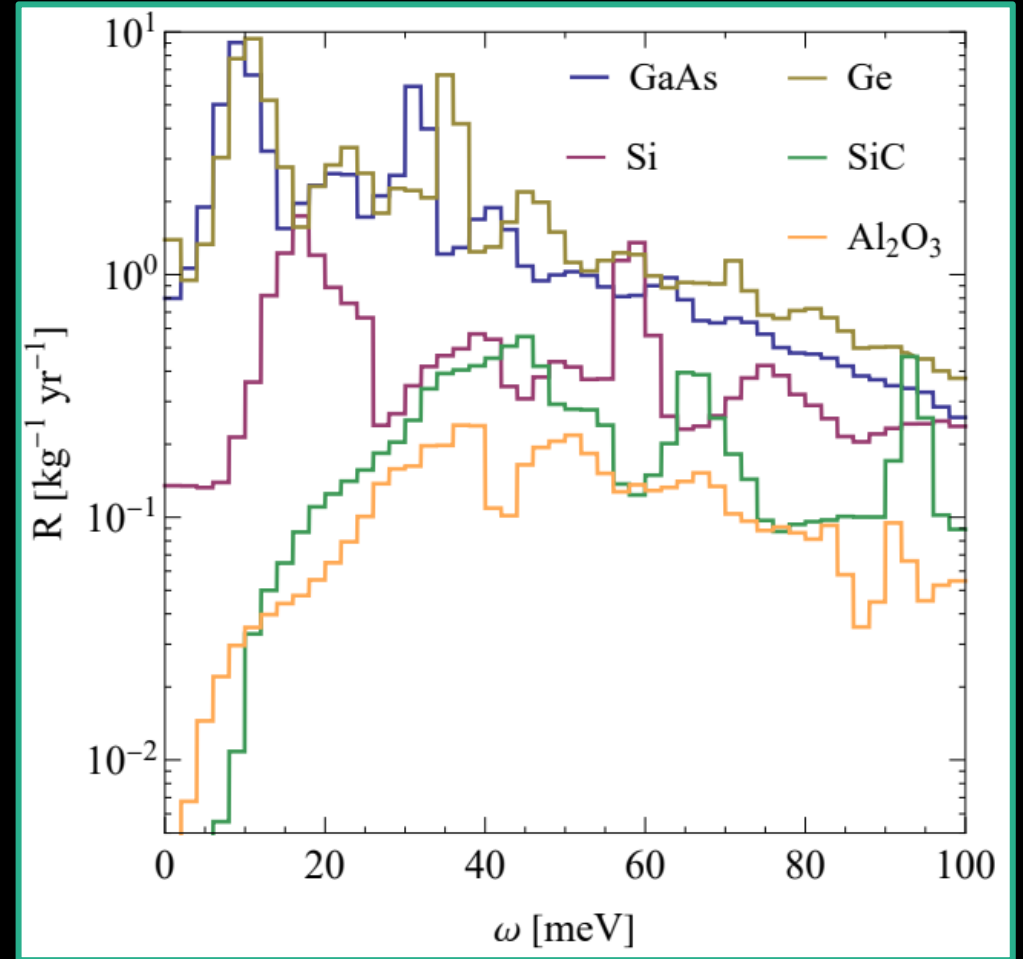
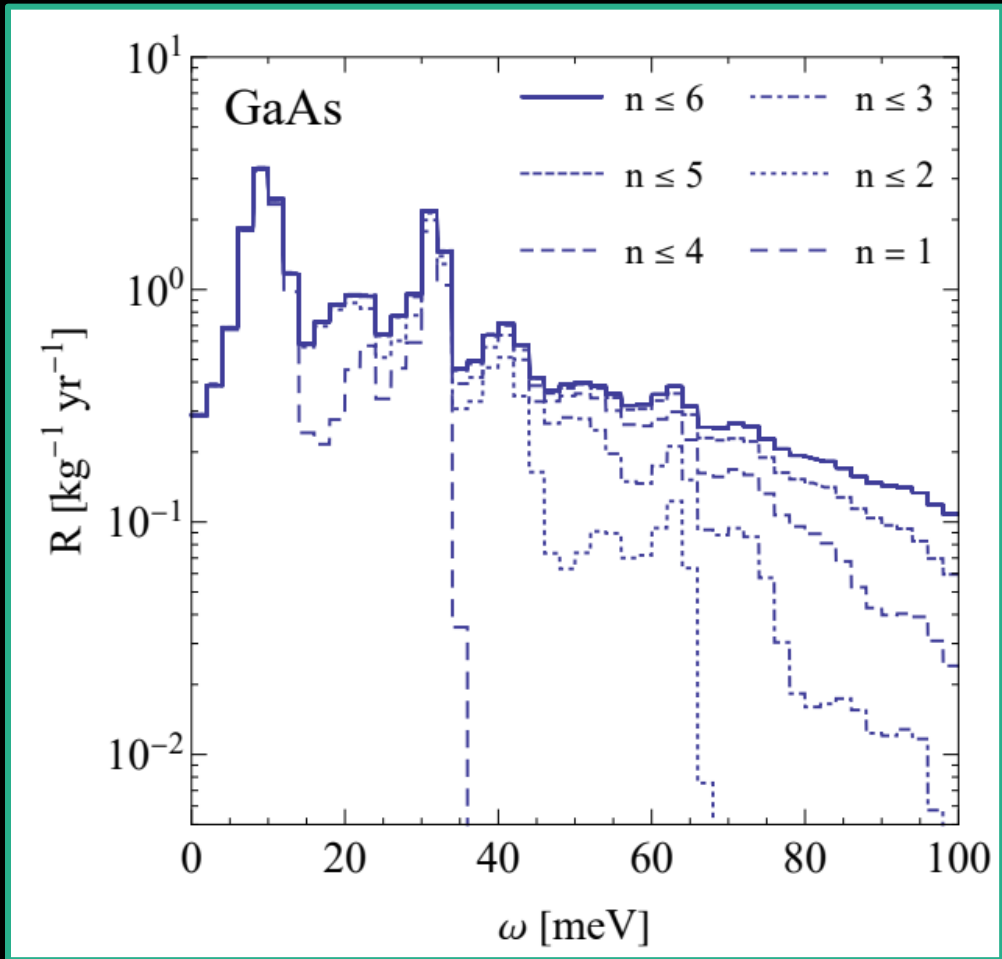
$$S^n(\mathbf{q}, \omega) \propto S^1 S^{n-1}$$



# Phonon Background Rates from Gamma-Rays



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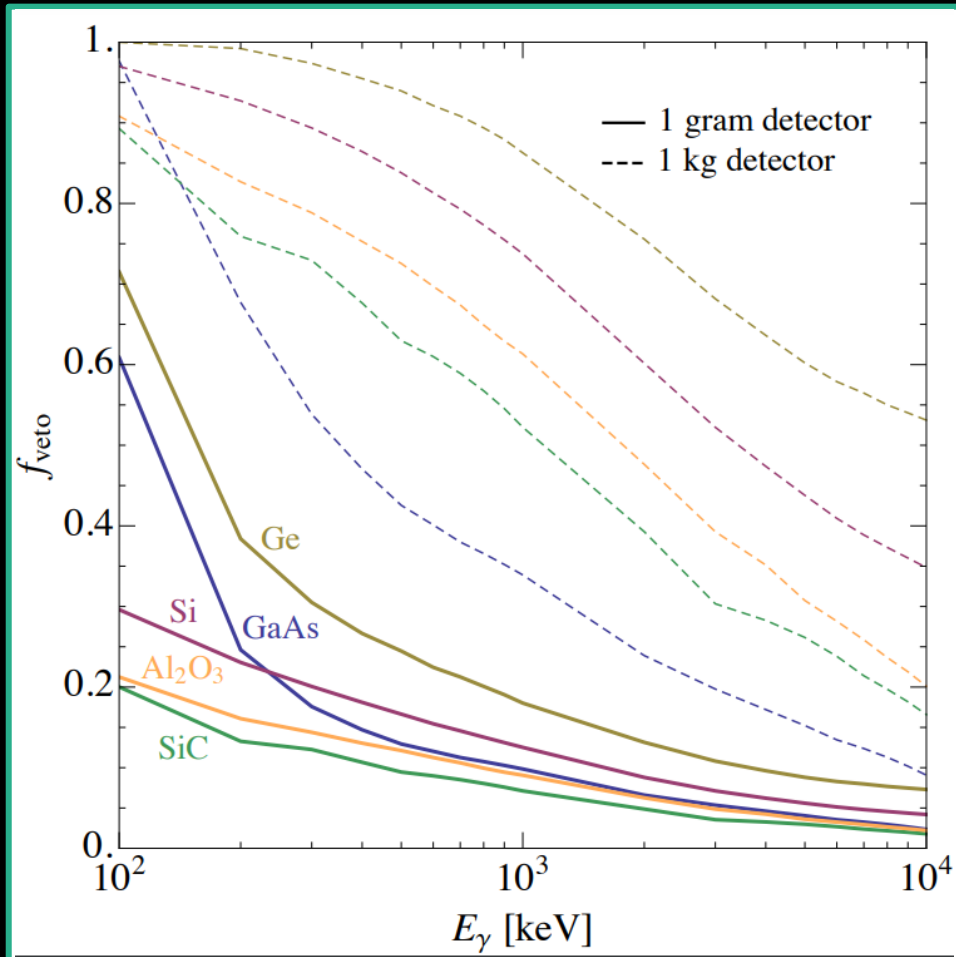
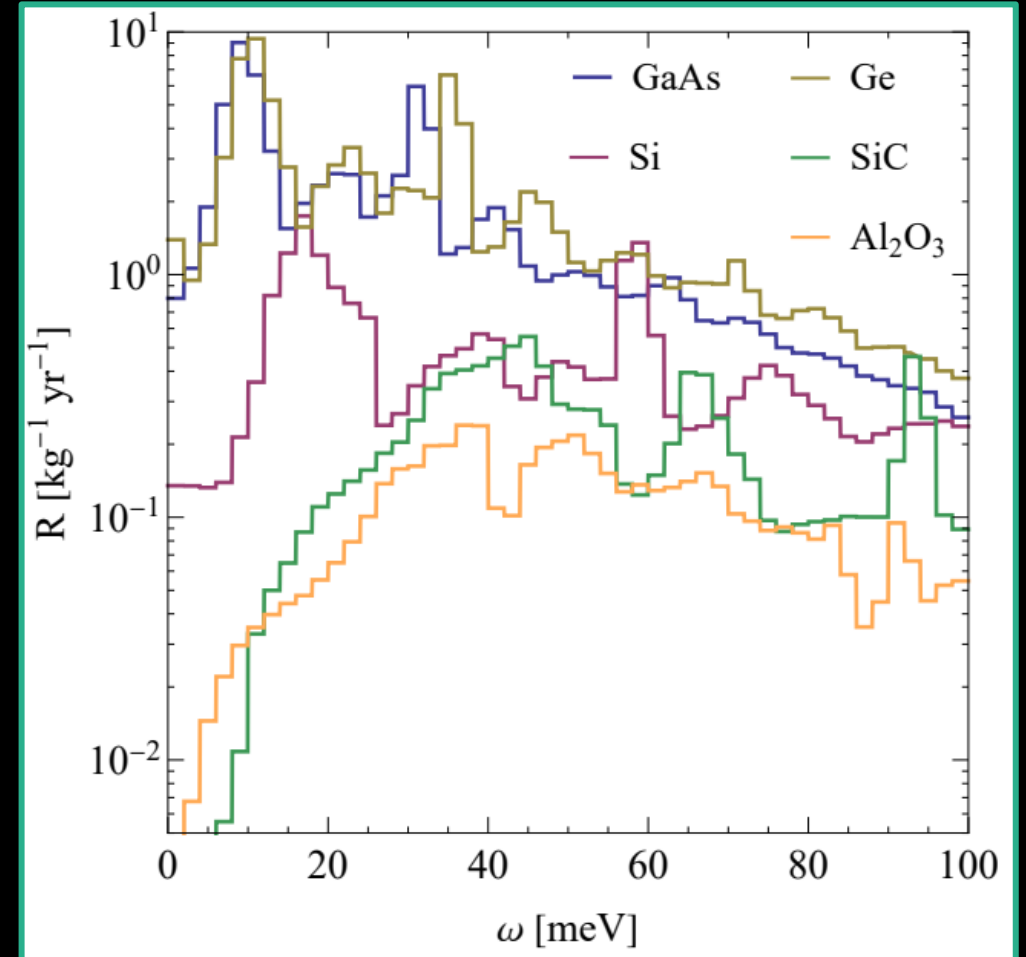
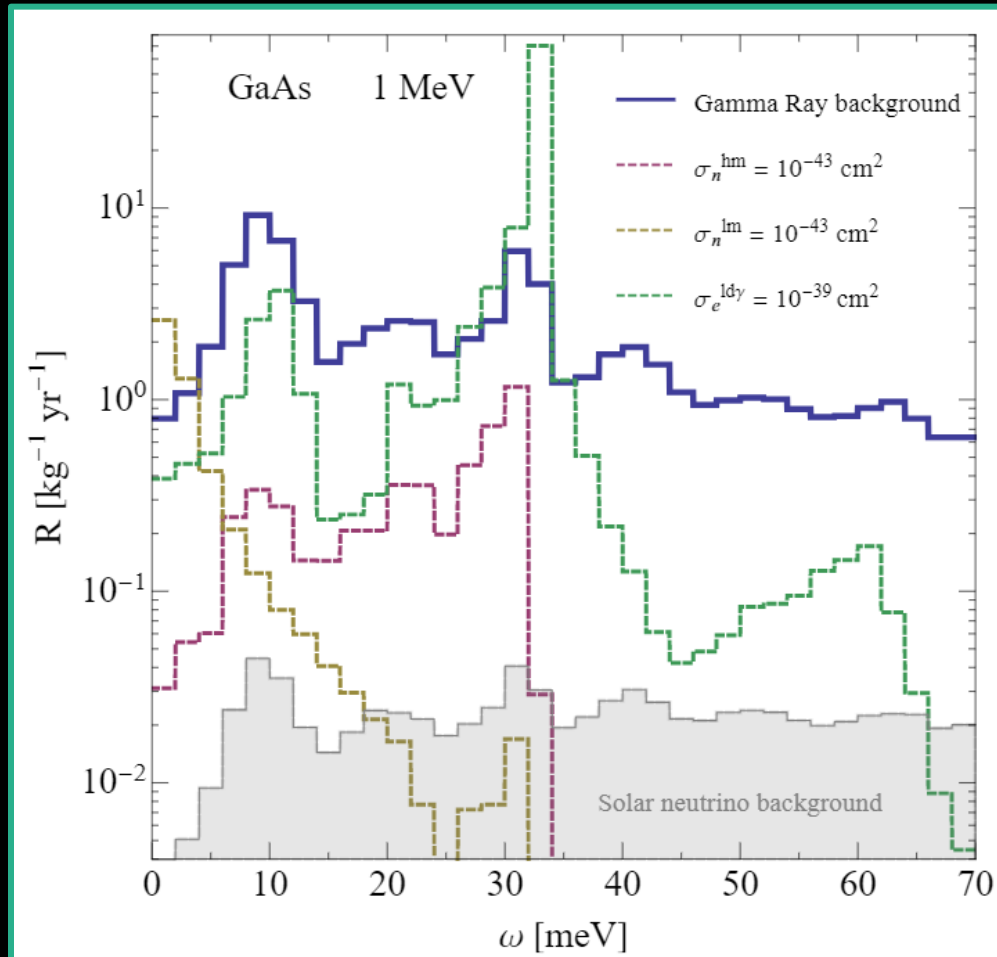


Figure credit to Mukul Sholapurkar



# Comparison to Signal Rates

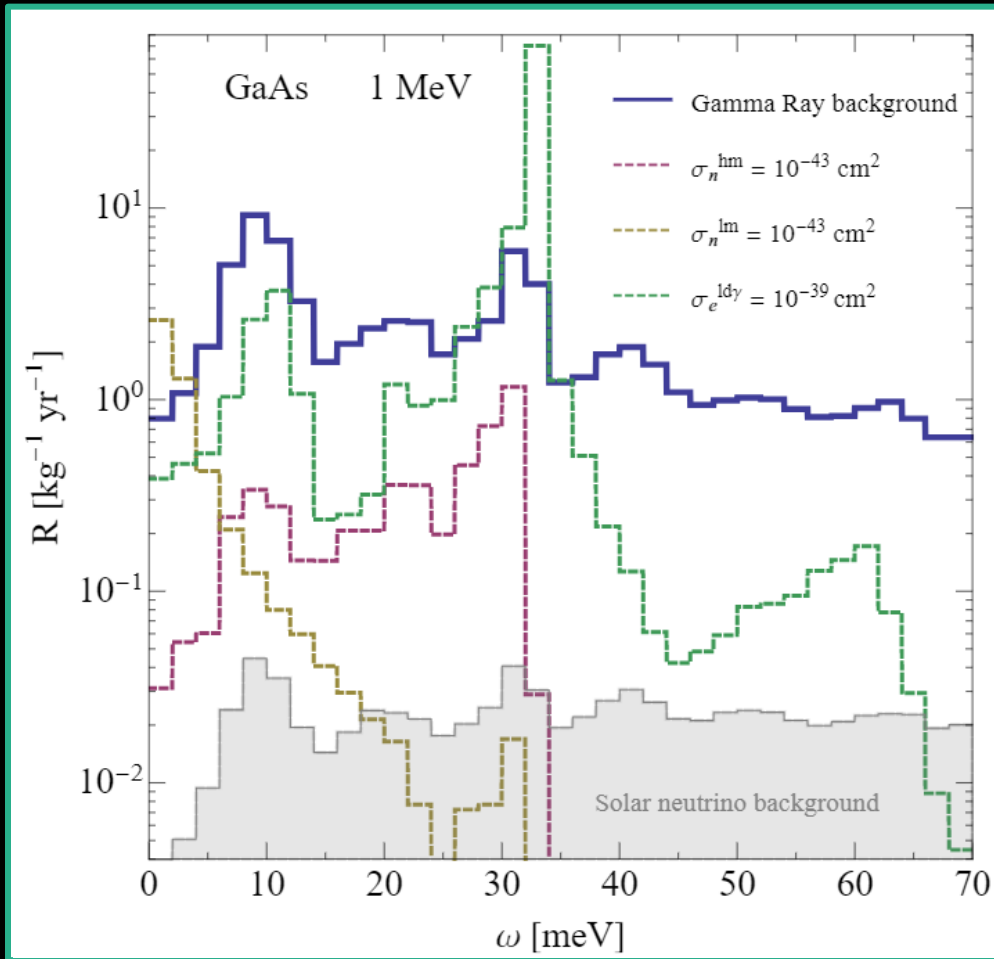


**hm**: heavy scalar mediator (Trickle et al., 2021)

**lm**: light scalar mediator (Trickle et al., 2021)

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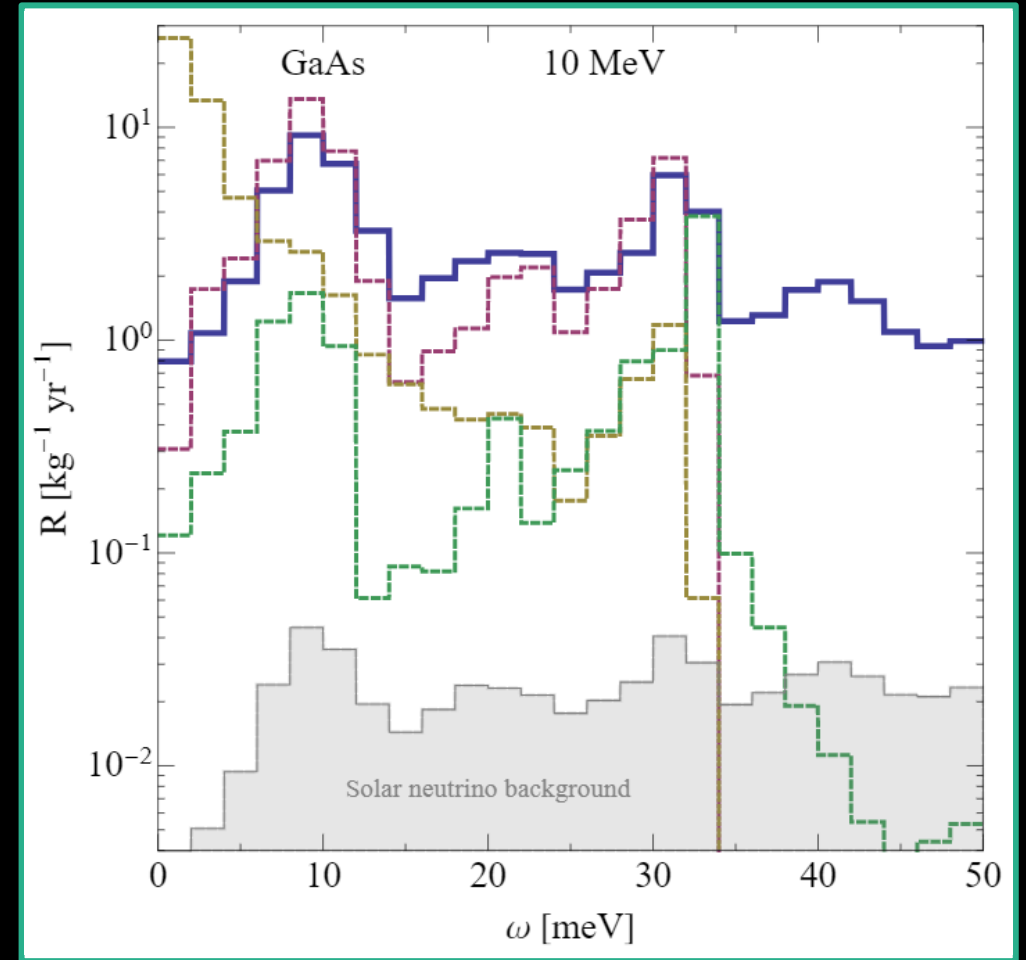
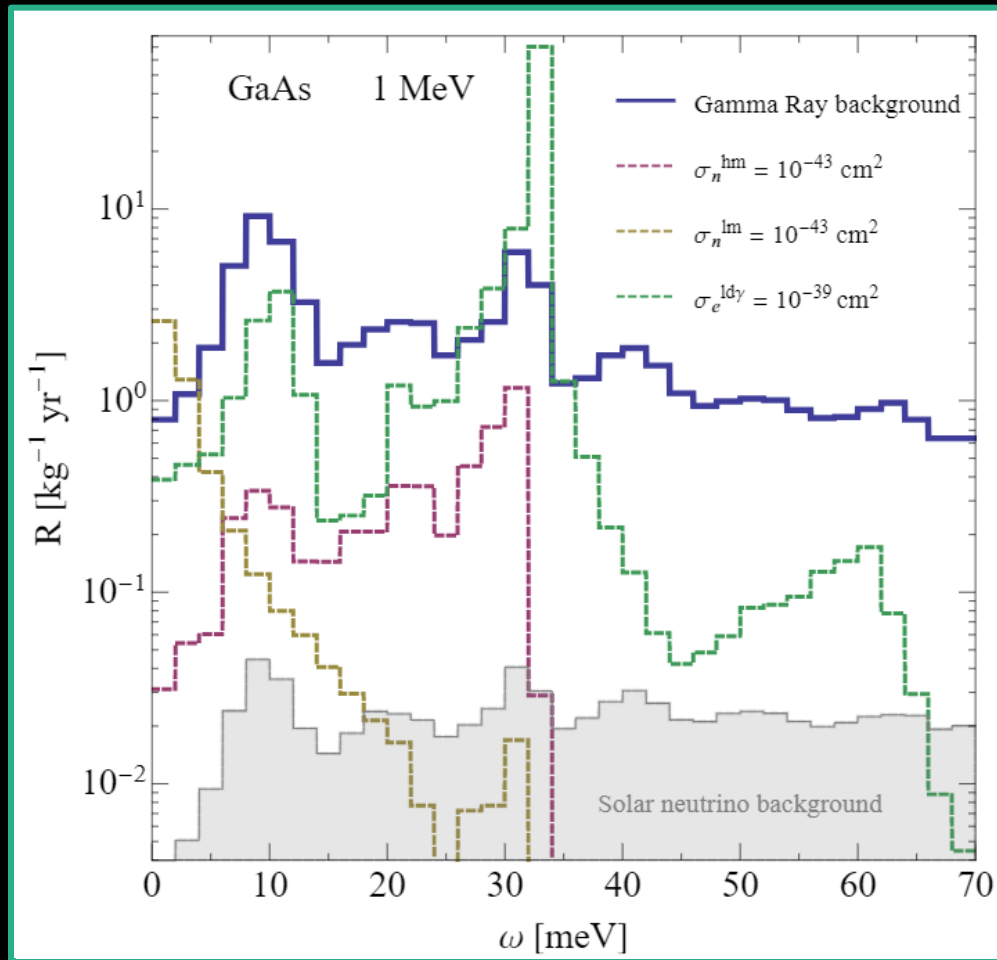
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**hm** and **lm** have the same  $S(\mathbf{q}, \omega)$  as the background

**ld $\gamma$**  model can polarize materials (different  $S(\mathbf{q}, \omega)$ )

**lm** gets enhancement at small  $\omega$  due to  $\frac{d\sigma}{dq} \propto \frac{q_0^4}{q^4}$

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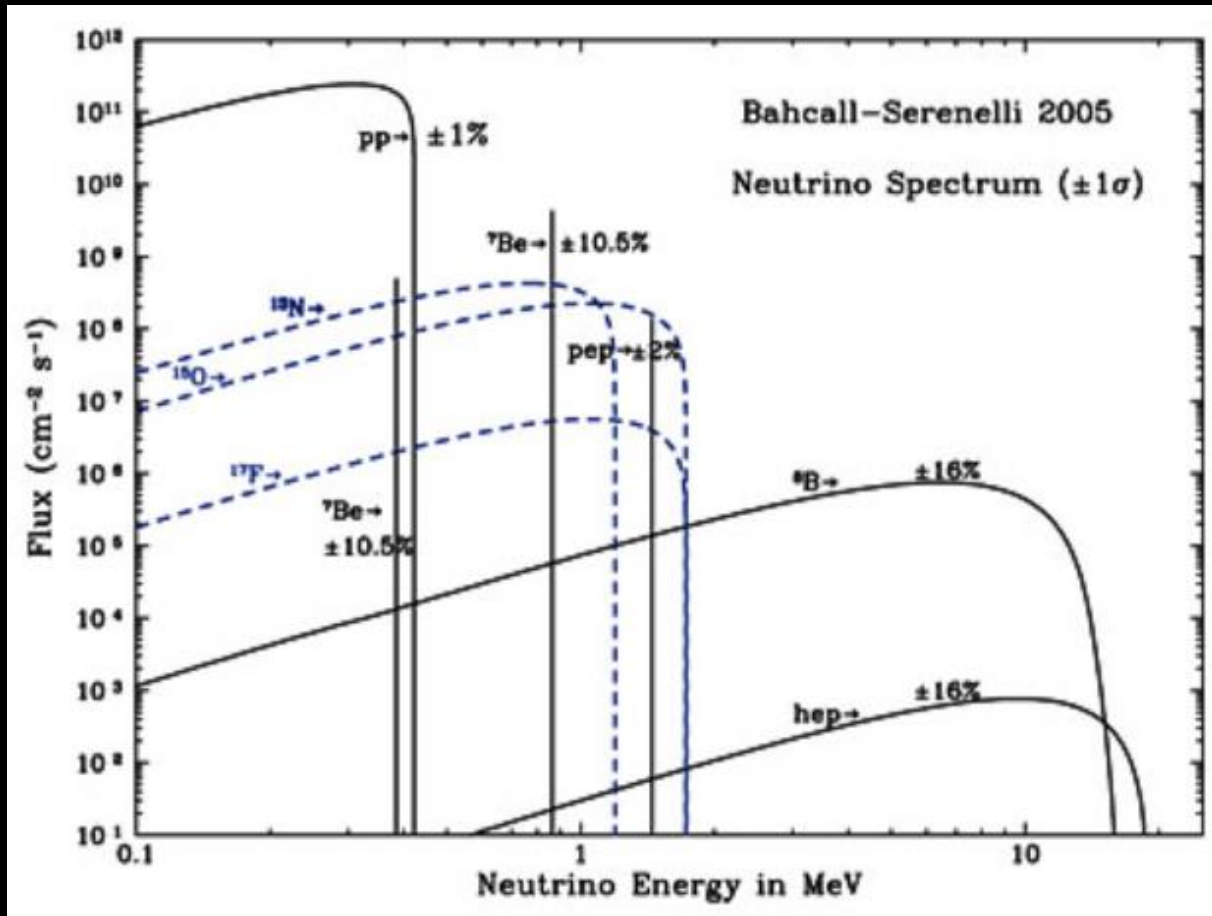
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- Calculations available publicly on Github: [https://github.com/KBerghaus/phonon\\_background](https://github.com/KBerghaus/phonon_background)

**Thank you!**

# Backup

# Solar Neutrino Flux

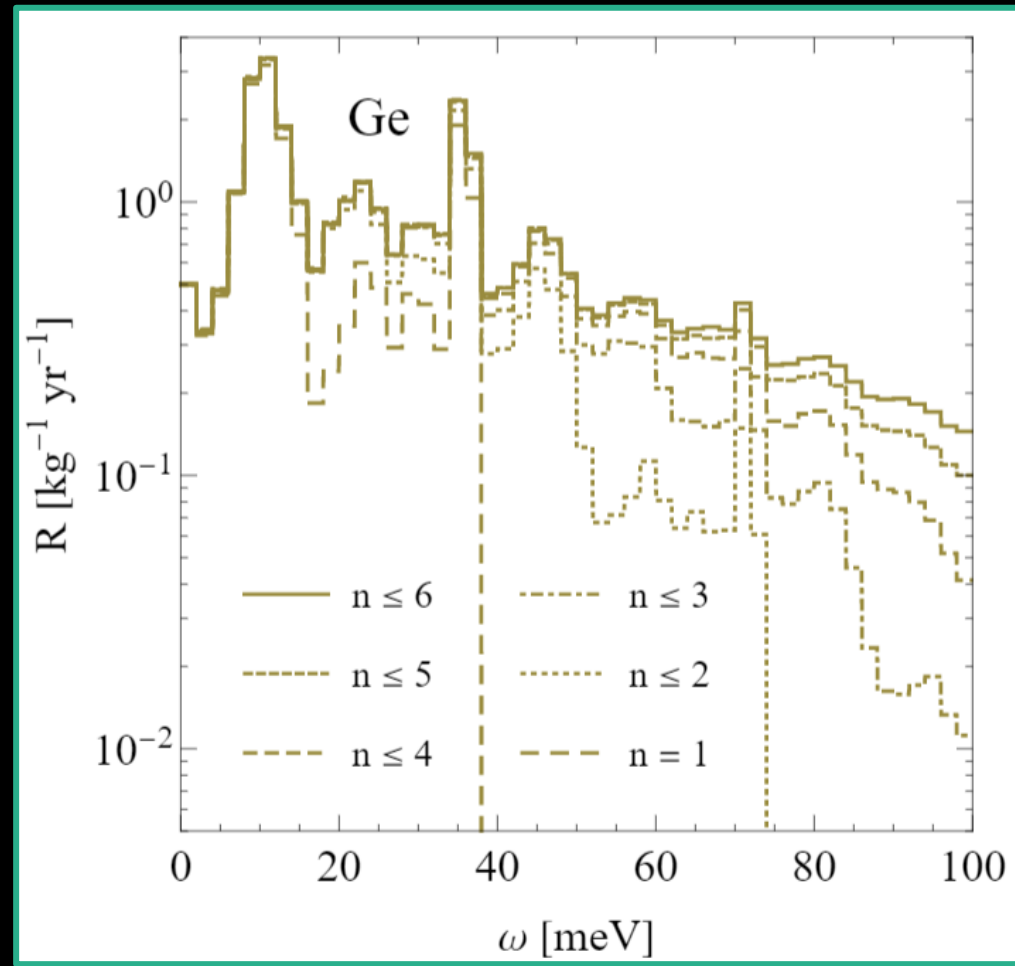
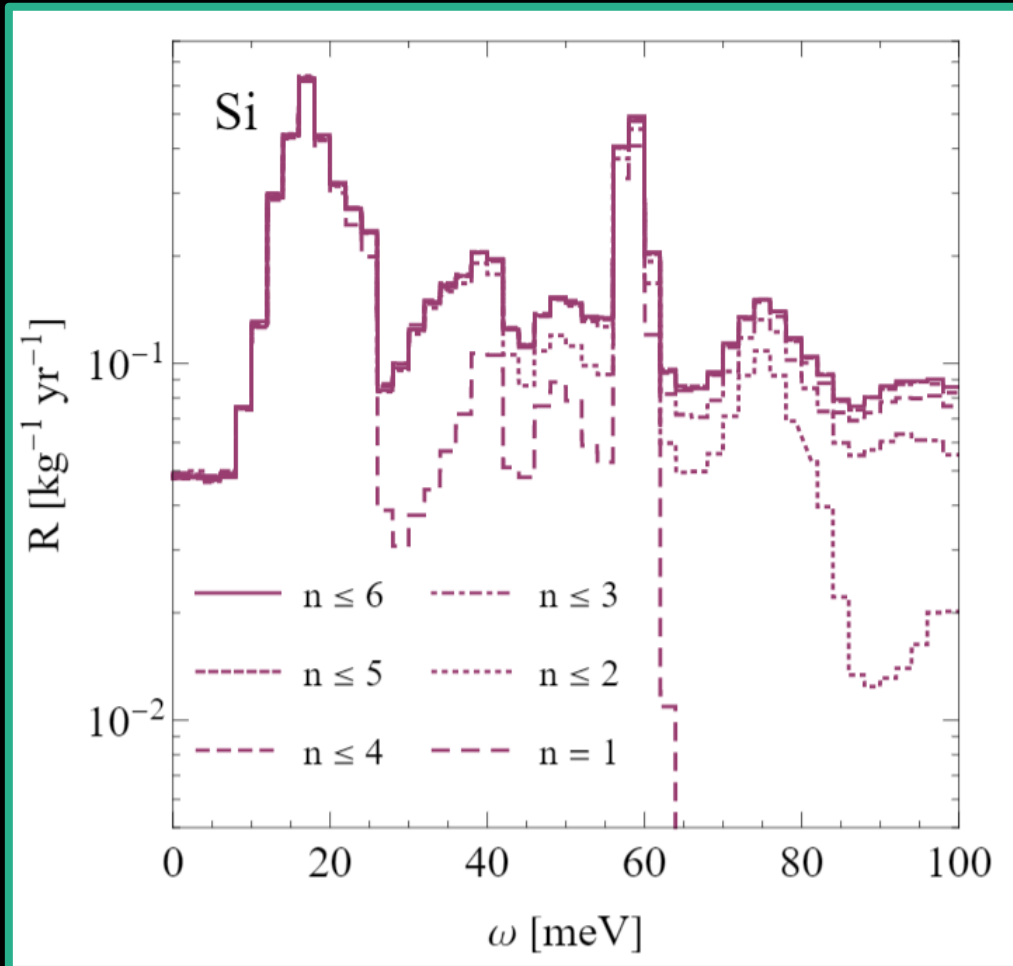


Bahcall, Serenelli, and Basu, *ApJ*, 621, L85 (2005)

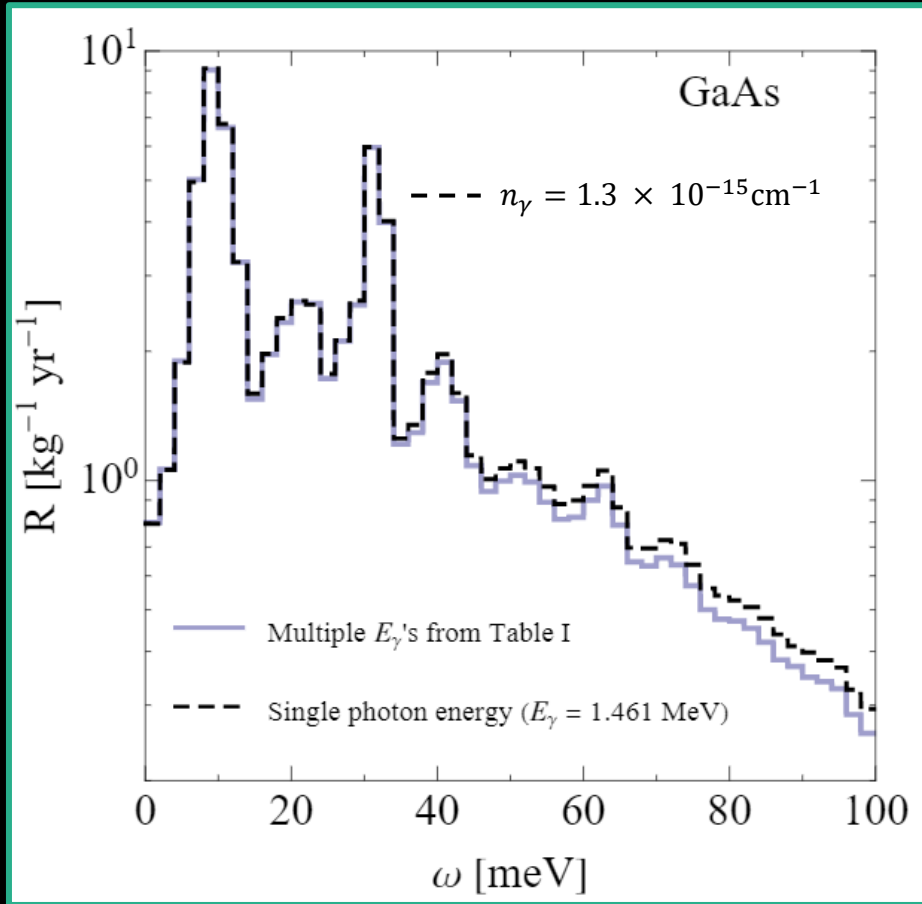
$$\frac{d\sigma_{N\nu}}{d\Omega}(\mathbf{q}, E_\gamma) = \frac{G_F^2}{4\pi} q Q_W^2 \left(1 - \frac{q^2}{4E_\gamma^2}\right) |F(q)|^2$$

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