

# DarkSide-50

## Characterisation of the LAr ionization response in the keV regime

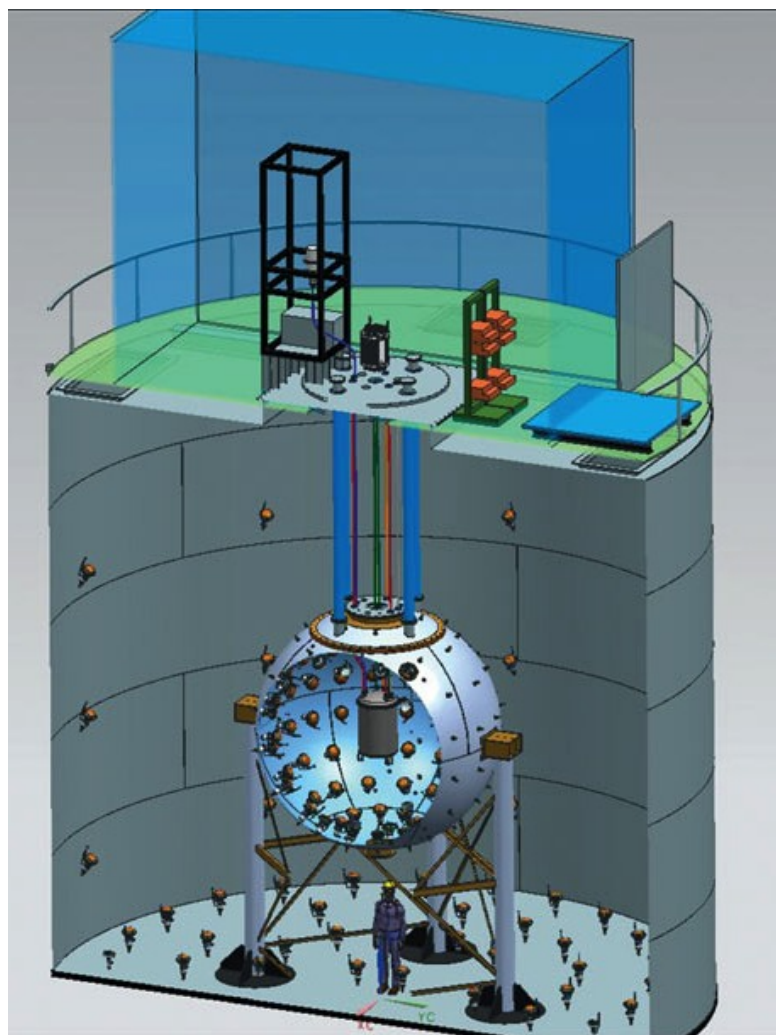
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*Julie Rode (LPNHE/APC) on behalf of the DarkSide experiment*  
16/02/2022

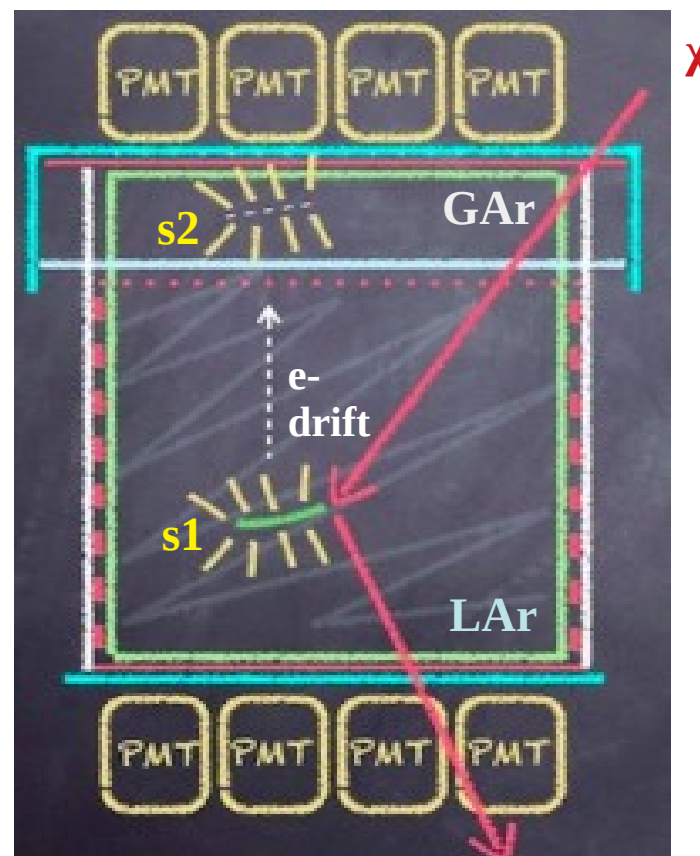
# DarkSide-50 experiment

- **50 kg** dual-phase Liquid **Argon** TPC
- Using Underground Argon: **depleted in  $^{39}\text{Ar}$**
- In a **30 tonnes** borated liquid scintillator **neutron veto**
- In a **1000 tonnes Water** Cherenkov Veto
- **Underground** in Gran Sasso National Lab, Italy



DOI:10.1088/1748-0221/12/12/P12011

## Dual-phase LAr TPC



<http://aris.in2p3.fr/doc/aris.pdf>

## DarkSide-50

- s1 light collection efficiency:  $0.16 \pm 0.01$  → *low efficiency*
- s2 yield:  $23 \pm 1$  pe/e- → *amplification*

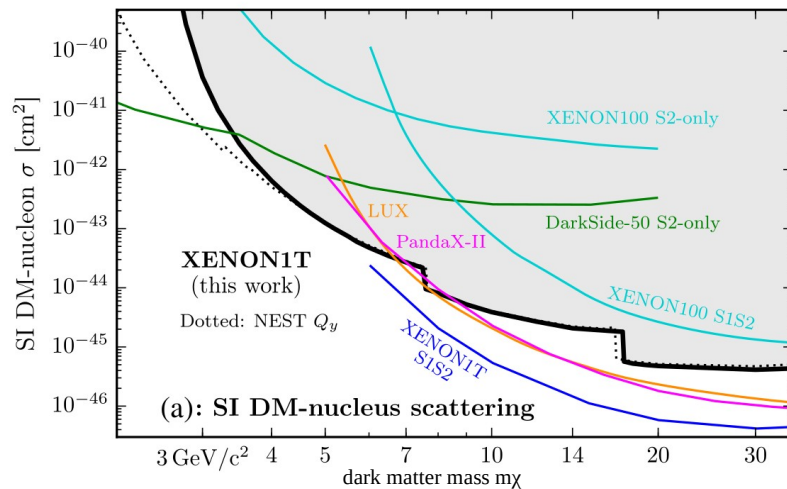
DOI: 10.1088/1748-0221/12/10/P10015

keV regime accessible using s2 only

# Motivation

**Lower detection threshold, higher sensitivity to light dark matter candidates:**

- WIMP-nucleon interaction with/without Migdal effect
- WIMP-electron interaction
- Sterile neutrinos
- Axion-like particles



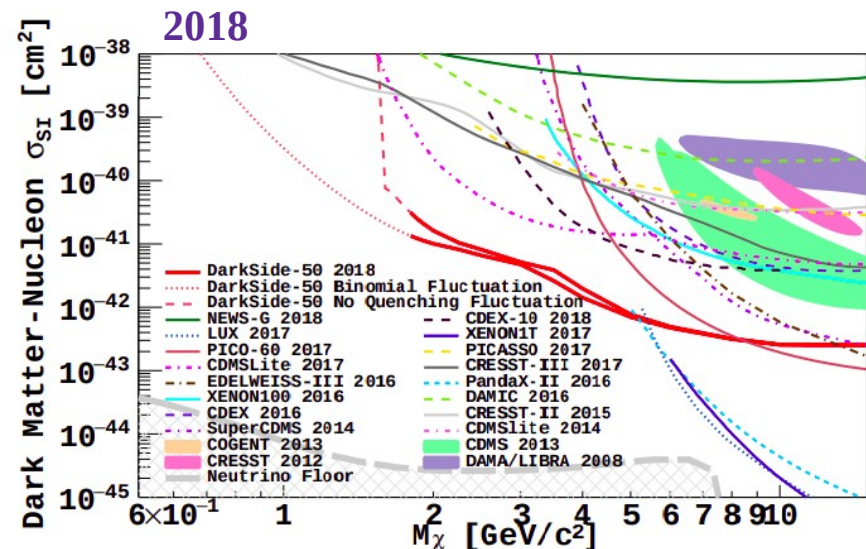
DOI: 10.1103/PhysRevLett.123.251801

**Liquid argon detectors:**

- Massive
- Radiopure
- High scintillation yield
- High ionization yield
- Low electron mobility
- Argon mass  $\ll$  Xenon mass
- Higher recoil energy (transferred momentum) wrt Xe at low energy

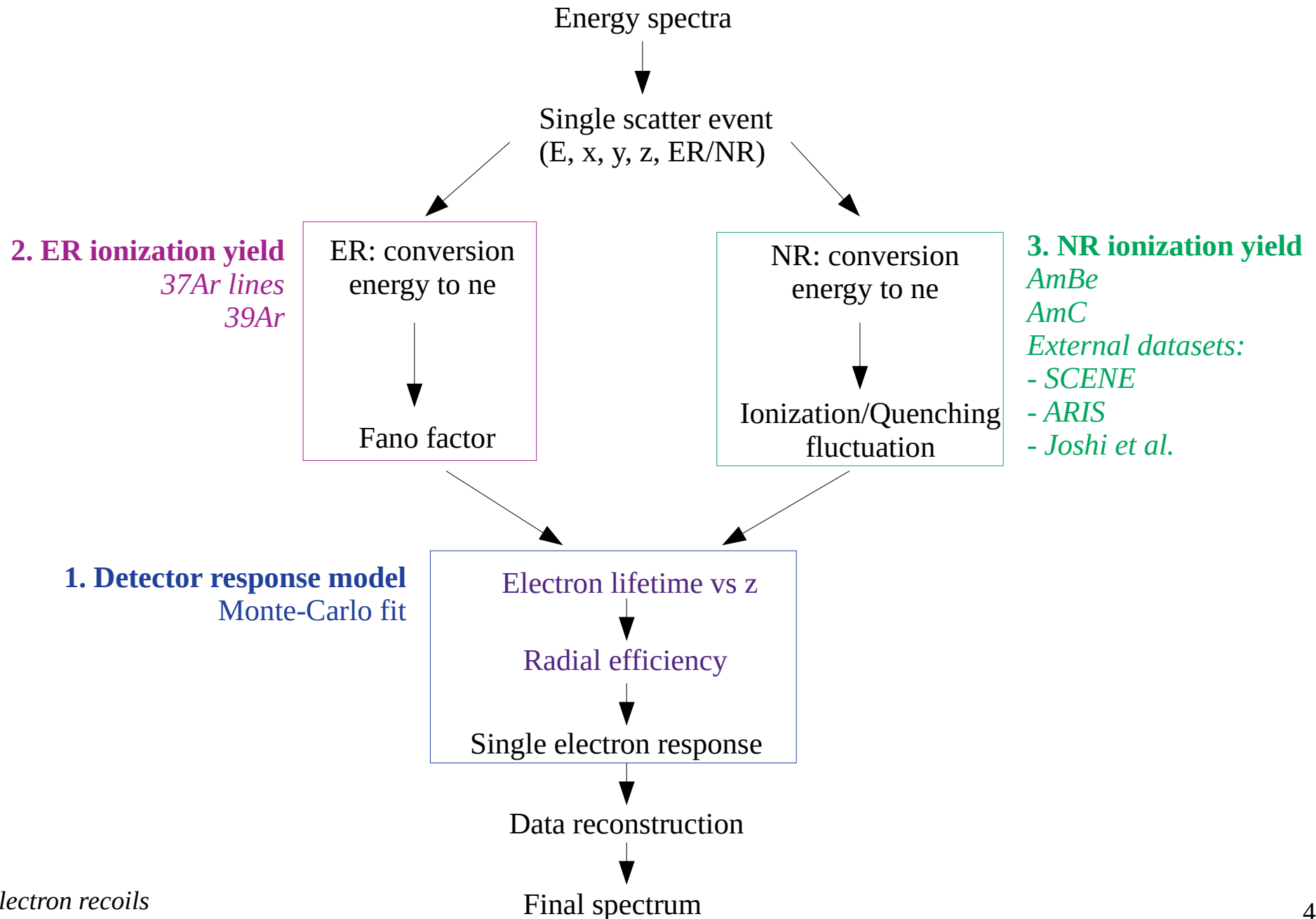
**Noble liquid detectors:**

- Efficient background discrimination
- Massive target
- High scintillation yield



DOI:10.1103/PhysRevLett.121.081307

# Response model



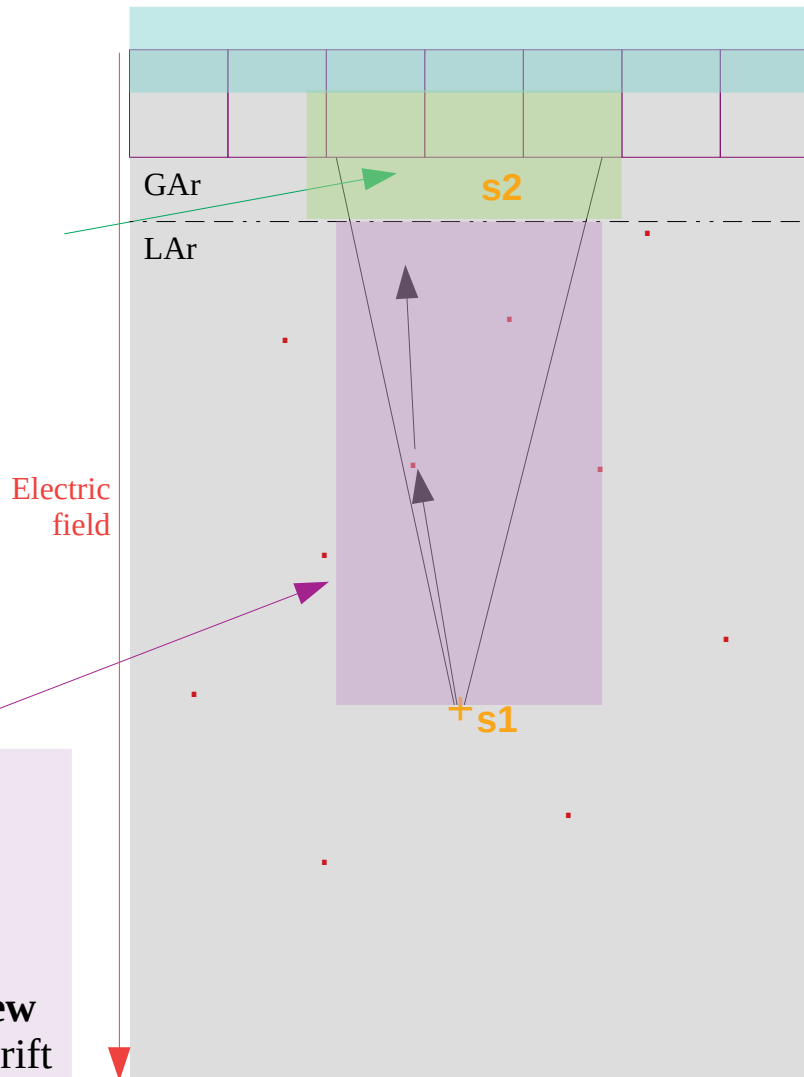
*ER: electron recoils*  
*NR: nuclear recoils*  
*ne: number of electrons*

# Instrumental effects

## Monte-Carlo modelling of the detector response in S2

### Resolution dominated by:

- Electroluminescence fluctuations
- PMT charge response



Electric field

### Electron lifetime

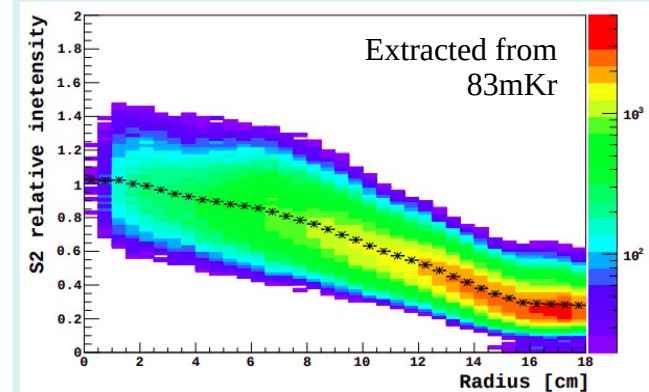
- Capture from **impurities** along the drift
- Measured  $> 10$  ms
- Maximal distortion of **a few percent** wrt the maximal tdrift

### Radial distortion

Low S2 pulse: XY reconstruction inefficient



XY estimator: top PMT position with max light fraction



DOI:10.1088/1748-0221/12/10/P10015

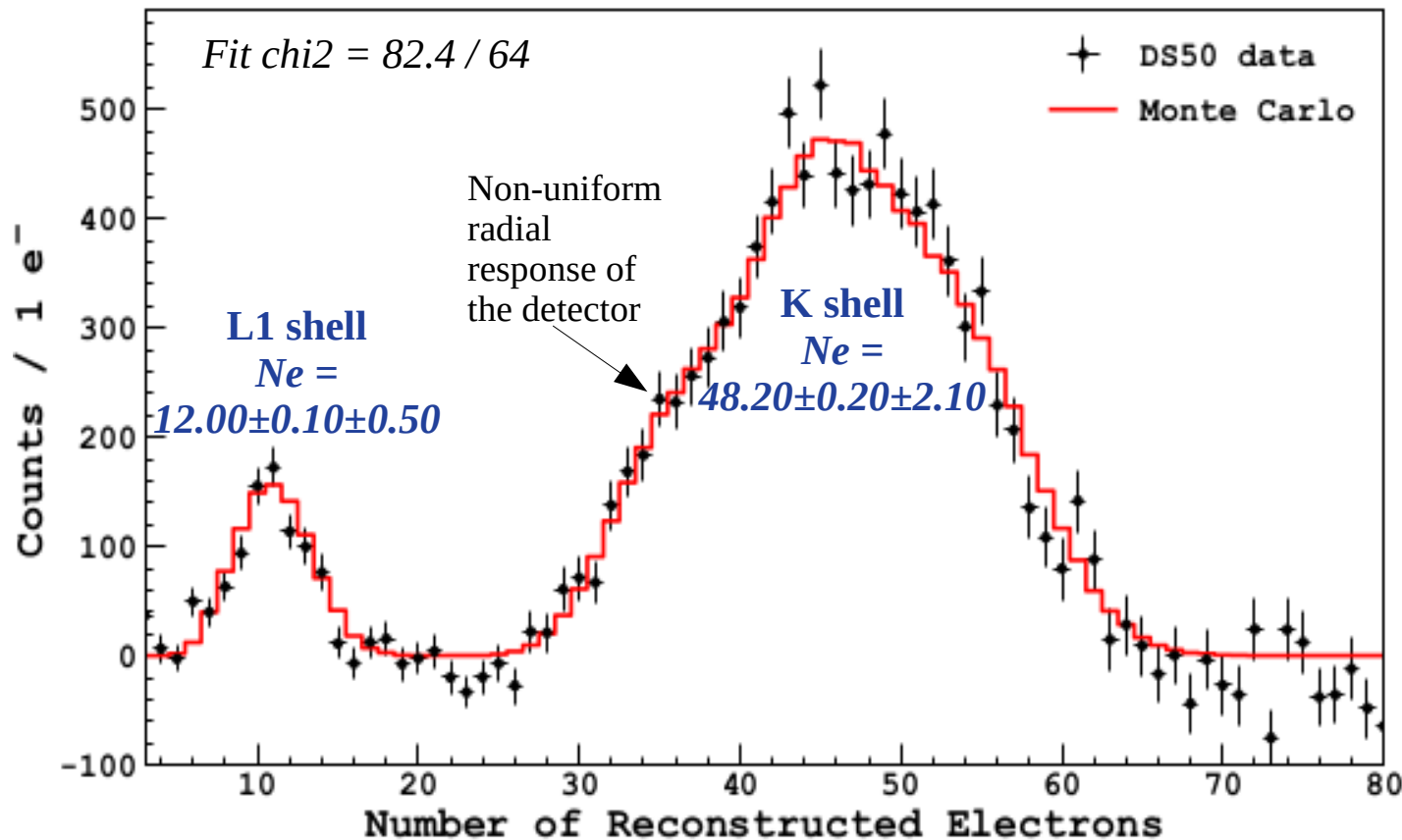
# $^{37}\text{Ar}$ K and L electron capture

## Sample selection

- Subtraction between the first  $\sim 100$  days from the latest  $\sim 500$  days of the UAr campaign
- $^{37}\text{Ar}$  almost entirely decayed in the last  $\sim 500$  days
- Samples normalized by their lifetime

## Fit performed with a **chi2 analysis**

- Free parameters: number of extracted electrons (both lines  $^{37}\text{Ar}$  + Fano factor)



## L1/K shells ratio:

→ measured =  $0.10 \pm 0.01$

→ expected =  $0.103 \pm 0.01$

DOI: 10.1103/PhysRev.120.2196

$0.102 \pm 0.01$

DOI: 10.1080/14786436208212179

$0.098 \pm 0.003$

DOI: 10.1007/BF01333365

## Fitted Fano factor:

→ measured =  $0.10 \pm 0.03$

→ expected =  $0.107$  (Schockley)

$0.116$  (Alkhazov)

DOI: 10.1016/0029-554X(76)90292-5

# 37Ar decay

Single electron capture transition ground state to ground state (half life of 35.01 days)

**Evaluation of emitted cascades** of electrons, X-rays and UV photons with **RELAX software** (EADL2017 library of atomic transition data)

- Atomic relaxation spectra of UV photons, X rays, Auger electrons (primaries)
- Primaries from bound state to bound state transitions for a single initial vacancy in the different sub-shells
- Deterministic propagation of the vacancies up to the valence shell and to the neutralization
- Consideration of atomic configurations

	L1-shell EC			K-shell EC		
Branching Ratio	8.4%		using <b>BetaShape</b> code <a href="https://doi.org/10.1016/j.apradiso.2019.108884">https://doi.org/10.1016/j.apradiso.2019.108884</a>	90.4%		
Total Released Energy	277			2829		
Mean number of primaries <sup>a</sup>	2.8			3.9		
	$\langle N \rangle$	$\langle E \rangle$ [eV]		$\langle N \rangle$	$\langle E \rangle$ [eV]	
K Auger electrons			Fit: $12.0 \pm 0.1(\text{stat.}) \pm 0.5(\text{syst.})$	0.905	2414	
K X-rays				0.095	2634	
L Auger electrons	0.9995	179		→ - 1	1.77	179
L X-rays	0.0005	207		→ neglected	8E-4	188
M Auger electrons	0.96	51		→ - (2±1)	0.35	51
UV photons (E>16 eV)	0.86	25		→ - 1	0.77	25
Undetectable via ionization	2.10	13		→ 0	3.26	13

Lack of model:  
complex event  
topology  
→ **exclude K  
shell**

Ionization electrons at 179eV: **8.2±1.3**



# Electron Recoil Ionization yield

Ionization yield per unit of ER energy following the custom model:

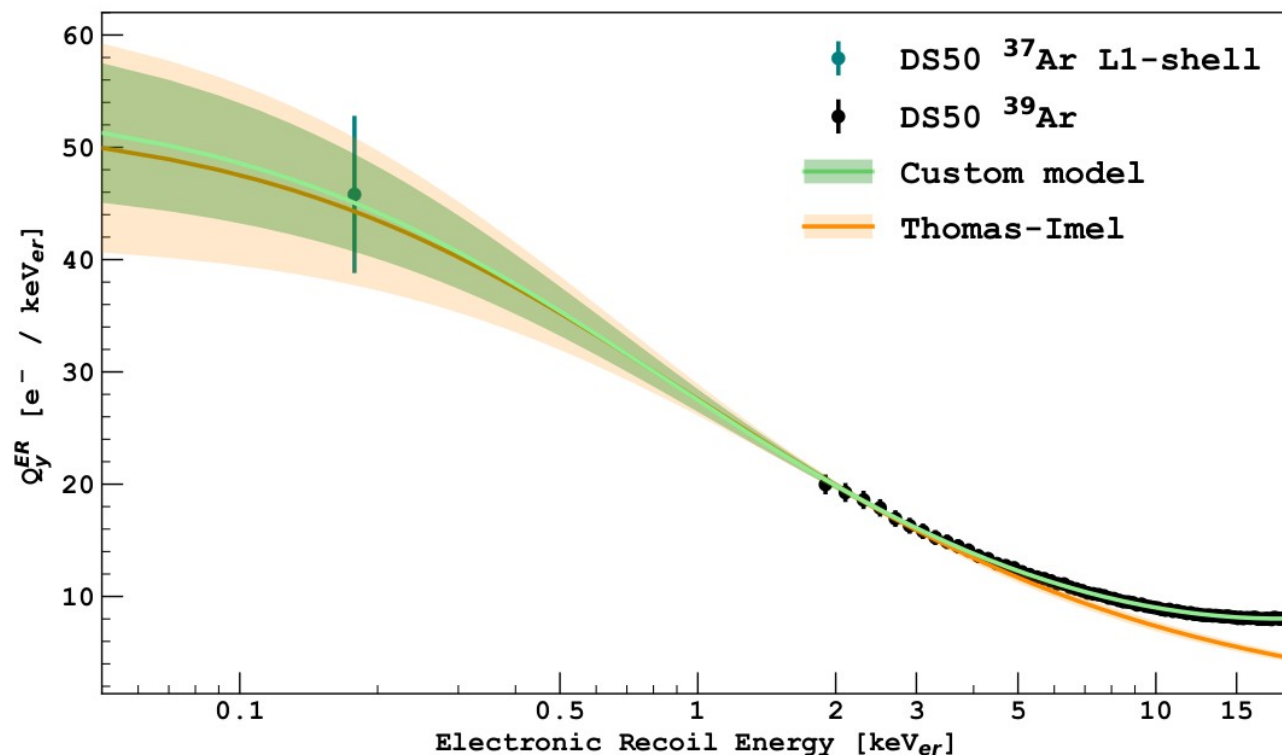
$$Q_y^{ER} = \frac{\ln(1 + \gamma \rho E_{er})}{E_{er}} \left( \frac{1}{\gamma} + p_0 (E_{er}/\text{keV}_{er})^{p_1} \right)$$

*Thomas-Imel model*

*Custom model:  
extension above 3keVer*

$\gamma = C_{box}/F$   
with F the drift field

→ Assumption of a constant excitation-to-ionization ratio



## Fit result

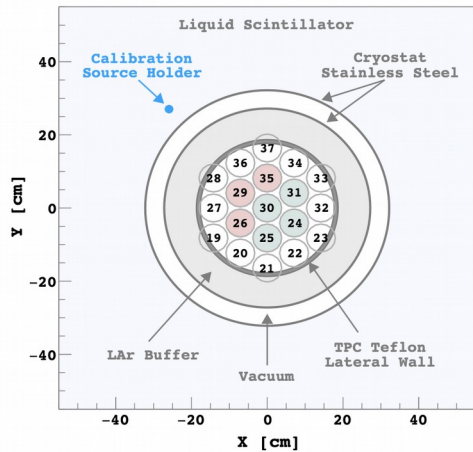
$C_{box} = 9.2 \pm 0.9 \text{ V/cm}$   
 $\rho = 54.4 \pm 7.3 \text{ keV}_{er}^{-1}$   
 $p_0 = 0.11 \pm 0.003$   
 $p_1 = 0.71 \pm 0.08$

Compatible with Thomas-Imel fit at  $1\sigma$   
up to 3keVer



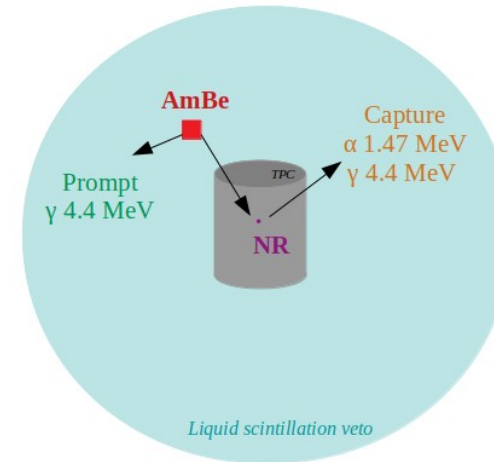
# Nuclear Recoil Ionization Yield

## $^{241}\text{Am}$ - $^{13}\text{C}$ (AmC)

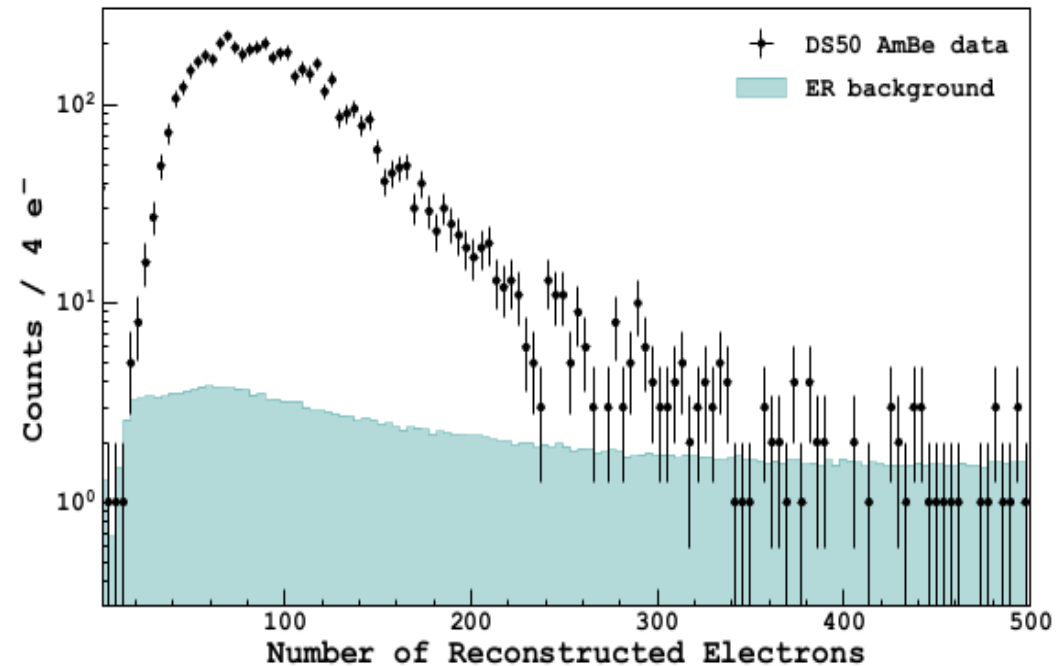
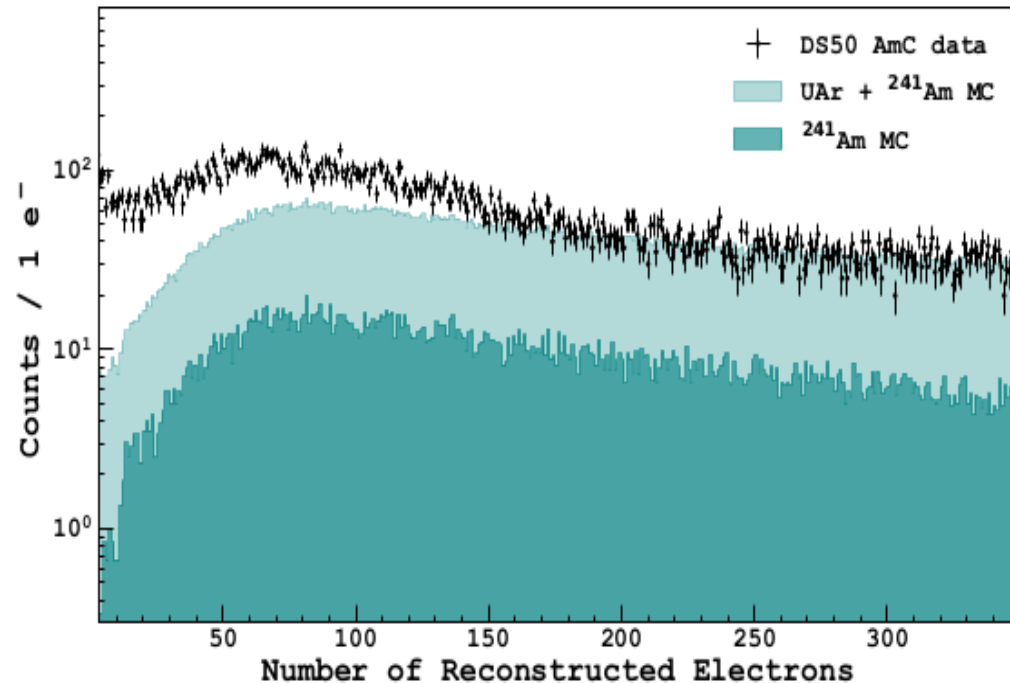


- Non-taggable source
- Stronger data selection to reduce Am  $\gamma$  contamination
- No coincidence with s1

## $^{241}\text{Am}$ - $^9\text{Be}$ (AmBe)



- Taggable source
- NR events selected by triple coincidence
- Coincidence with s1



# Simultaneous fits of Internal and External Data

## External datasets

### \* SCENE

- 4 ionization yields between 16.9 and 57.3keV,
- Drift field:  $g2=3.1\pm 0.3pe/e^-$
- Results normalized to DarkSide-50 response by the  $g2$  ratio

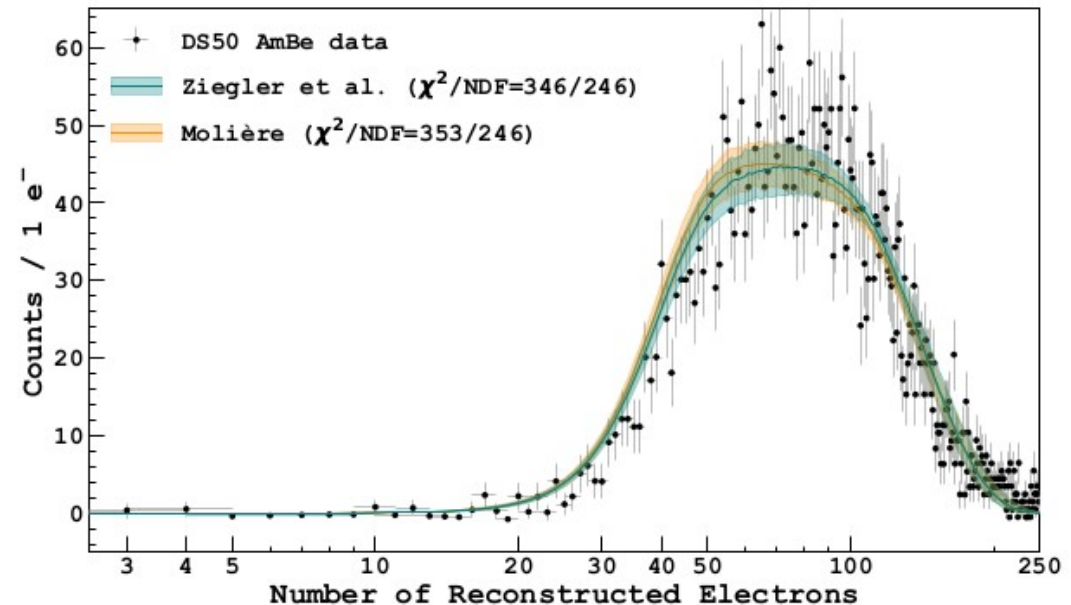
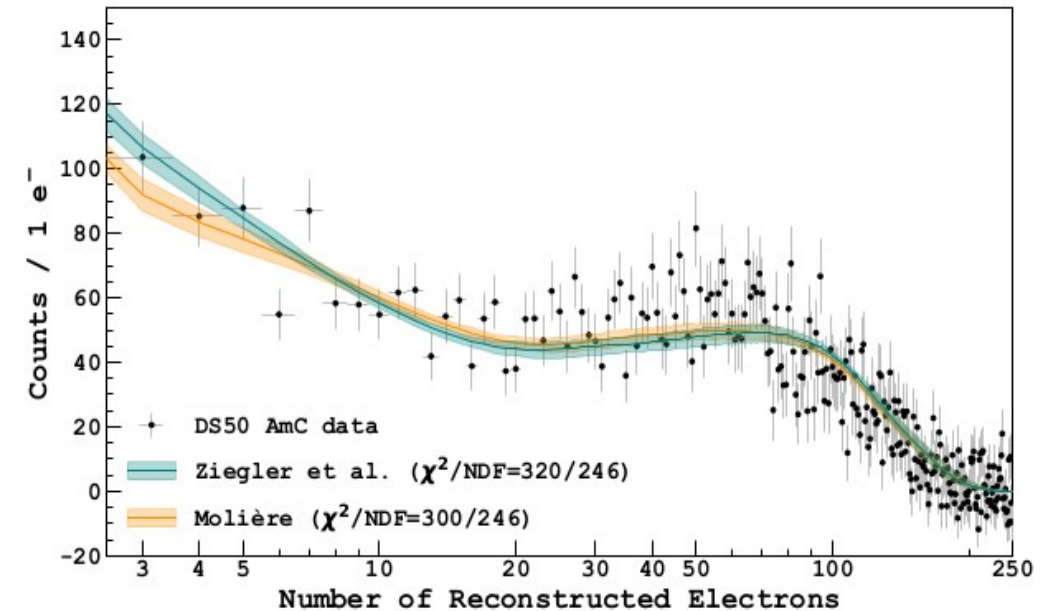
### \* ARIS

- 8 scintillation responses between 7.1 and 117.8 keVnr
- Same drift field than DarkSide-50
- Results normalized to DarkSide-50 by the ratio between field-off S1 yields
- S2 by the NR S2/S1 ratio within the AmBe dataset (MC simulations)

### \* Joshi et al.

- Ionization yield at 6.7keV
- Correction from the initial publication
- Compared to the final result only

## Internal datasets

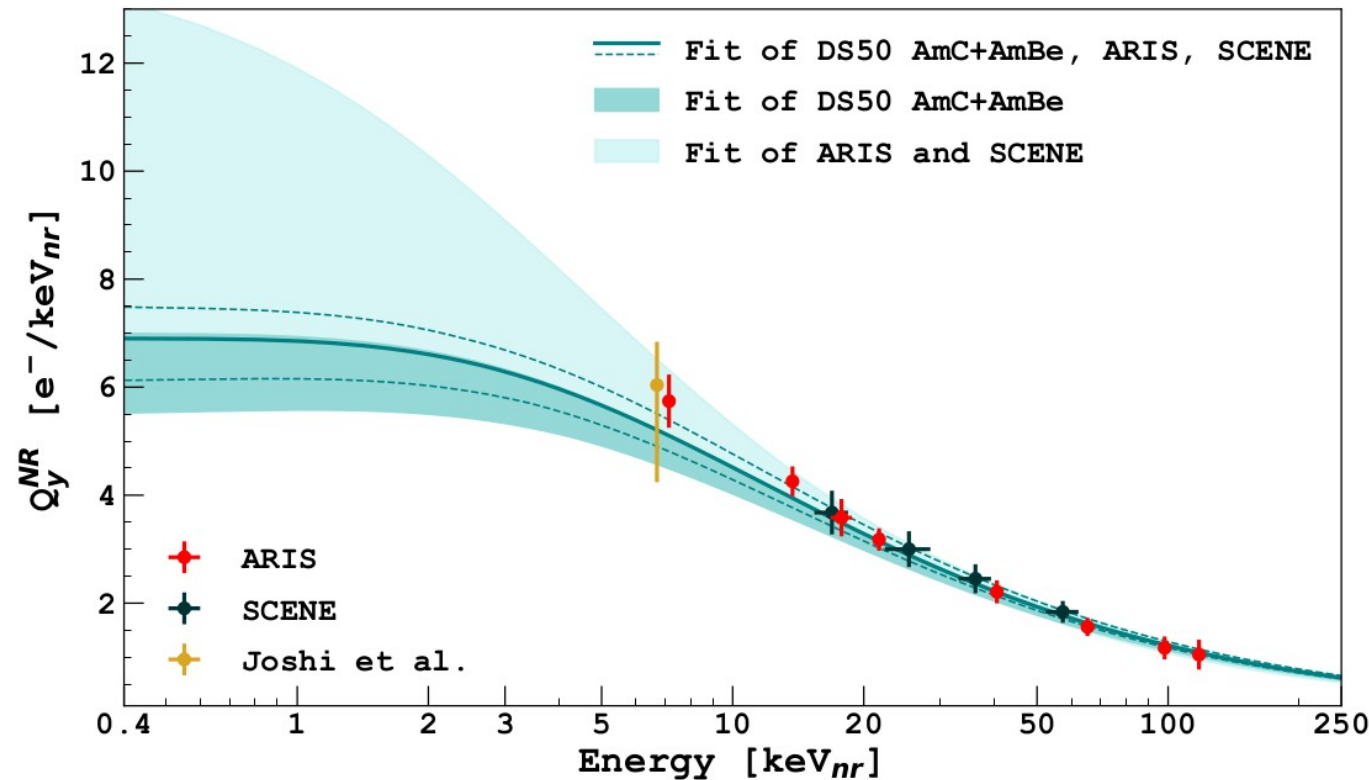


# Combined fit

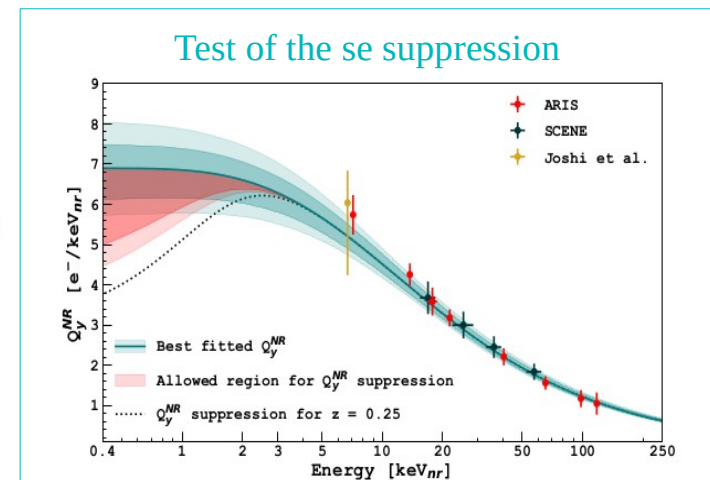
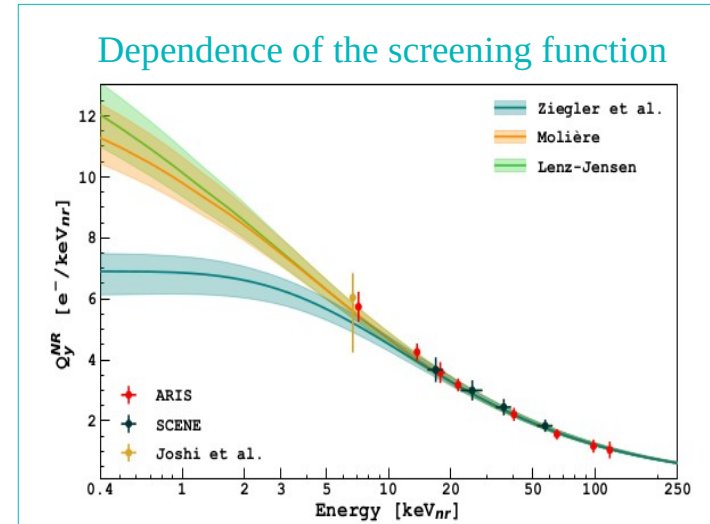
## Best parameters:

$C_{\text{box}} = 8.1^{+0.1}_{-0.2}$  V/cm

$B = 6.8^{+0.1}_{-0.3} \cdot 10^3$



- Lowest NR threshold in LAr: **3 electrons**
- Errors from statistical uncertainties and systematics of  $g_2$



See: [10.1103/PhysRevD.104.082005](https://arxiv.org/abs/10.1103/PhysRevD.104.082005)

*More details on:  
10.1103/PhysRevD.104.082005*

## Detector response validation

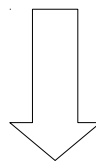
LAr ionization response calibration

**ER**

Down to  $180eV_{er}$   
Extrapolated down to  
*few tens of eV*

**NR**

Down to  $\sim 500eV_{nr}$   
(model dependant)



**New DarkSide-50 lowmass analysis in progress,  
stay tuned!**

Thank you for your attention!

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# 39Ar

Additional points from rotated energy as:

→ AAr campaign

$$E_{er} = w \left( \frac{S1}{g1} + \frac{S2}{g2} \right)$$

with  $w = 19.5 \pm 1.0$  eV

→ from [doi.org/10.1143/JJAP.41.1538](https://doi.org/10.1143/JJAP.41.1538)

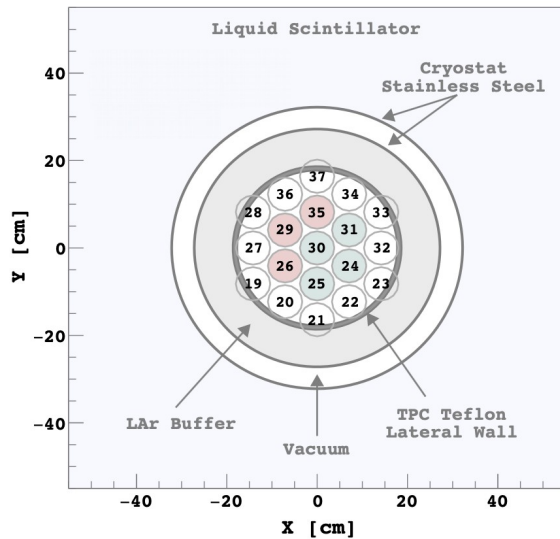
$g1 = 0.16 \pm 0.01$

→ fit of spectral shapes of  $^{133}\text{Ba}$  and  $^{57}\text{Co}$  [[arXiv:1707.05630v3](https://arxiv.org/abs/1707.05630v3)]

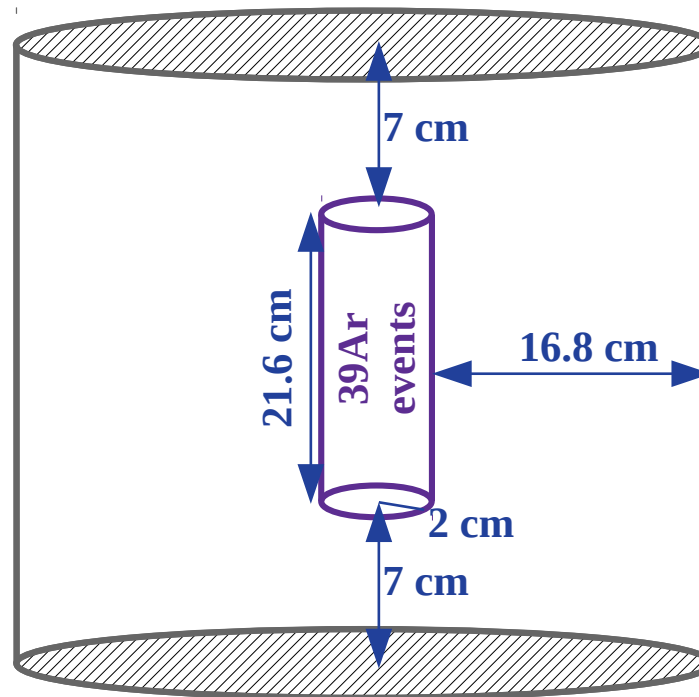
$g2 = 23 \pm 1$  pe/e- (in central PMTs)

→ from s2 echoes

## Event selection:



Fiducialisation of the volume



+ Number of S2  $\leq 1$

# Nuclear Recoil Ionization Yield

Ionization yield using Thomas-Imel

$$Q_y^{NR} = \frac{N_{i.e.}}{E_{nr}} = \frac{(1-r)N_i}{E_{nr}}$$

Recombination using Thomas-Imel

$$1-r = \frac{1}{\gamma N_i} \ln(1 + \gamma N_i) \quad \text{with } \gamma = \boxed{\text{Cbox}}/F$$

Numbers of ion as a function of electronic and nuclear stopping power

$$N_i = \beta \kappa(\epsilon) = \boxed{\beta} \frac{\epsilon s_e(\epsilon)}{s_n(\epsilon) + s_e(\epsilon)} \quad \text{Free parameters of the model}$$

Reduced energy

$$\epsilon = \frac{a}{2e^2 Z^2} E_{nr}/\text{keV} \simeq 0.0135 E_{nr}/\text{keV}$$

Electronic stopping power

$$s_e(\epsilon) = \frac{0.133 Z^{2/3}}{A^{1/2}} \sqrt{\epsilon} \simeq 0.145 \sqrt{\epsilon}$$

Nuclear stopping power based on Universal Screening Function

*Other functions tested*

$$s_n(\epsilon) = \frac{\ln(1 + 1.1383 f_Z \epsilon)}{2[f_Z \epsilon + 0.01321(f_Z \epsilon)^{0.21226} + 0.19593(f_Z \epsilon)^{0.5}]}$$

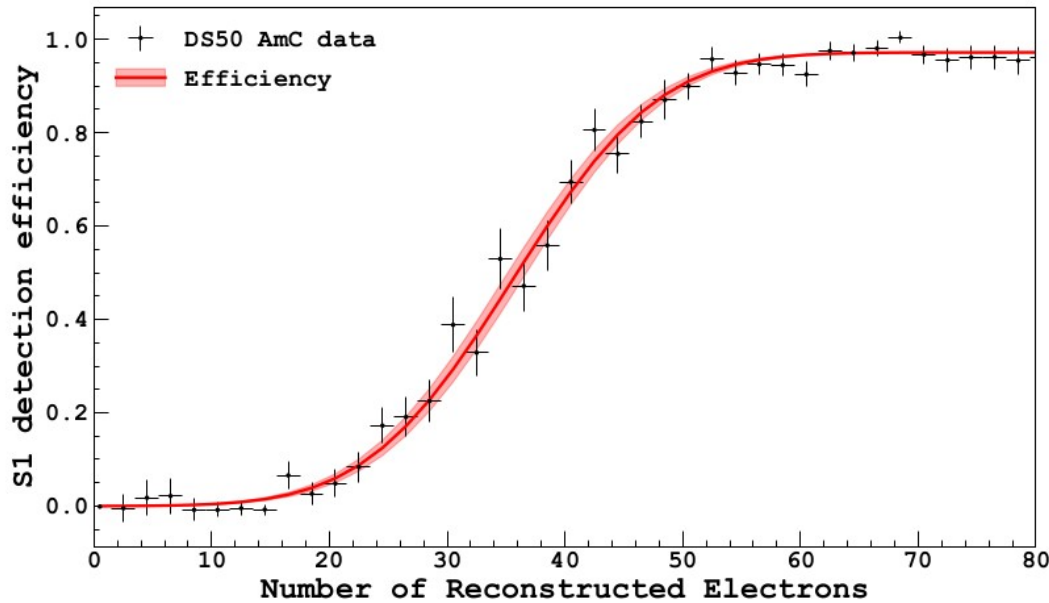
## Test of two extreme models for intrinsic resolution

- 1/ Binomial fluctuations in energy quenching
- 2/ No fluctuations in energy quenching

—————▶ Negligible difference: use of 2/

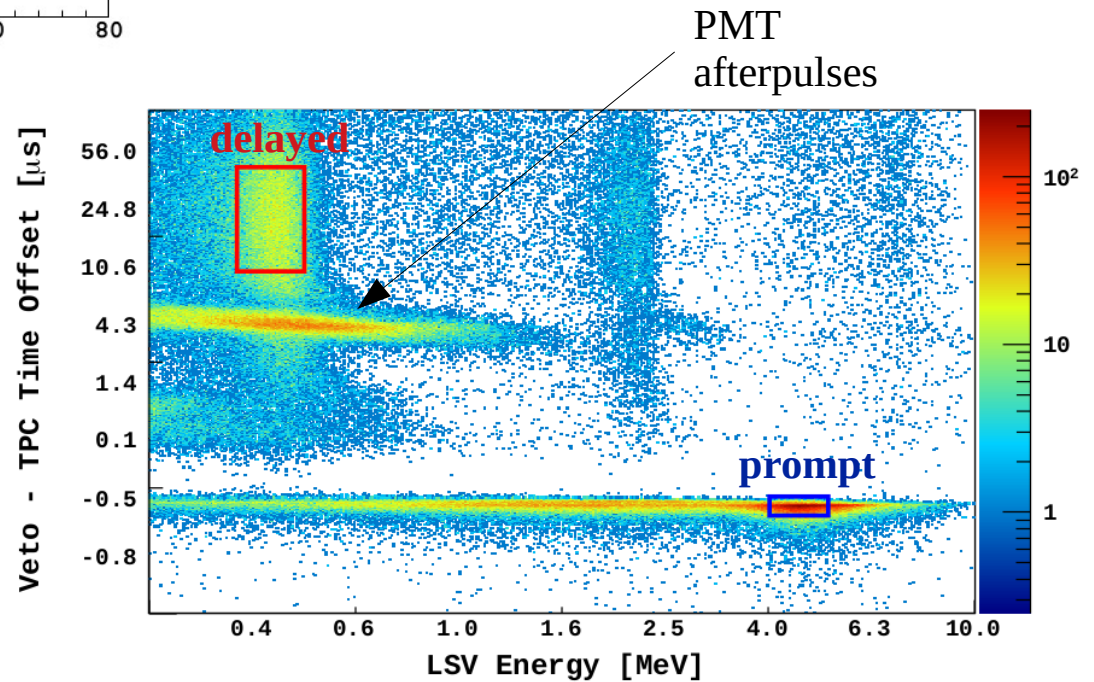


# Nuclear Recoil Ionization Yield: AmBe



## S1 detection efficiency:

- Ratio of ER contamination subtracted to AmC samples for S2only and for S1+S2
- Error function fit

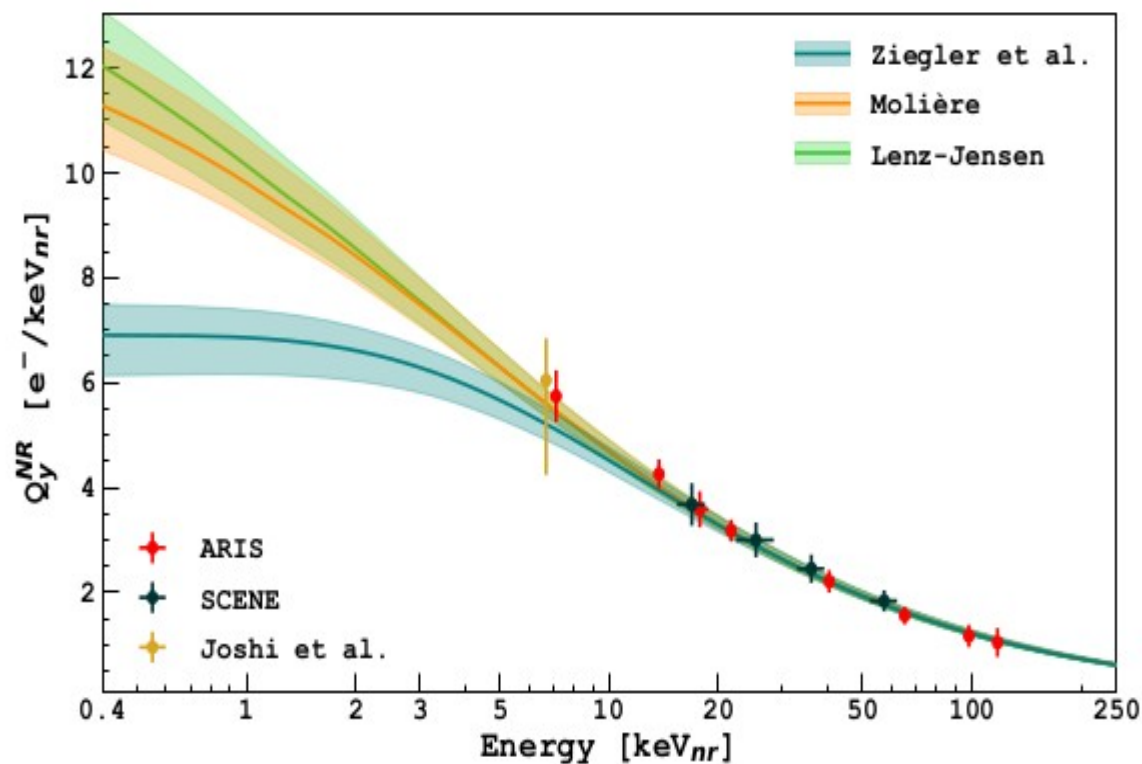


(Detectors difference trigger offset: -550 ns)

# Dependence on the screening function

## Different nuclear stopping power models

$$s_n(\epsilon) = \frac{1}{\epsilon} \int_0^\epsilon f(\eta) d\eta \quad \text{with} \quad f(\eta) = \frac{\lambda \eta^{1-2m}}{(1 + [2\lambda \eta^{2(1-m)}]^q)^{1/q}}$$



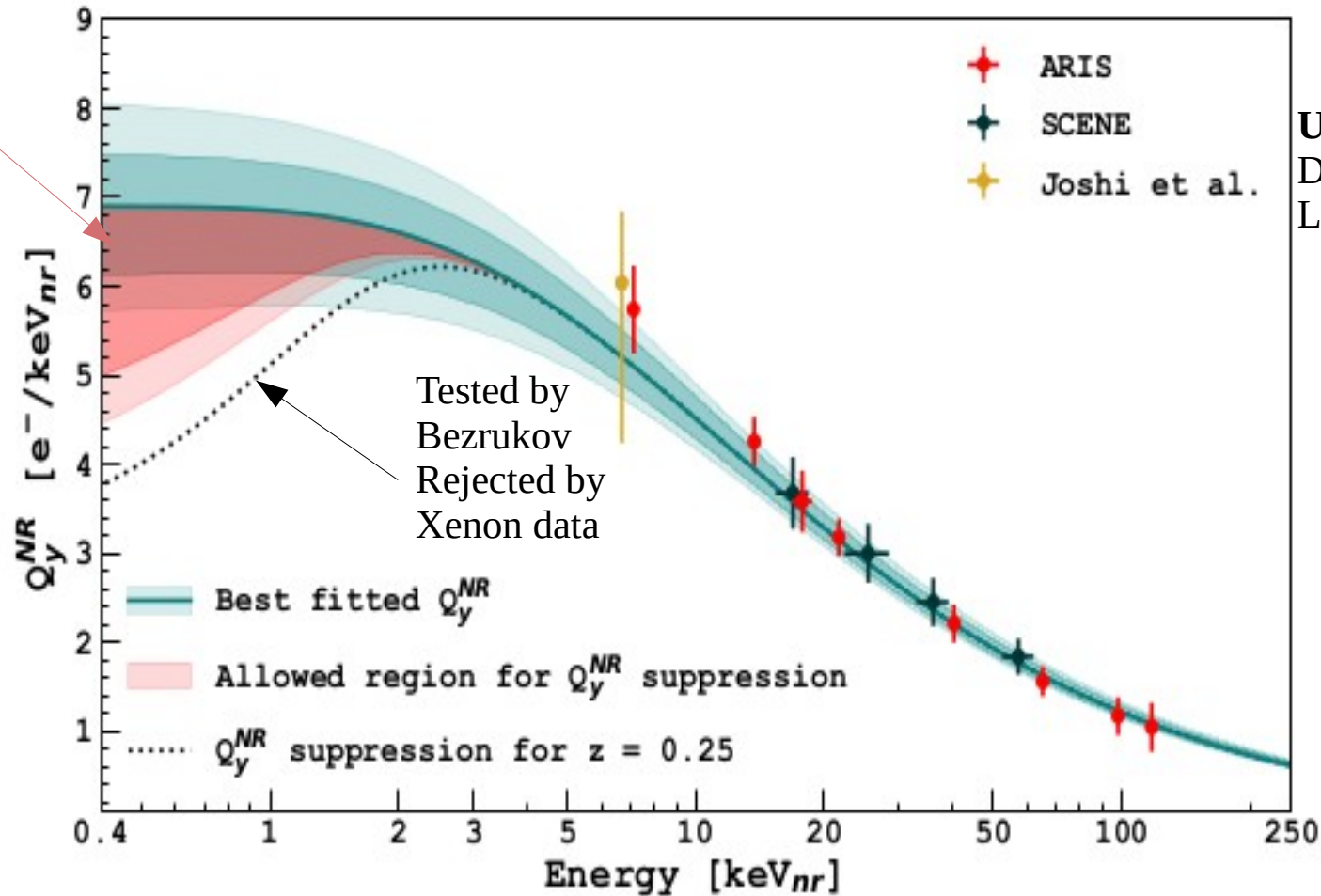
Ziegler et al. yields the **lowest ionization yield** in the WIMP's interest region  
→ most conservative result

# Se suppression

Impact of low energy se suppression (from arXiv:1011.3990)

$$F(v/v_0) = 1/2 (1 + \tanh(50 \epsilon - z)) \quad \text{with} \quad F(v/v_0) \rightarrow 1 \text{ for } z \rightarrow -\infty$$

and  $z = 0.25$ : hypothesis of Coulomb effects inside se



→ Suppression **not compatible** with the LXe (arXiv:1011.3990) and AmC dataset ( $z > 0.04$  excluded at  $2\sigma$ )