

Exactly Stable Protons with a Muonic Force

Anders Eller Thomsen

Based on work with J. Davighi, A. Greljo, and P. Stangl

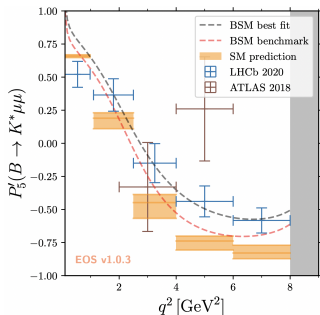
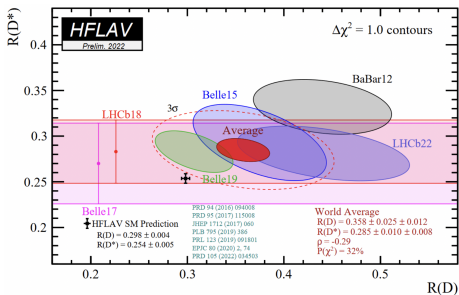
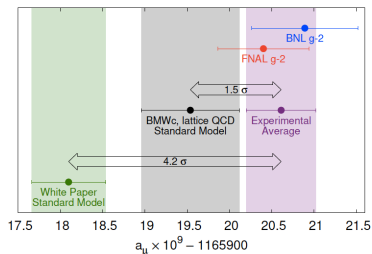
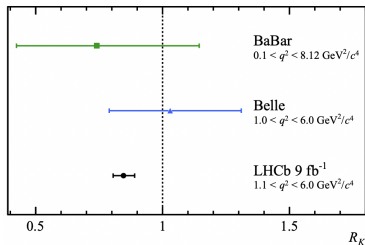
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FOR FUNDAMENTAL PHYSICS

DISCRETE 2022, 7–11 November

TeV-scale new physics?



Gubernari et al. [2206.03797]

Selection rules in the SM

The Standard Model sans Higgs

Symmetries:

$SU(3)_c \times SU(2)_L \times U(1)_Y \times \text{Poincaré}$

Matter fields:

$q_i, u_i, d_i, \ell_i, e_i$

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Explicit breaking by the Higgs couplings

$$\mathcal{L}_{\text{yuk}} = -\bar{q}y_u\tilde{H}u - \bar{q}y_dHd - \bar{\ell}y_eHe + \text{H.c.}$$

$$\hookrightarrow G_F \longrightarrow U(1)_B \times U(1)_{L_e} \times U(1)_{L_\mu} \times U(1)_{L_\tau}$$

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Consequences:

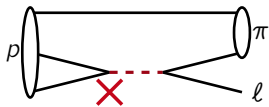
- Conservation of B number
- Conservation of LF

Exceptions:

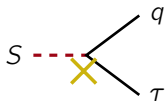
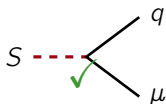
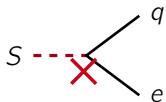
- Non-perturbative violation of $B + L$
- Neutrino oscillations

Selection rules for leptoquarks and muonic forces

Diquark interactions



LQ interactions

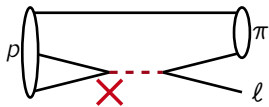


Problem: TeV-scale scalar LQ interactions are strongly constrained, e.g.,

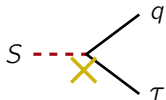
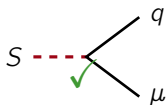
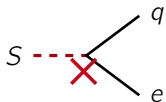
- Proton lifetime $\tau_p \gtrsim 10^{34}$ yr
- LFV in muon decays: $\text{BR}(\mu \rightarrow e\gamma) < 10^{-12}$

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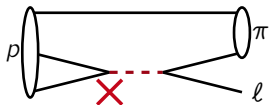
Solution: *Lepton-flavored gauged* $U(1)_X$

Hambye, Heeck [1712.04871]; Davighi, Kirk, Nardecchia [2007.15016]; Greljo, Stangl, AET [2103.13991]; Greljo, Soreq, Stangl, AET, Zupan [2107.07518]; Davighi, Greljo, AET [2202.05275]; Heeck, Thapa [2202.08854]; Crivellin *et al.* [2203.10111]; Greljo, Stangl, AET, Zupan [2203.13731]

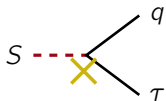
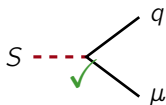
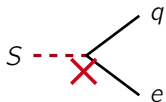


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"Muquark"

The accidental symmetries of \mathcal{L}_{LQ} are those of the SM, $U(1)_B \times U(1)_{L_e} \times U(1)_{L_\mu} \times U(1)_{L_\tau}$, and the leptoquarks have charge $S \sim (-1/3, 0, -1, 0)$

The $U(1)_X$ model

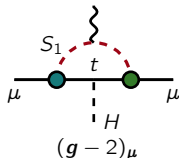
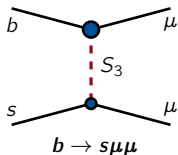
- Initial $U(1)_X$ model: $X = B - 3L_\mu \sim (m = 1, n = 3)$
- The generalization: $X = 3m(B - L) - n(2L_\mu - L_e - L_\tau)$, $\gcd(m, n) = 1$

	Fields	$U(1)_X$
Quarks	q_i, u_i, d_i	m
Electrons & taus	$\ell_{1,3}, e_{1,3}, \nu_{1,3}$	$n - 3m$
Muons	ℓ_2, e_2, ν_2	$-2n - 3m$
Higgs	H	0
Leptoquarks	S_3, S_1	$2m + 2n$
SM singlets	$\phi_{e\tau}$ ϕ_μ	$6m - 2n$ $6m + n$

Muonquark fit

Muonquark mediated anomalies

Crivellin, Müller, Ota [1703.09226]; Gherardi, Marzocca, Venturini [2008.09548]

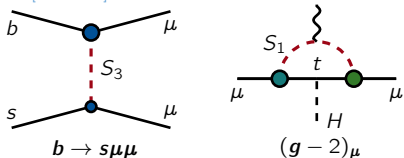


- Direct searches give only modest constraints: $M_{1,3} \gtrsim 1.7 \text{ TeV}$
ATLAS collaboration [2006.05872] [2210.04517]
- Decoupling limit ($\begin{smallmatrix} v_X \rightarrow \infty \\ g_X \rightarrow 0 \end{smallmatrix}$) ensures NP contribution exclusively from $S_{1,3}$
- Approximate U(2) flavor symmetry
Kagan et al. [0903.1794]; Barbieri et al. [1105.2296]
- *Global fit* with `smelli` (also using `wilson` and `flavio`)

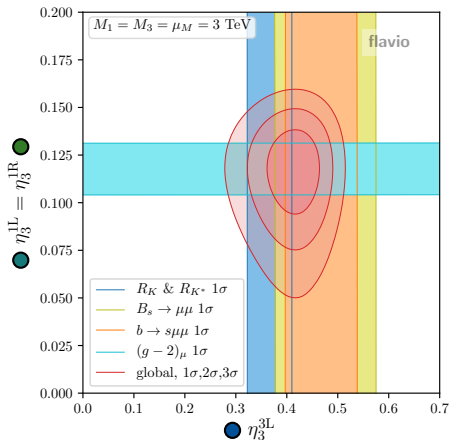
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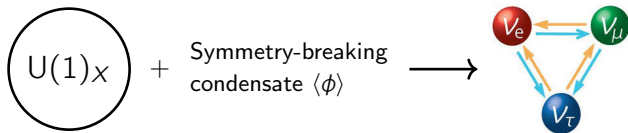


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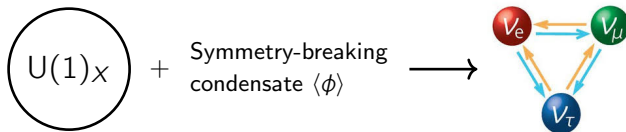


Greljo, Stangl, AET [2103.13991]

Neutrino masses



Neutrino masses

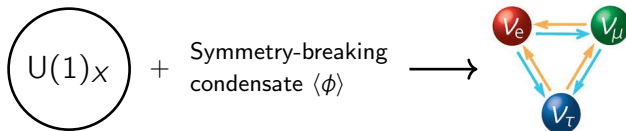


$$\mathcal{L} \supset -y_e^{ij} \bar{\ell}_i e_j H - y_\nu^{ij} \bar{\ell}_i \nu_j H^*$$

$$y_{e,\nu} \sim \begin{pmatrix} \times & 0 & \times \\ 0 & \times & 0 \\ \times & 0 & \times \end{pmatrix}$$

Fields	$U(1)_X$
q_i, u_i, d_i	m
$\ell_{1,3}, e_{1,3}, \nu_{1,3}$	$n - 3m$
ℓ_2, e_2, ν_2	$-2n - 3m$
H	0
S_3, S_1	$2m + 2n$
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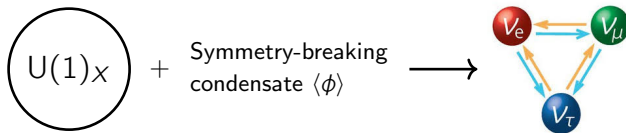
$$\mathcal{L} \supset -y_e^{ij} \bar{\ell}_i e_j H - y_\nu^{ij} \bar{\ell}_i \nu_j H^* - \bar{\nu}^i c \nu^j (\xi_{e\tau}^{ij} \phi_{e\tau} + \xi_\mu^{ij} \phi_\mu)$$

$$y_{e,\nu} \sim \begin{pmatrix} \times & 0 & \times \\ 0 & \times & 0 \\ \times & 0 & \times \end{pmatrix} \quad \frac{M_\nu}{\langle\phi\rangle} \sim \begin{pmatrix} \times & \times & \times \\ \times & 0 & \times \\ \times & \times & \times \end{pmatrix}$$

\Rightarrow Seesaw neutrino masses: $\langle\phi\rangle \lesssim 10^{11}$ TeV

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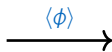
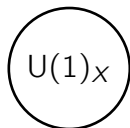
⇒ Seesaw neutrino masses: $\langle\phi\rangle \lesssim 10^{11}$ TeV

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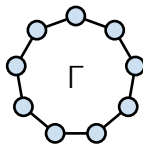
What happens to the accidental symmetries in the U(1)_X broken phase? Proton decay? cLFV?

Discrete IR remnant

Gauged UV symmetry



Exact IR remnant



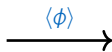
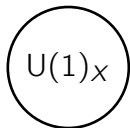
$$U(1)_X \longrightarrow \Gamma = \left\{ e^{i\alpha} \in U(1)_X : e^{i\alpha[\phi]_X} \phi = \phi \right\} \cong \mathbb{Z}_k,$$

Fixed by neutrino charges

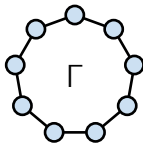
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Diquark operators are banned by Γ *iff*

$$(m, n) = (3a + r, 9b + 3r), \quad \text{for } r \in \{1, 2\},$$

$$(a, b) \in \mathbb{Z}^2, \quad \text{and } \text{gcd}(3a + r, b - a) = 1$$

Davighi, Greljo, AET [2202.05275]

$$\Gamma \cong \begin{cases} \mathbb{Z}_9, & \text{for } b + r \in 2\mathbb{Z} + 1 \\ \mathbb{Z}_{18}, & \text{for } b + r \in 2\mathbb{Z} \end{cases}$$

- No proton decay mediated by LQs
- $m, n \neq 0 \implies$ lepton-flavored $U(1)_X$ symmetry

IR selection rules

$b + r \pmod{2}$	Γ	ℓ	q	S	$qS\ell$	qS^*q
0	\mathbb{Z}_{18}	9	$3a + r$	$6a + 8r$	0	$12r$
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- Low-energy EFT contains only q and $\ell \implies \Delta B = 0 \pmod{3}$!
 - Exact proton stability, no matter other NP (all orders in EFT)
 - Neutron–antineutron oscillations forbidden
 - EW sphalerons are allowed (leptogenesis is still possible)

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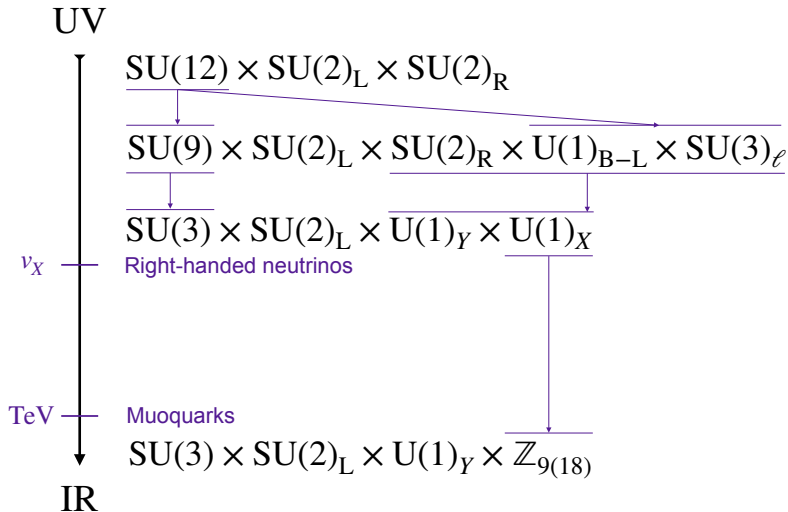
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Neutrino masses are compatible with high-quality IR selection rules
from a muonic forces

Unification of lepton-flavored $U(1)_X$



Davighi, Greljo, AET [2202.05275]

- Muonic (lepton-flavored) forces provide a robust organizing principle for NP explanations of flavor anomalies
- Neutrino masses are compatible with high-quality IR selection rules from a muonic forces
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Thank You!