

Simulation of quark/gluon jets

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Outline

1. Motivation and definitions
2. LH Quark/gluon jets discrimination [[Les Houches arXiv:1605.04692](#)]
3. What we have learnt [[JHEP 1707 \(2017\) 091](#)]
4. How we improved the q/g jets simulation in Herwig [[arXiv:1708.01491](#)]
5. Outlook

Motivation

BSM searches: often signature for a BSM signals: many quark, backgrounds: QCD gluons

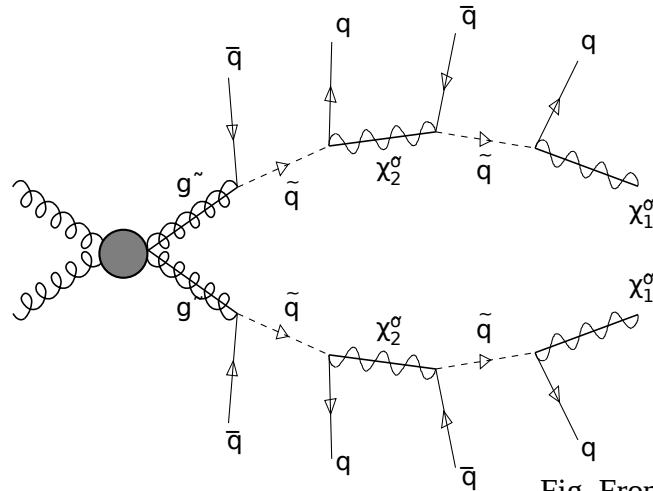
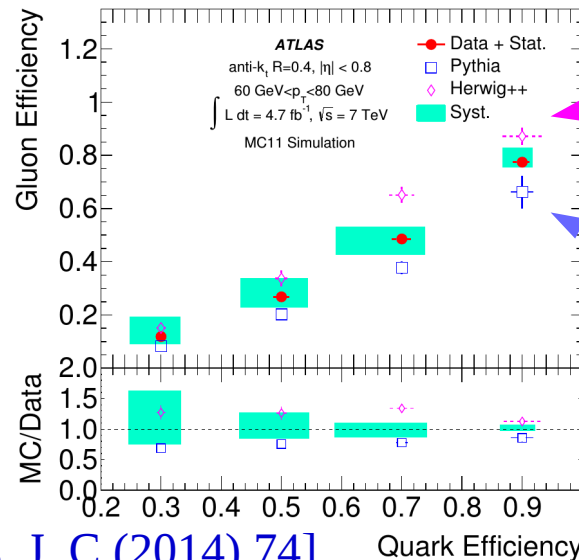


Fig. From J. Gallicchio and M. D. Schwartz, Phys. Rev. Lett.107 (2011)

Problem: Q/G jets LHC data show discrepancy with the predictions from MC generators



Herwig++ too pessimistic,
 Quark and gluon jets looks
 more the same than in data.

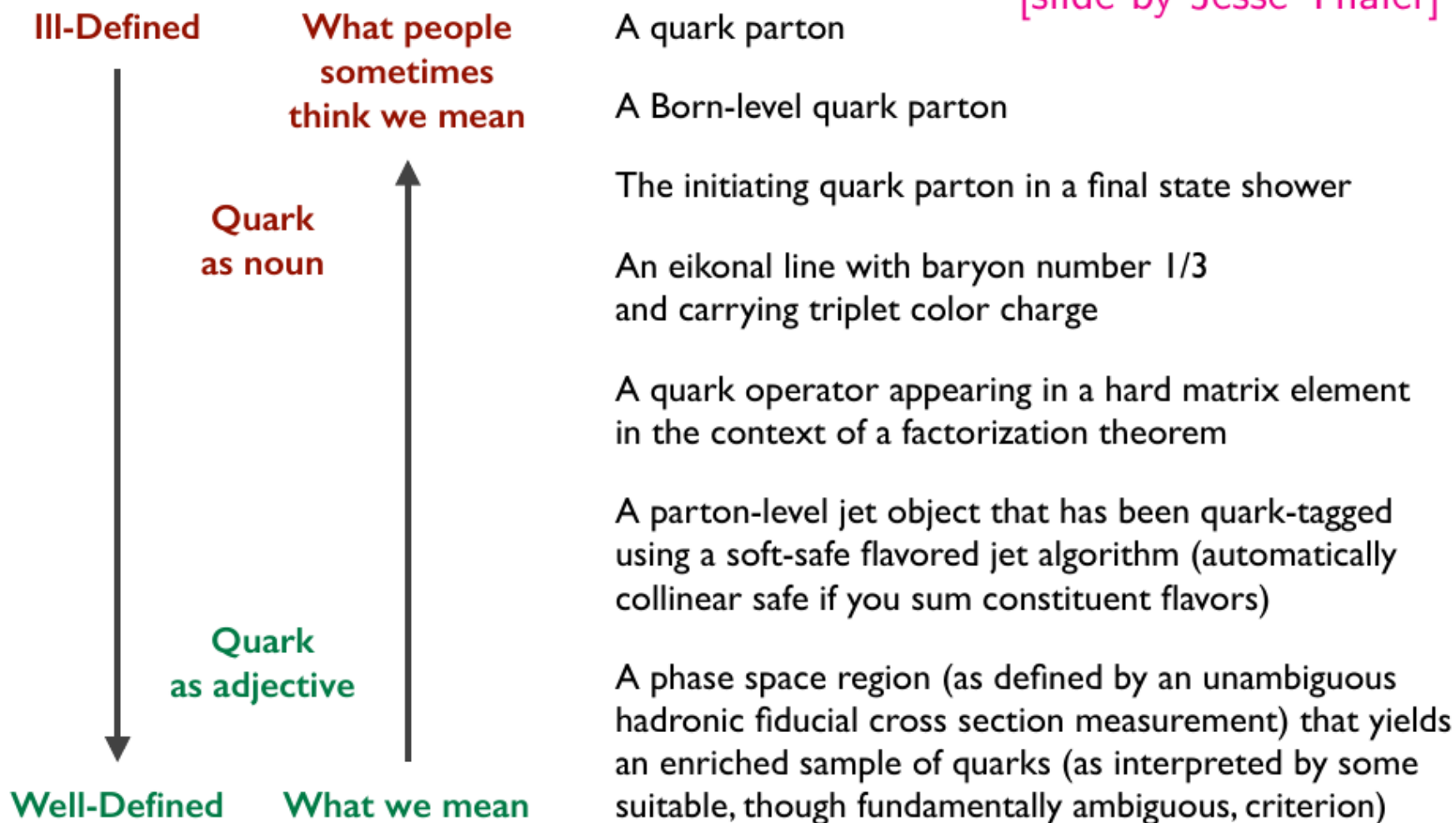
Pythia too optimistic,
 Quark and Gluon jets are
 too similar compared to
 data.

[ATLAS, Eur. Phys. J. C (2014) 74]

What is a Quark Jet?

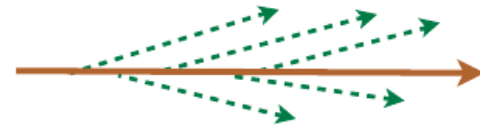
From lunch/dinner discussions

[slide by Jesse Thaler]



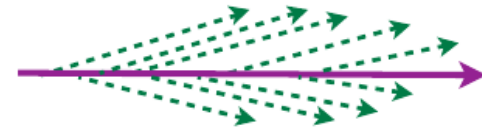
Definition

Cartoon:



Quark: $C_F = 4/3$

vs.



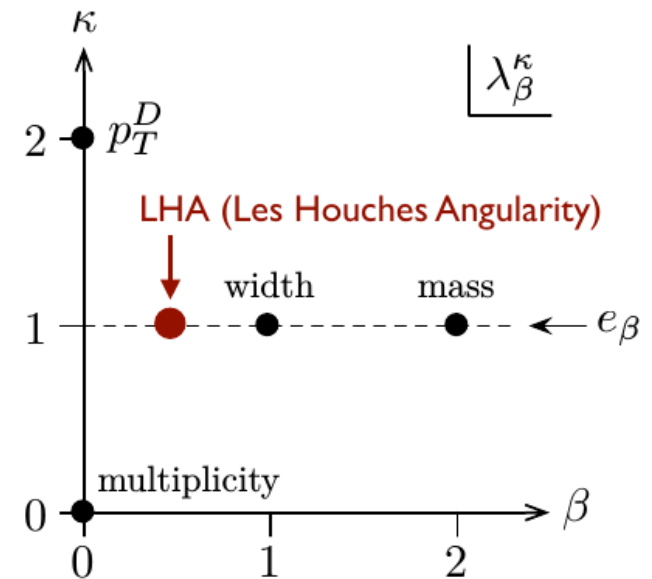
Gluon: $C_A = 3$

Probe radiation pattern with
e.g. Generalized Angularities

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta}$$

↑ momentum fraction ↑ angle to recoil-free axis

$(\lambda_{\beta}^{\kappa})_{\text{quark}} < (\lambda_{\beta}^{\kappa})_{\text{gluon}}$



[Larkoski, Salam, Thaler, 13]
[Larkoski, Thaler, Waalewijn, 14]

Framework

Processes:

- Quark: $e^+e^- \rightarrow (\gamma/Z)^* \rightarrow u\bar{u}$
- Gluons: $e^+e^- \rightarrow H^* \rightarrow gg$

Different settings:

- Changing the collision energy Q
- Changing the jet radius R

Different Monte-Carlo generators at parton and hadron level:

- Pythia 8 (v8.205)
- Herwig++ (v2.7.1)
- Sherpa (v2.1.1)

Additionally different Parton Shower algorithms

- Vincia (v1.201 - plugin to Pythia)
- Deductor (v1.0.2 + hadronization from Pythia)
- Ariadne (v5.0.β + hadronization from Pythia)

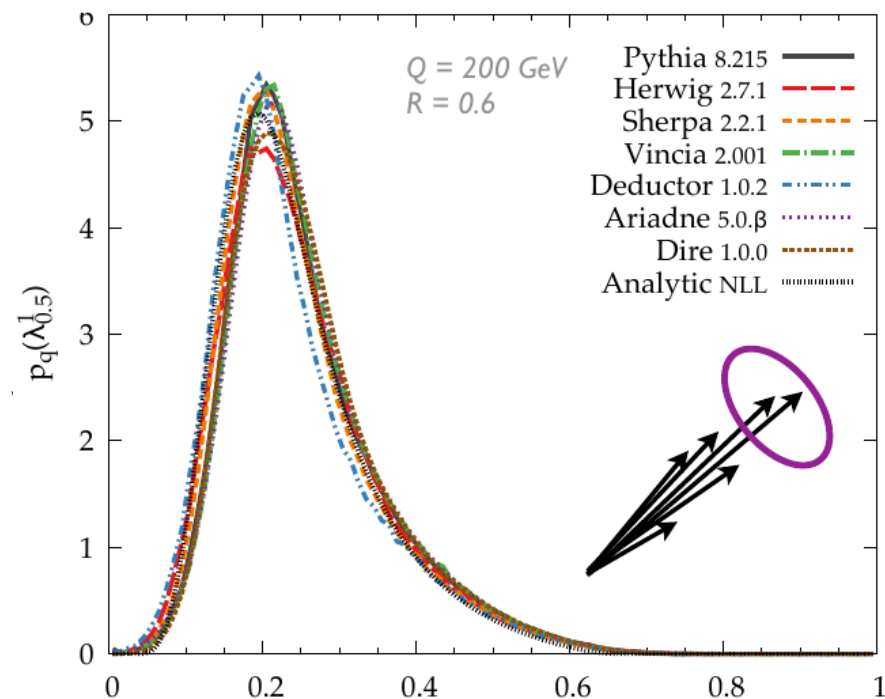
LHA – Idealized Quark/Gluon distributions

[Gras, Hoeche, Kar, Larkoski, Lönnblad, Plätzer, AS, Skands, Soyez, Thaler, JHEP 1707 (2017) 091]

$e^+e^- \rightarrow \text{quarks } (C_F = 4/3)$

VS.

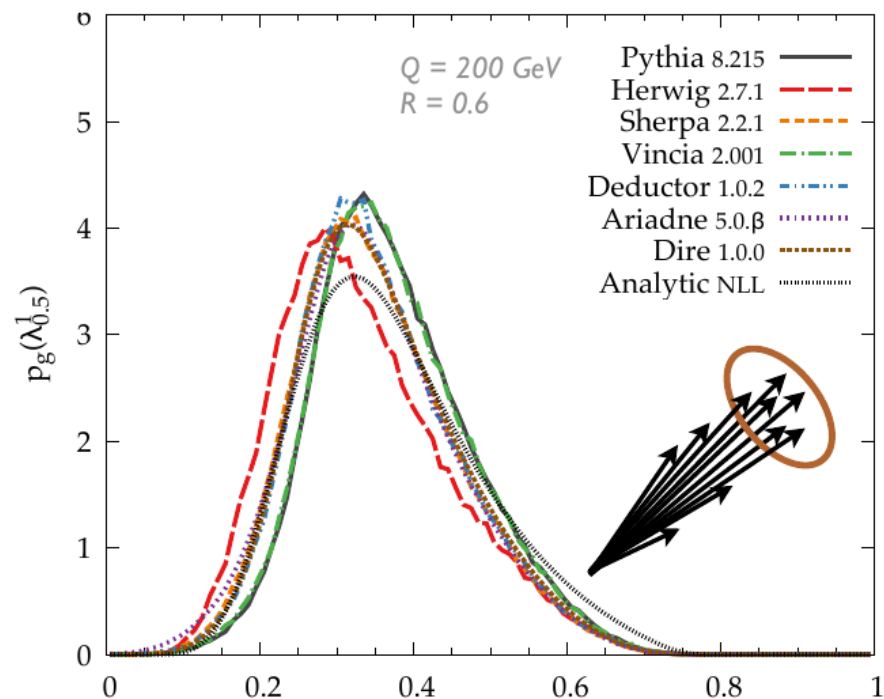
$e^+e^- \rightarrow \text{gluons } (C_A = 3)$



$$\text{LHA} = \sum_i z_i \sqrt{\theta_i}$$

Small spread

Constrained by LEP



$$\text{LHA} = \sum_i z_i \sqrt{\theta_i}$$

Large spread

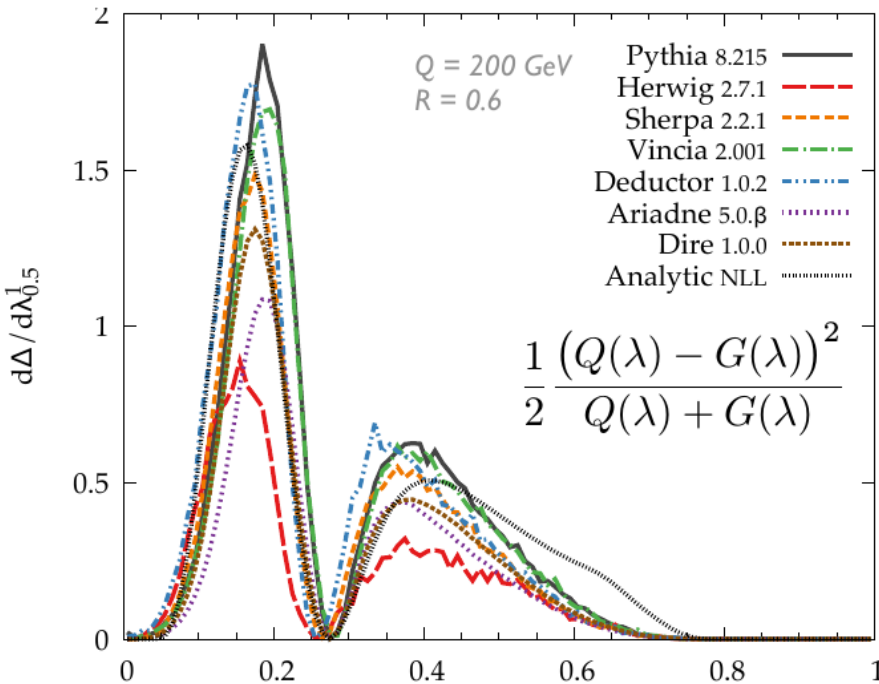
Up to now no e^+e^- data has been used to constrain it.

LHA – Separation power

$$\Delta = \frac{1}{2} \int d\lambda \frac{(p_q(\lambda) - p_g(\lambda))^2}{p_q(\lambda) + p_g(\lambda)}$$

$\Delta = 0$ - corresponds to no discrimination power.
 $\Delta = 1$ - corresponds to perfect discrimination power.

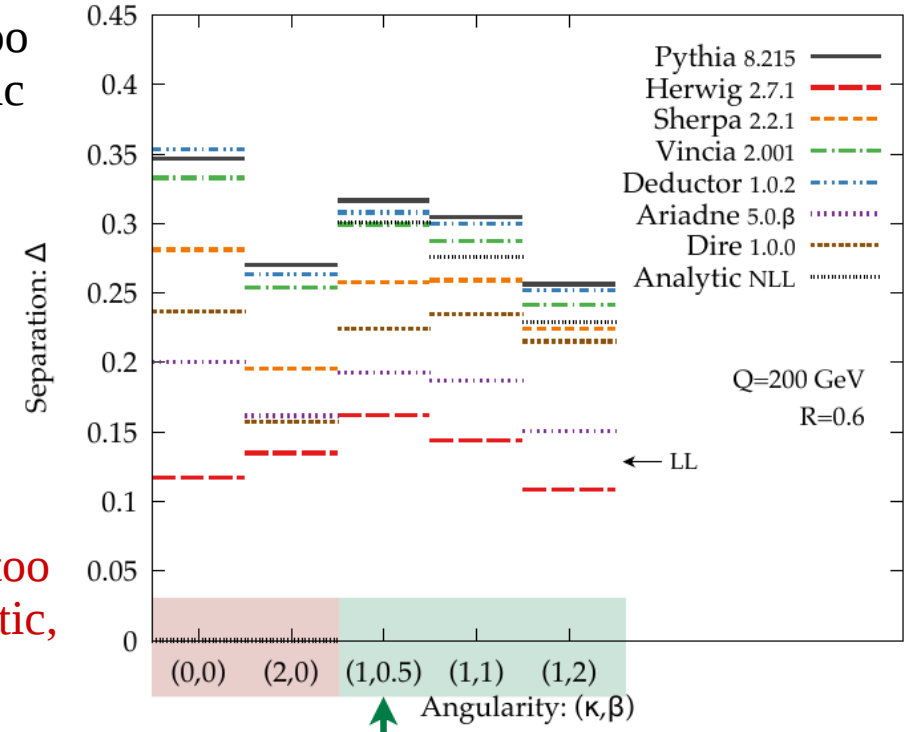
Differential



Pythia too optimistic

Herwig too pessimistic,

Integrated Values



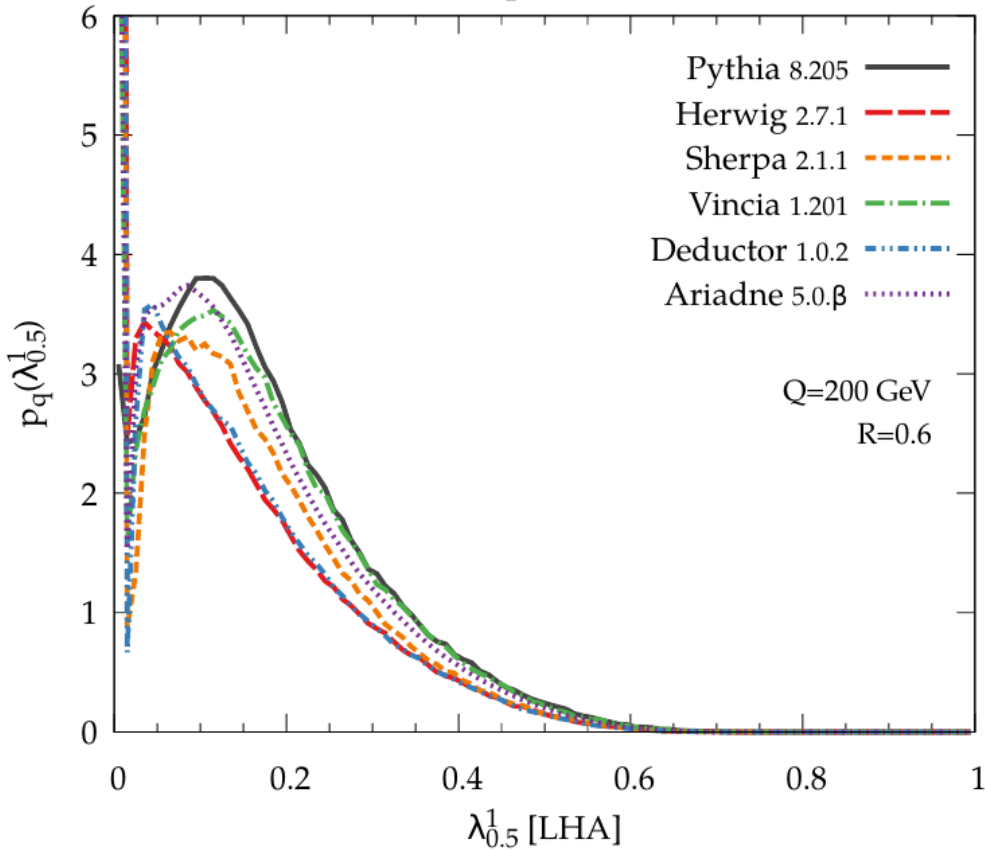
$$\text{LHA} = \sum_i z_i \sqrt{\theta_i}$$

Affects both **IRC unsafe** and **IRC safe** observables

Separation power – non-perturbative effects

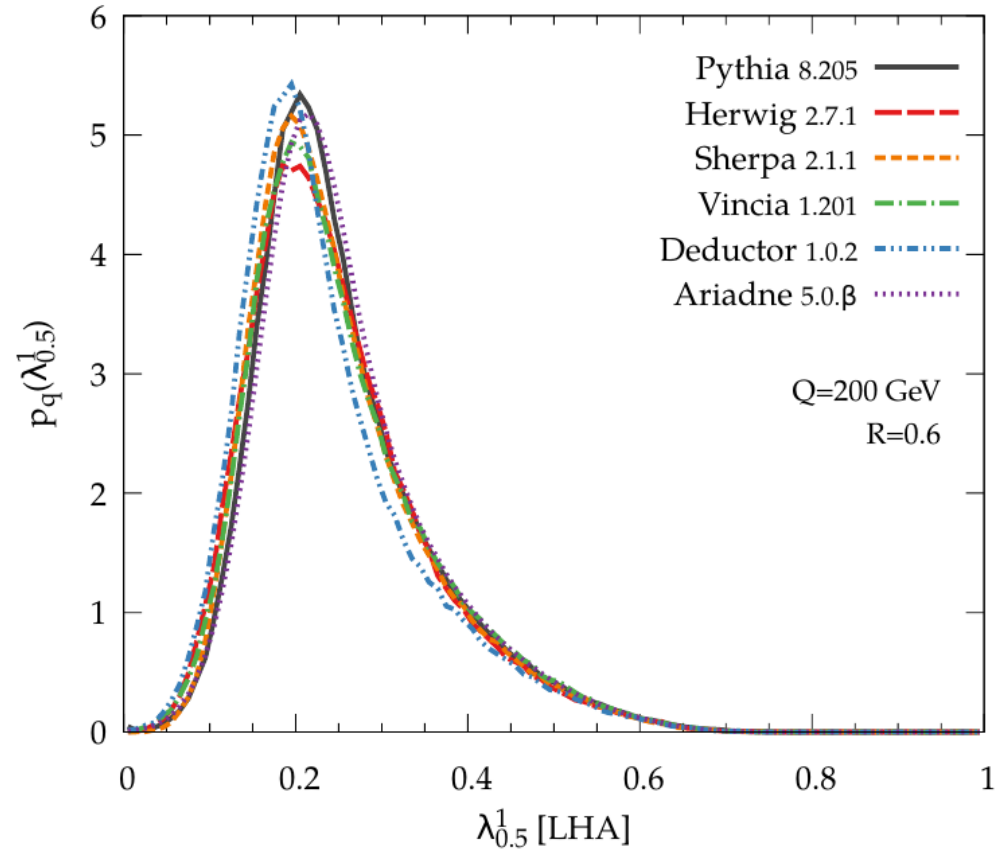
Parton level

Quark, parton-level



Hadron level

Quark, hadron-level

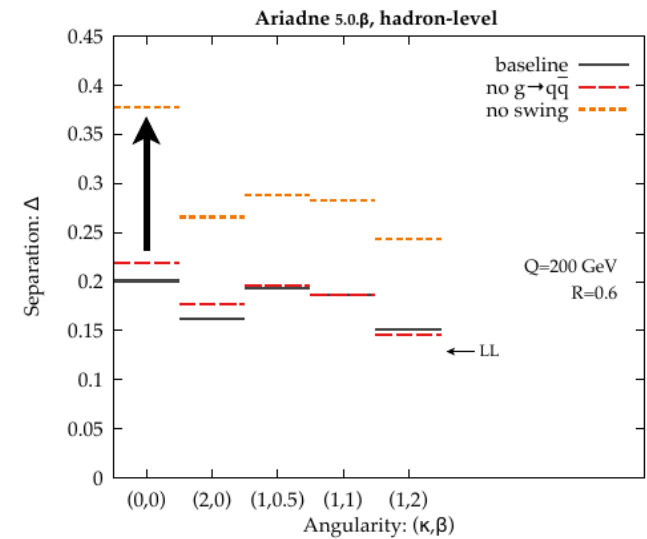
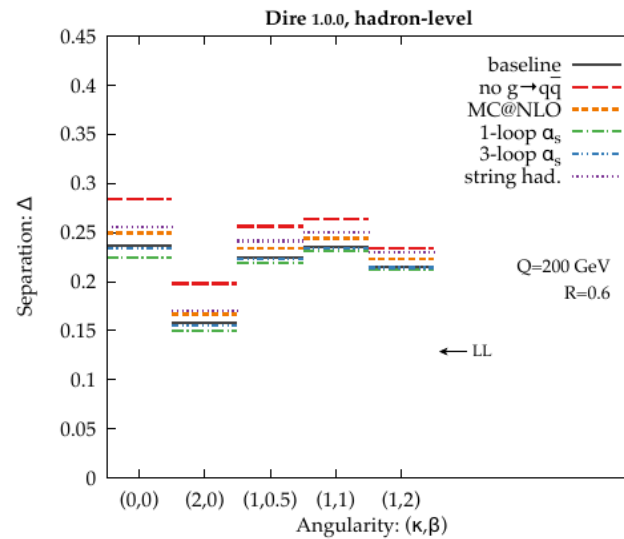
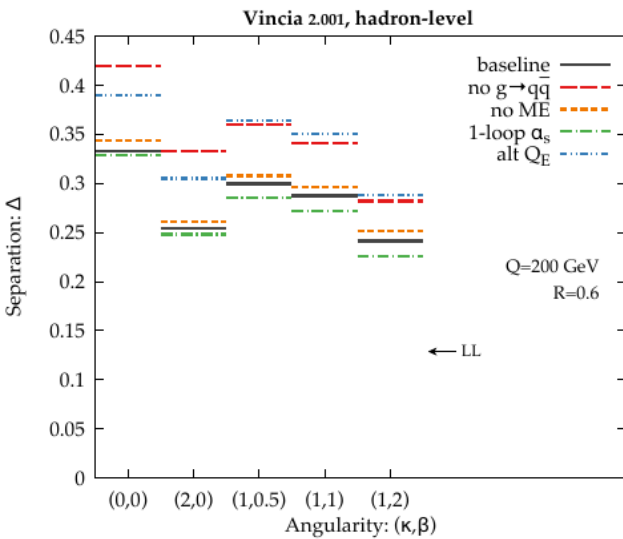
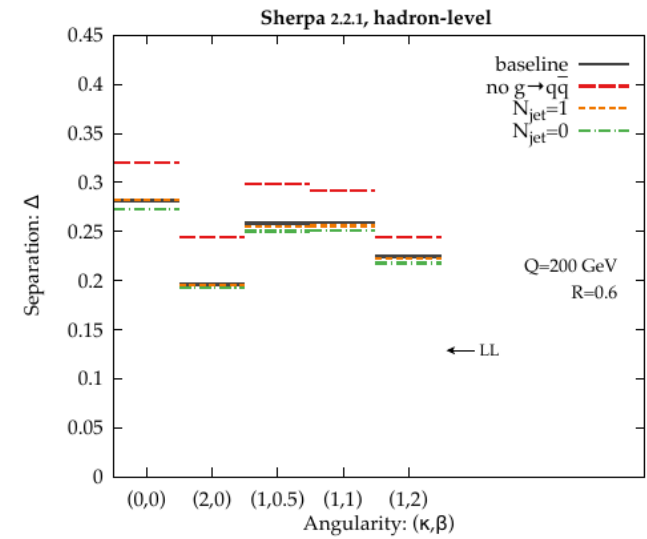
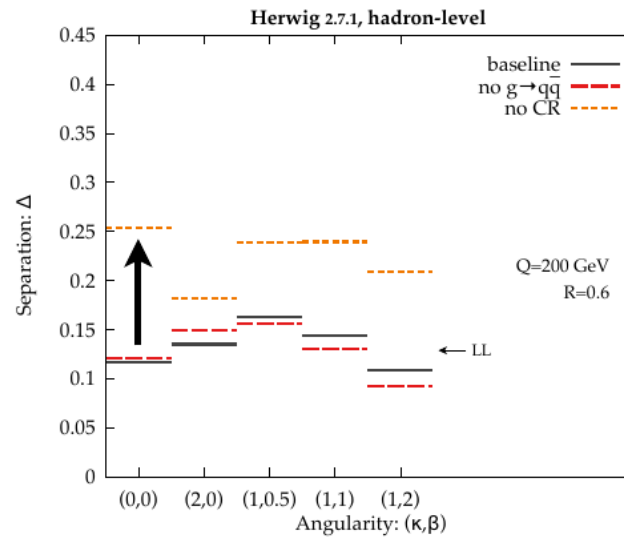
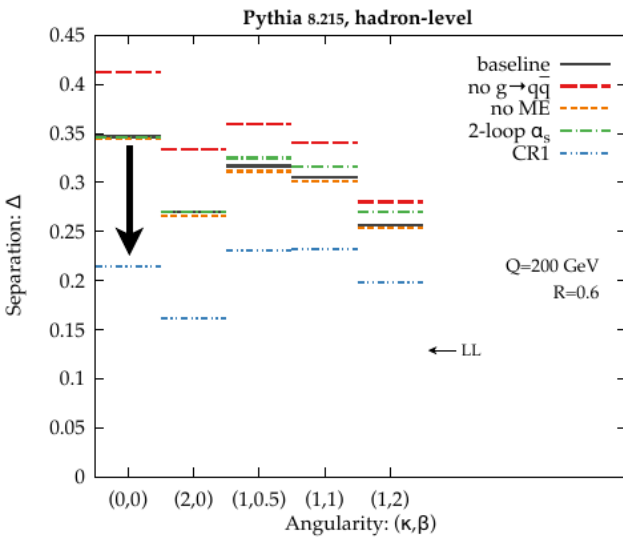


Large hadronisation effects (here for quarks)

Large differences between MCs also seen at parton level.

Interplay of perturbative and non-perturbative effects => challenge for both pQCD and NP models

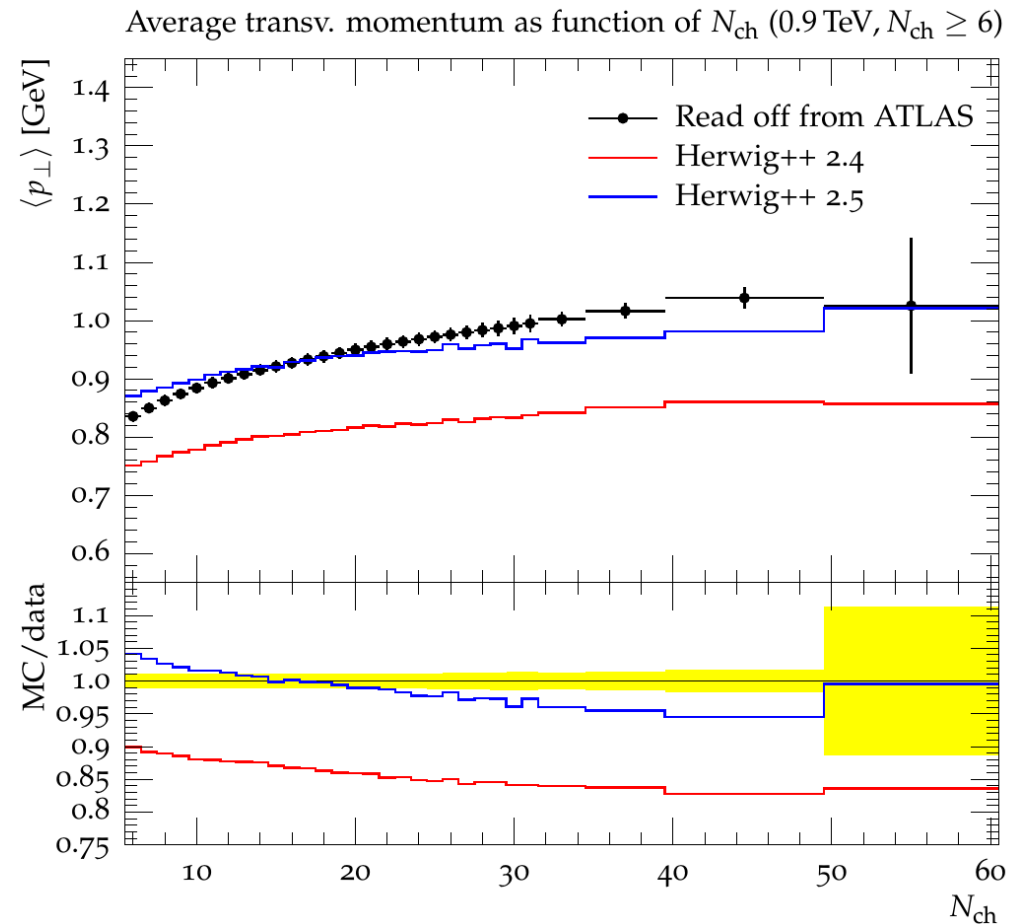
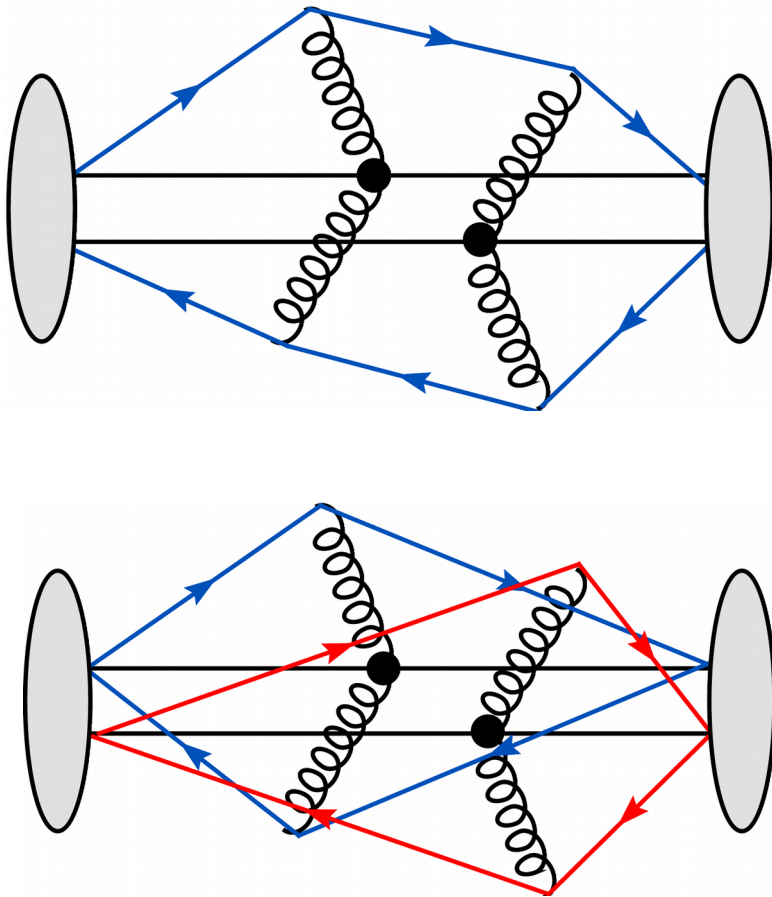
Separation power – sensitivity to MC options



Huge effect of color reconnection – very unexpected!

Herwig – Colour reconnection

- The least understood part of the Multiple Particle Interaction models.
- Needed to describe the Underlying Event and Min Bias data (sensitive to MPI phenomena)
- Crucial to constrain it, important for top mass, g/q gluon, ...



Herwig – Colour reconnection

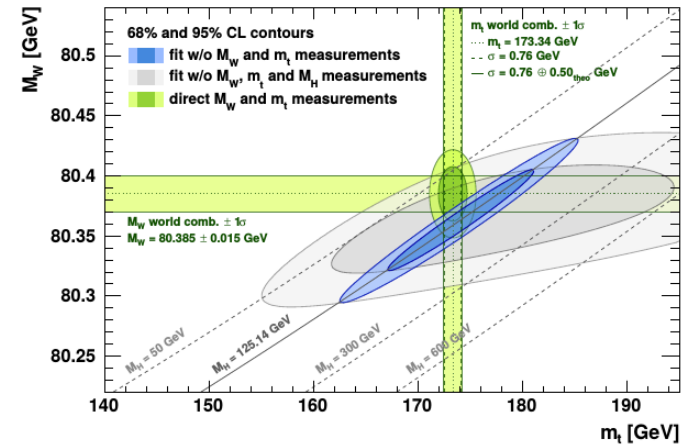
Top quark mass: precision matters

Precision tests of the Standard Model:

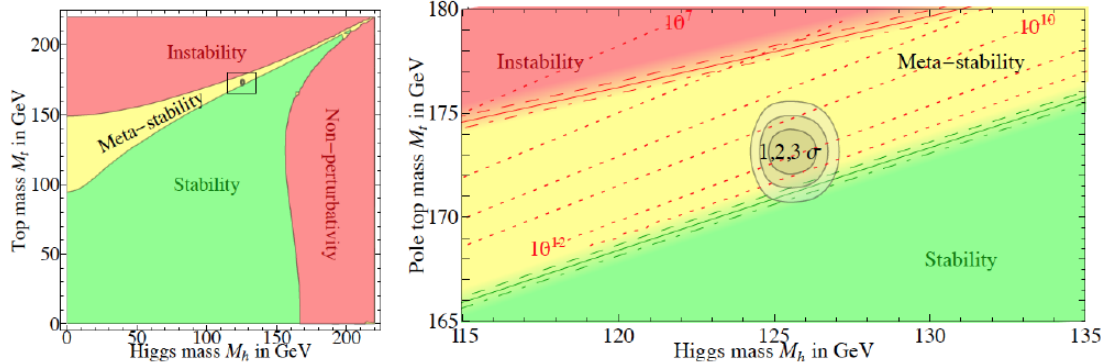
global EW fit *Rieman et al., Baak et al., ...*

↔ check self-consistency through
 m_t, m_W, m_H correlations

Gfitter Collab., Eur.Phys.J. C74 (2014) 3046



Degrassi et al, JHEP 1208 (2012) 098



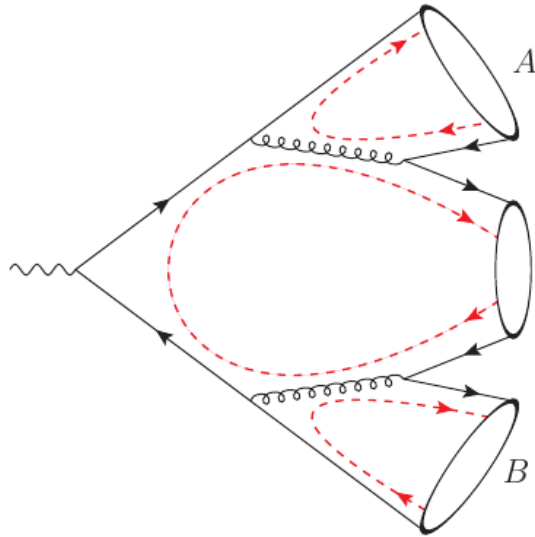
Stability of EW vacuum:
stable or meta-stable?

Different sources of uncertainties in m_t extraction via MC: accuracy of ME's, parton shower + hadronization, **color reconnection**, b -quark fragmentation ...

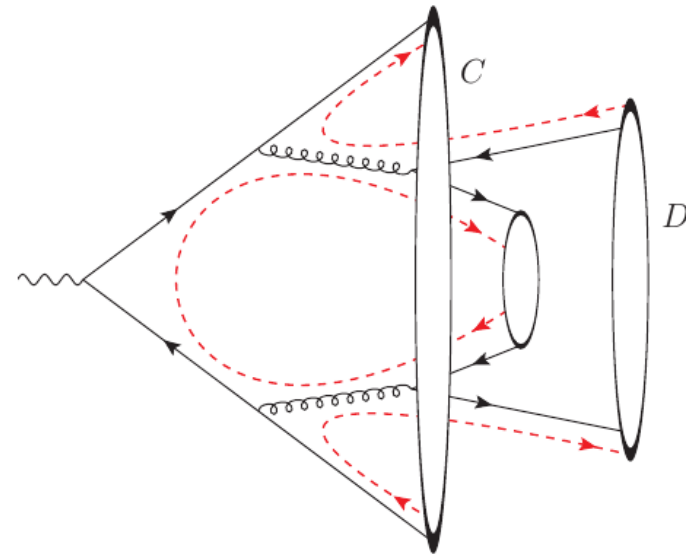
dominant source of uncertainty

Herwig – Colour reconnection

Cluster hadronization [Webber, Nucl. Phys. B238 (1984) 492]



- ▶ perturbative QCD provides *preconfinement* [Amati, Veneziano, Phys. Lett. B83 (1979) 87]
- ▶ i.e. small cluster masses
 $M_{cl} \gtrsim M_{parton 1} + M_{parton 2}$



- ▶ improved description of soft events/UE at hadron colliders: manually **reduce cluster masses**
- ▶ if $M_C + M_D < M_A + M_B$ accept alternative clustering with probability p_{reco} (model parameter)

[Gieseke, Rohr, AS Eur.Phys.J. C72 (2012) 2225]

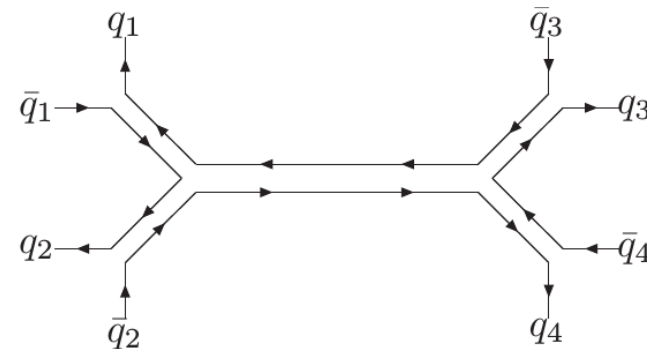
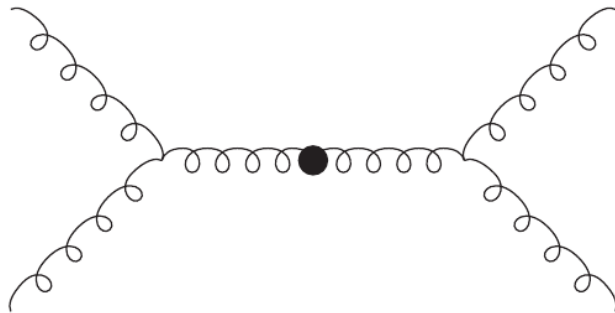
Strategy:

Data which has not previously been used.

1. Search for the LEP and LHC data sensitive to gluon jets.
 - Data on gluon jets in $e+e-$ collisions from the OPAL experiment [G. Abbiendi, et al.,: Phys.Rev.D69, 032002 and Eur. Phys. J. C37 (1), 25 (2004)]
 - In pp collisions from ATLAS [G. Aad, et al., Eur. Phys. J. C76 (6), 322 (2016)]
2. Improve the non-perturbative color reconnection model.
3. Improve the perturbative Parton Shower kinematics.

Herwig – Improvements of Color reconnection

- Possible that the color lines of a gluon produced at any stage of the shower can be reconnected leading to the production of a color-singlet object (see example below)

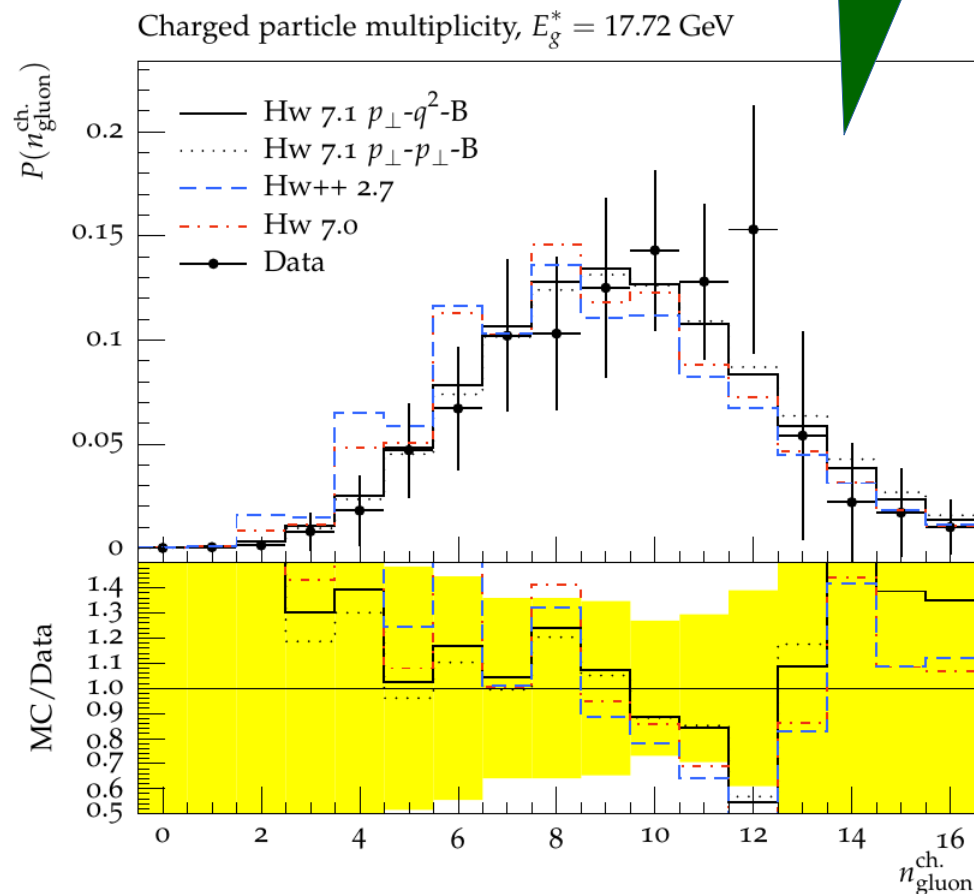
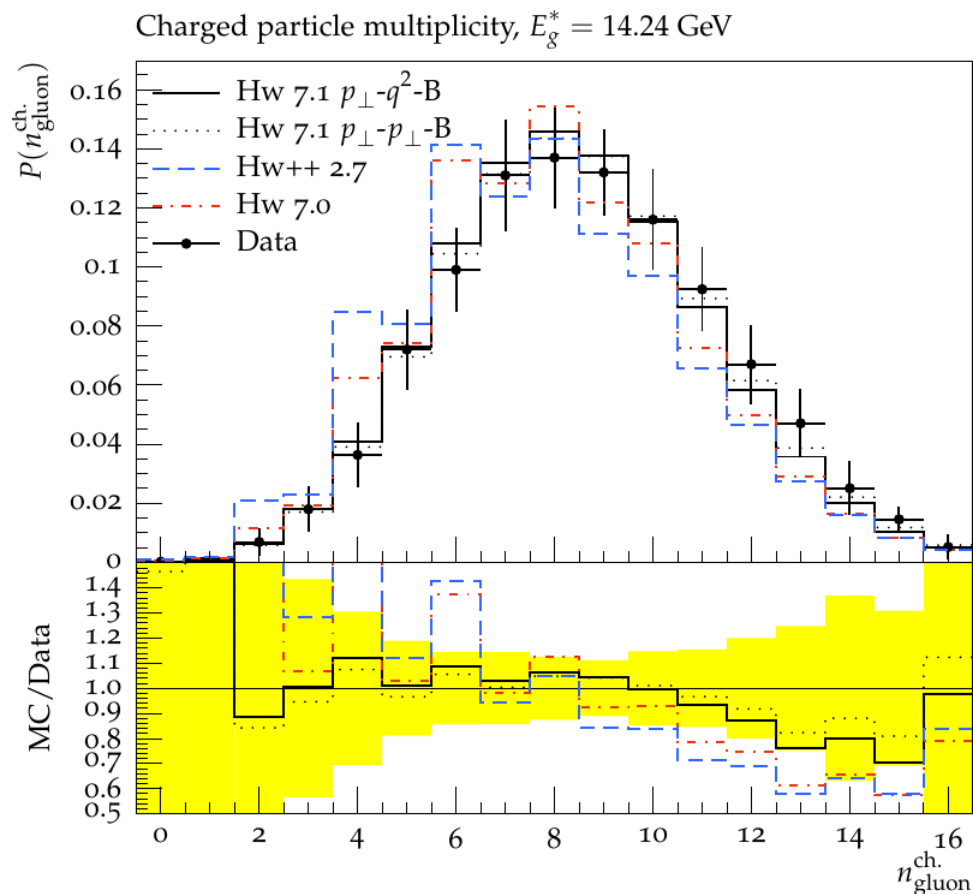


- Clusters containing partons from the parton shower of each of the original gluons, i.e. q_1, \bar{q}_3 and q_4, \bar{q}_2 , will have large masses and the rearrangement to give the clusters q_1, \bar{q}_2 and q_4, \bar{q}_3 will be kinematically favoured, although it means the original gluons will effectively become colour singlets rather than octets.
- this is physically possible we would expect that it occurs at a rate which is suppressed in the number of colours, N_c , as $1/N_c^2=1/9$, not the much higher reconnection rate $2/3$ which is current default value.
- We forbid the reconnection which would lead to a gluon produced in any stage of the parton-shower evolution becoming a colour-singlet after hadronization.

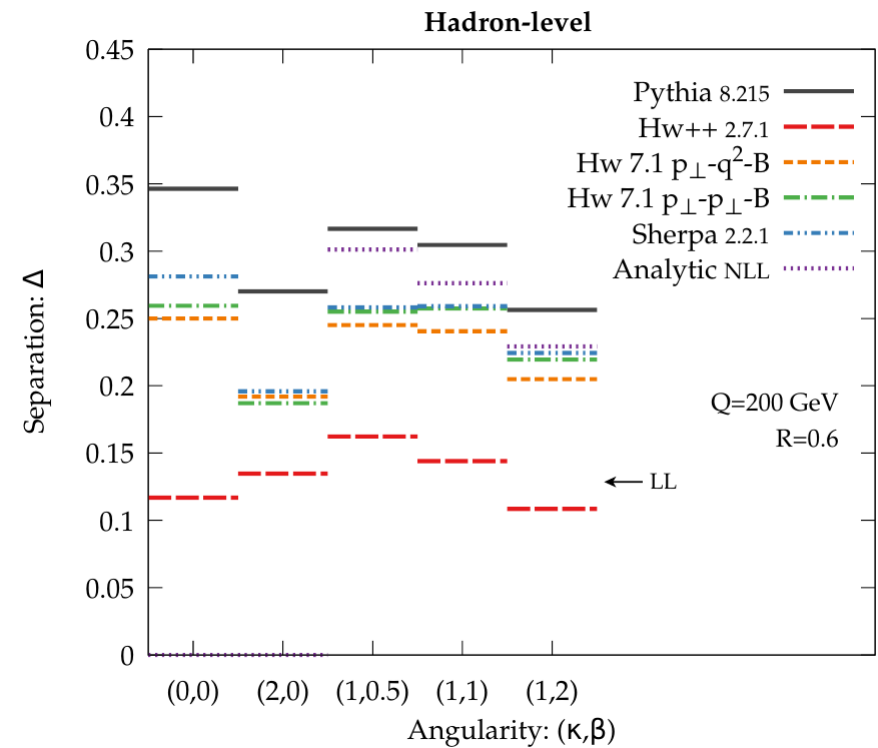
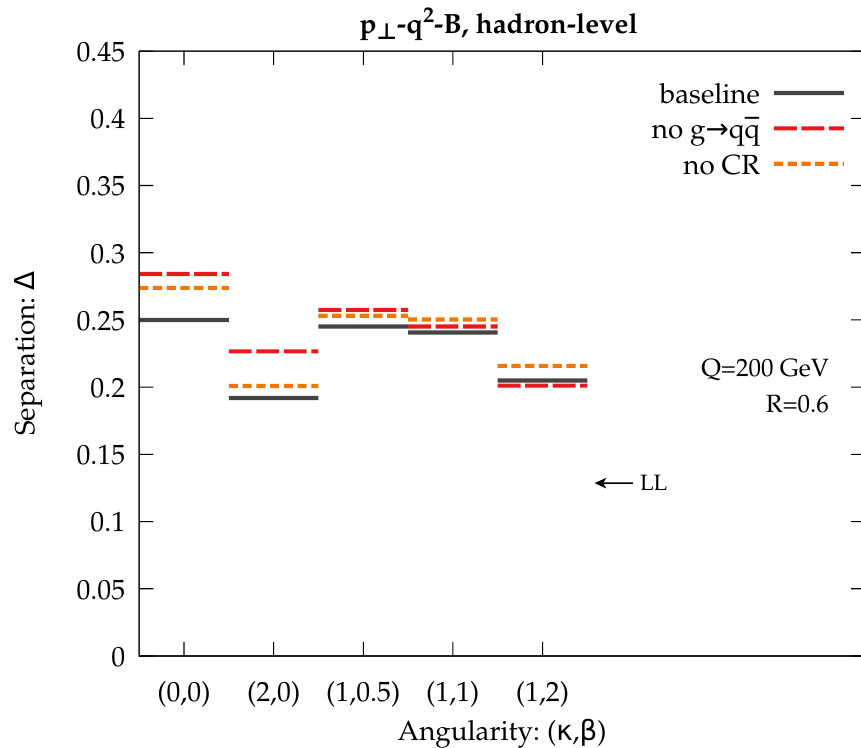
OPAL

Data which has not previously been used for tuning.

Multiplicity distribution of charged particles in gluon jets for two different gluon energies.



Idealized Quark/Gluon distributions



- Sensitivity to CR is gone especially for IRC safe observables – as expected.

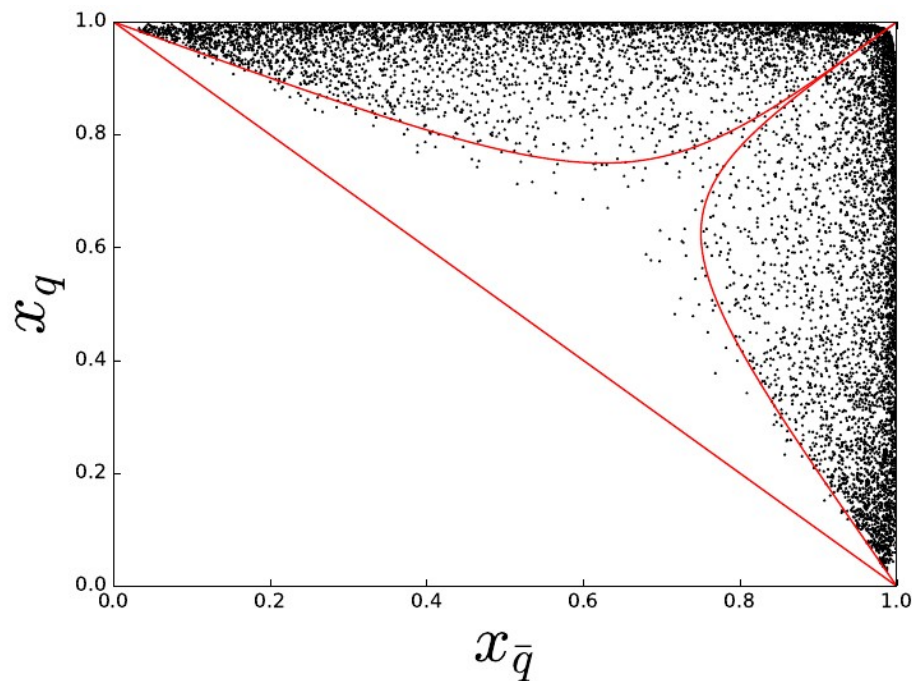
- Herwig is now more optimistic when it comes to distinguishing q/g jets.
- Spread of predictions is reduced.

Herwig – Parton Shower kinematics

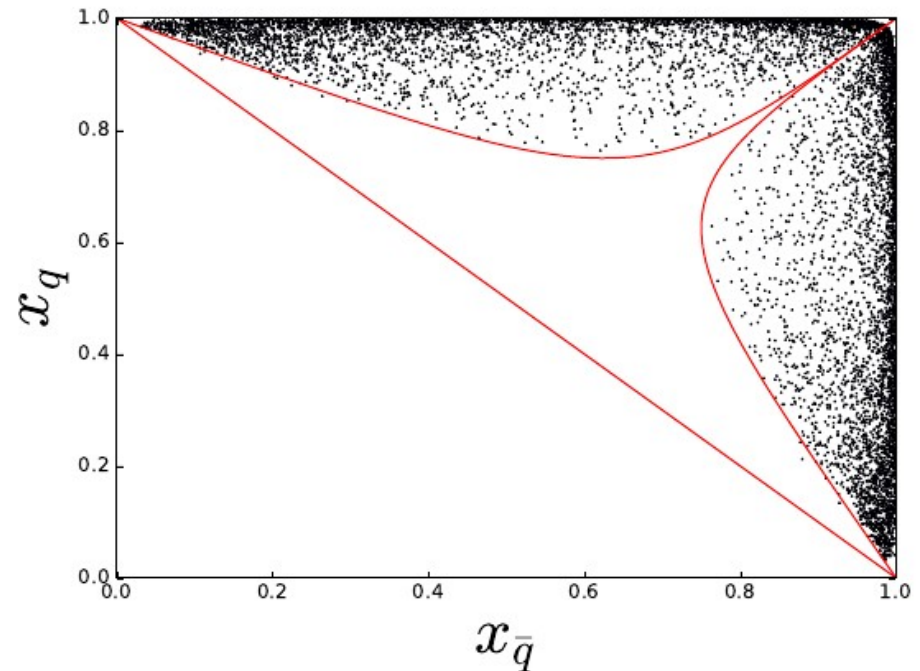
Choices that are formally subleading but can have a large effect on physical observables.

- Choice of minimal scale not fixed: we investigated p_T and q cut-off.
- Kinematics: choice of whether to preserve p_T or virtuality q^2 during the subsequent evolution.

Preserve p_T



Preserve virtuality q^2



Tuning

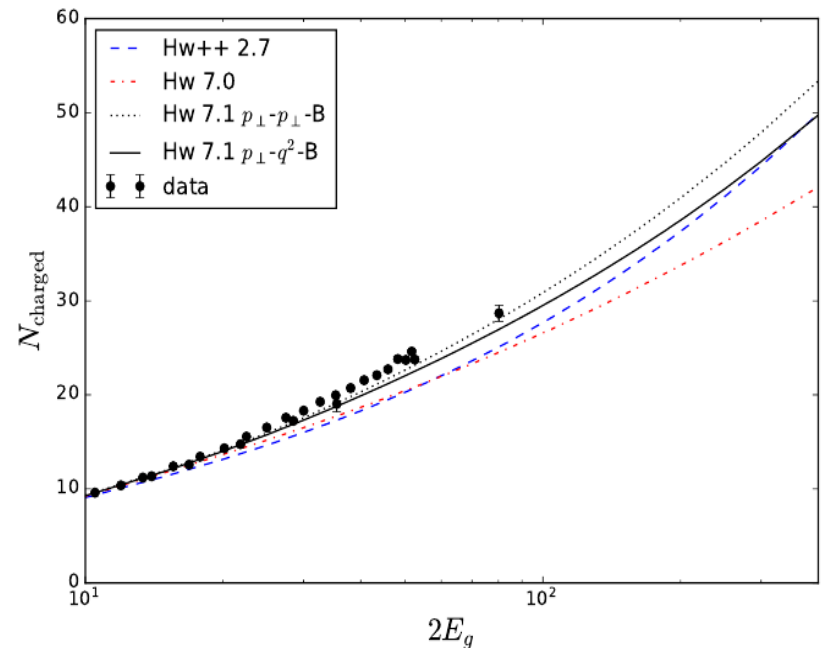
Unfortunately, when the PS is changed we need to retune the hadronization model:

12 tunes for different PS options and weights w on low-energy charged multiplicity data A($w=0$), B($w=100$), C($w=1000$)

Cut-Off Preserved Tune	p_{\perp}						Virtual Mass					
	A	p_{\perp} B	C	A	q^2 B	C	A	p_{\perp} B	C	A	q^2 B	C
	Tuning Observables											
Light quarks	4.4	4.3	6.7	3.0	2.9	4.2	7.8	7.6	6.9	4.6	4.3	3.6
Charm quarks	3.2	2.8	5.8	3.6	3.5	3.9	4.5	4.6	6.4	3.9	3.9	7.4
Bottom quarks	4.0	3.4	3.6	5.4	4.9	3.4	3.4	3.3	3.4	4.1	4.1	4.9
Gluons	1.1	1.1	1.5	1.1	1.1	1.4	1.2	1.2	1.2	1.3	1.2	1.5
	N_{charged}											
Gluon	14.2	18.6	22.6	26.9	37.1	60.0						
All quarks	4.6	2.7	2.7	3.4	2.5	5.2						
Light quarks	2.2	1.7	2.8	1.7	1.8	4.4						
Charm quarks	2.8	2.0	1.1	2.2	1.6	1.0						
Bottom quarks	20.4	18.1	15.8	24.1	21.3	15.7						
ATLAS Jets	3.2	0.9	4.3	13.3	10.1	7.8						

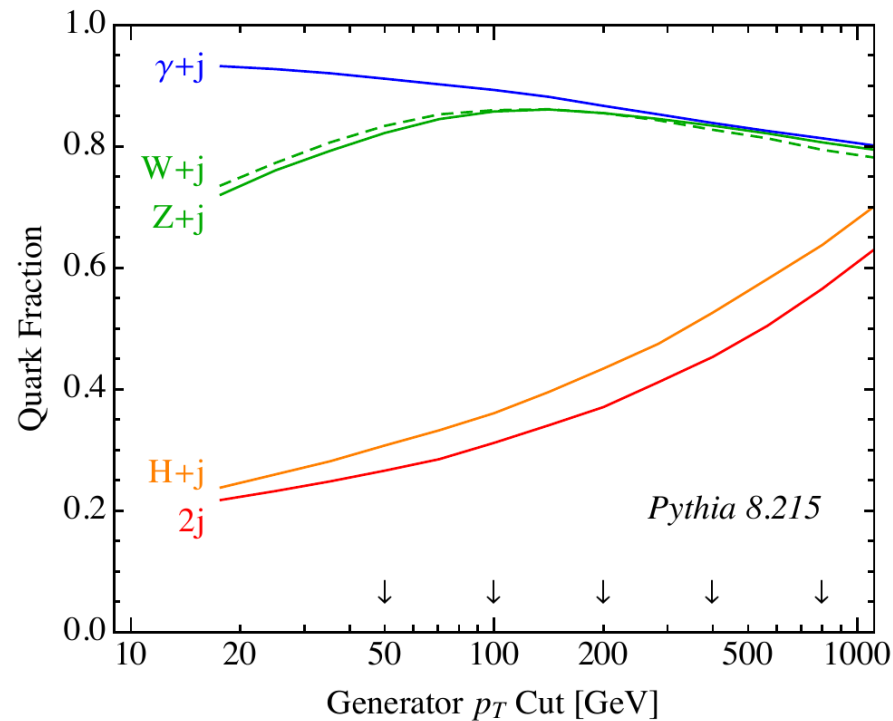
χ'^2

- the data on light quark jets, in particular event shapes measured at LEP favour preserving q^2 the data on the charged particle multiplicity in gluon jets favours preserving the p_T of the branching.



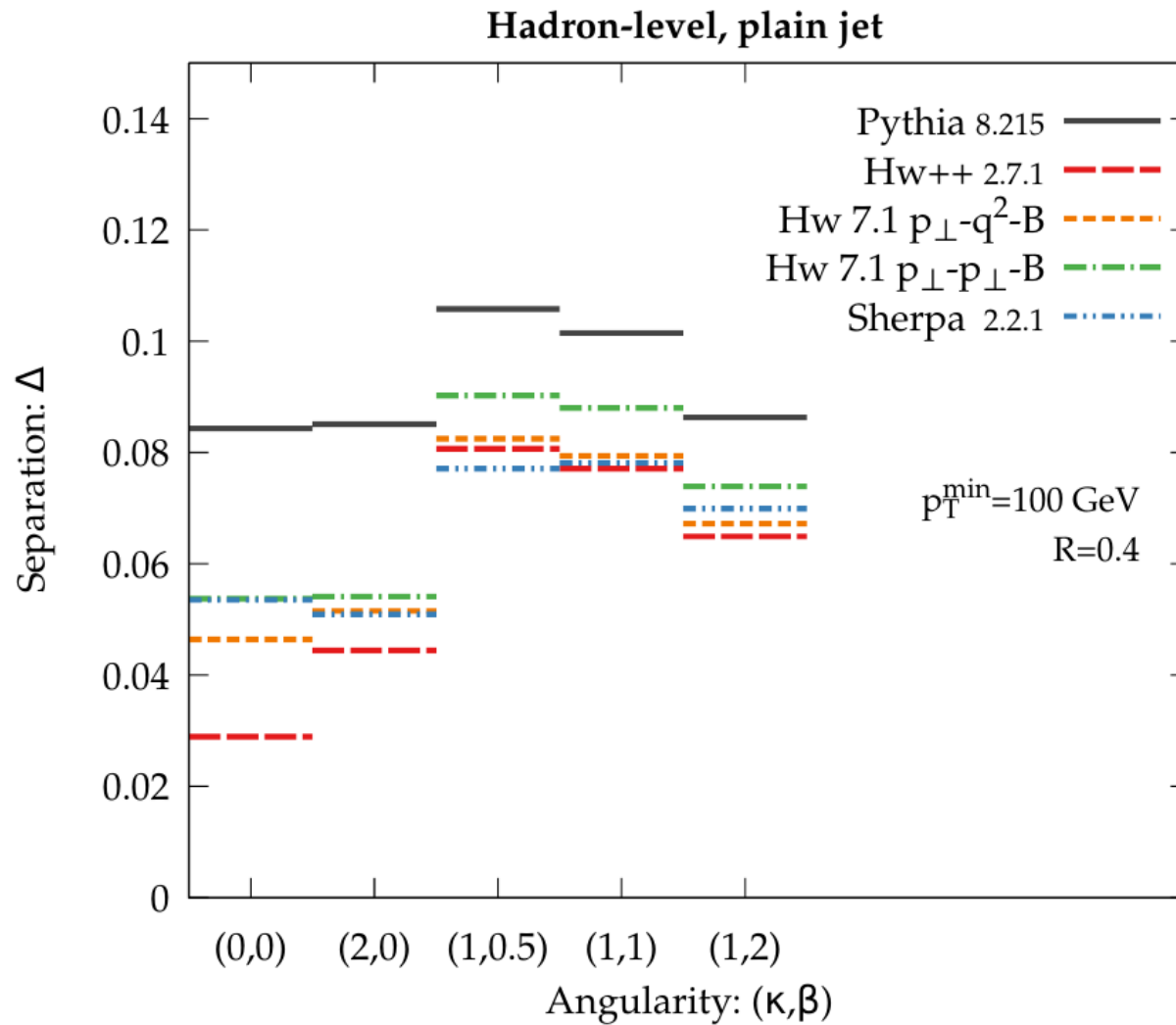
Evolution of # charged particles in gluon jets vs twice the energy of the gluon jet.

LHC 13 TeV



$$pp \rightarrow Z + j \text{ ("quark-enriched")} : \quad p_T^Z > p_T^{\min}, \quad \frac{p_T^{\text{jet}}}{p_T^Z} > 0.8, \quad |y_{\text{jet}} - y_Z| < 1.0.$$

$$pp \rightarrow 2j \text{ ("gluon-enriched")} : \quad \frac{p_{T,1} + p_{T,2}}{2} > p_T^{\min}, \quad \frac{p_{T,2}}{p_{T,1}} > 0.8, \quad |y_1 - y_2| < 1.0.$$



- Improvements of Herwig led to better discrimination power at the LHC (interesting would be to check against more q/g data, however most of them are not available to us).
- Spread of prediction reduced especially for IRC unsafe observables

Summary and outlook

1. The properties of quark and gluon jets, and the differences between them, are increasingly important at the LHC.
2. Quark jets well constrained by the LEP data, this was not the case for gluon jets.
3. We have performed a tuning the Herwig 7 event generator using data on gluon jets from LEP for the first time.
4. Improvements of perturbative and non-perturbative aspects of the simulation led to significantly better description of gluon jets, in particular their charge particle multiplicity.
5. However still there is a tension between the data on charged particle multiplicities, for both quark and gluon initiated jets, and the data on event shapes and particle spectra from LEP.
6. The tension might be resolved by improvements of the non-perturbative hadronization modelling.

Causality violation

[Information](#)[References \(53\)](#)[Citations \(1\)](#)[Files](#)[Plots](#)

A case study of quark-gluon discrimination at NNLL' in comparison to parton showers

Jonathan Mo (Amsterdam U. & NIKHEF, Amsterdam), Frank J. Tackmann (DESY), Wouter J. Waalewijn (NIKHEF, Amsterdam & Amsterdam U.)

Aug 2, 2017 - 10 pages

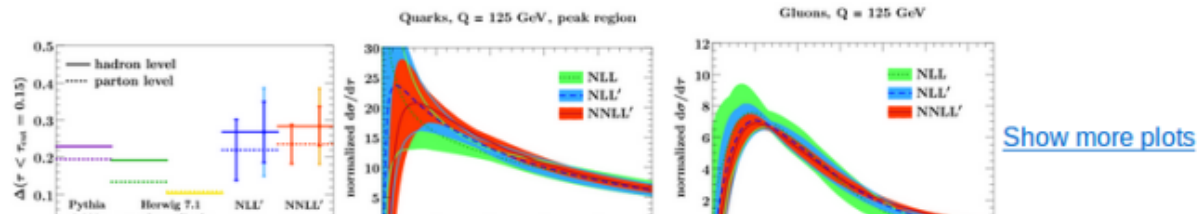
DESY-17-111, NIKHEF-2017-031
e-Print [arXiv:1708.00867](https://arxiv.org/abs/1708.00867) [hep-ph] | [PDF](#)

Abstract (arXiv)

Predictions for our ability to distinguish quark and gluon jets vary by more than a factor of two between different parton showers. We study this problem using analytic resummed predictions for the thrust event shape up to NNLL' using $e^+e^- \rightarrow Z \rightarrow q\bar{q}$ and $e^+e^- \rightarrow H \rightarrow gg$ as proxies for quark and gluon jets. We account for hadronization effects through a nonperturbative shape function, and include an estimate of both perturbative and hadronization uncertainties. In contrast to previous studies, we find reasonable agreement between our results and predictions from both Pythia and Herwig parton showers. We find that this is due to a noticeable improvement in the description of gluon jets in the newest Herwig 7.1 compared to previous versions.

Note: 10 pages, 5 figures

Keyword(s): INSPIRE: [parton: showers](#) | [gluon: jet](#) | [hadronization: effect](#) | [HERWIG](#) | [quark](#) | [event shape analysis](#) | [nonperturbative](#) | [quark gluon](#) | [thrust](#)

[Information](#)[References \(96\)](#)[Citations \(0\)](#)[Files](#)[Plots](#)

Improving the Simulation of Quark and Gluon Jets with Herwig 7

Daniel Reichelt (Dresden, Tech. U.), Peter Richardson (CERN & Durham U., IPPP), Andrzej Siodmok (Cracow, INP)

Aug 4, 2017 - 14 pages

CERN-TH-2017-174, IFJPAN-IV-2017-16, MCNET-17-13
e-Print [arXiv:1708.01491](https://arxiv.org/abs/1708.01491) [hep-ph] | [PDF](#)

Abstract (arXiv)

The properties of quark and gluon jets, and the differences between them, are increasingly important at the LHC. However, Monte Carlo event generators are normally tuned to data from e^+e^- collisions which are primarily sensitive to quark-initiated jets. In order to improve the description of gluon jets we make improvements to the perturbative and the non-perturbative modelling of gluon jets and include data with gluon-initiated jets in the tuning for the first time. The resultant tunes significantly improve the description of gluon jets and are now the default in Herwig 7.1.

Thank you for your attention!

Tuning

Unfortunately, when the PS is change we need to retune the hadronization model which is a big effort.

Cut-Off	p_{\perp}						Virtual Mass					
Preserved Tune	A	p_{\perp} B	C	A	q^2 B	C	A	p_{\perp} B	C	A	q^2 B	C
Bottom quark hadronization parameters												
CIMaxBottom	4.655			3.911			4.0612			4.163		
CIPowBottom	0.622			0.638			0.9475			0.590		
PSplitBottom	0.499			0.531			1.9568			1.881		
CISmrBottom	0.082			0.020			0.04			0.040		
SingleHadronLimitBottom	0.000			0.000			0.0204			0.000		
Charm quark hadronization parameters												
SingleHadronLimitCharm	0.000			0.000			0.078			0.012		
CIMaxCharm	3.551			3.638			3.805			3.885		
CIPowCharm	1.923			2.332			2.242			2.452		
PSplitCharm	1.260			1.234			1.895			1.767		
CISmrCharm	0.000			0.000			0.000			0.000		
Light quark hadronization and shower parameters												
AlphaMZ ($\alpha_s^{\text{CMW}}(M_Z)$)	0.1094	0.1087	0.1126	0.1260	0.1262	0.1265	0.1221	0.1218	0.1184	0.1314	0.1317	0.1254
pTmin	1.037	0.933	0.809	1.301	1.223	0.992	N/A			N/A		
aParameter	N/A			N/A			0.367			0.234		
cutoffKinScale	N/A			N/A			2.939	2.910	2.294	3.277	3.279	1.938
CIMaxLight	3.504	3.639	4.349	3.058	3.003	3.197	3.328	3.377	3.846	3.414	3.427	3.477
CIPowLight	2.576	2.575	1.226	1.513	1.424	2.786	1.286	1.318	2.063	2.766	2.792	2.35
PSplitLight	1.003	1.016	0.855	0.885	0.848	0.648	1.198	1.185	1.277	1.346	1.333	2.015
PwtSquark	0.552	0.597	1.167	0.602	0.666	1.024	0.721	0.741	0.782	0.626	0.646	1.15
PwtDIquark	0.369	0.344	0.181	0.416	0.439	0.512	0.277	0.273	0.246	0.321	0.328	0.366