### Light long-lived particles at colliders

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# **Fermilab**

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based on

ongoing work, arXiv:2303.XXXXX with Bogdan Dobrescu

JHEP 12 (2022) 005, arXiv:2203.08824 with Kai Böse, Torben Ferber, Christopher Hearty, Felix Kahlhoefer, Alessandro Morandini and Kai Schmidt-Hoberg

### New physics off the beaten track

- LHC has collected an enormous amount of data already:
  - 137  $fb^{-1}$  in Run 2; on its way to collecting 300  $fb^{-1}$  in Run 3
- Still, this is **only a small fraction** of the final luminosity goal



Active frontier with a great deal of recent progress:

### **Long-lived particles**

i.e. particles with macroscopic decay lengths of  $\mathcal{O}(mm)$  to  $\mathcal{O}(km)$ 

### Landscape of displaced decays

Organize approximately by mass and decay position



see e.g. Alimena et al., 1903.04497 for review

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## Light long-lived particles

In this talk: Two different ways to close this gap,

illustrated with two models





Light LLPs from vectorlike leptons at the LHC, in particular **in the muon chambers**  Light LLPs from strongly interacting dark sector at **Belle II** 

- In both cases, detectors are used in new, creative ways that they were not originally intended for.
- In both cases, light particles are long-lived due to a mass hierarchy.

### What heavy new physics could drive a macroscopic LLP lifetime?

- Heavy new **bosons**: second part of this talk
- Heavy new fermions: have to be vectorlike

(left- and right-handed fields have the gauge charges)

due to stringent constraints on new chiral fermions

- Vectorlike fermions among the simplest possible additions to the SM
- Novel form of matter
- Part of a large number of top-down BSM models (grand unification, compositeness, extra dimensions ...)

Most minimal possibility: Singlet vectorlike lepton  ${\mathscr C}$ 

### Vectorlike leptons: standard story

• Gauge eigenstates: vectorlike lepton  $\mathscr{E}_L$ ,  $\mathscr{E}_R$ 

(1, 1, -1) under SU(3)xSU(2)xU(1)

see e.g. Kumar & Martin, 1510.03456

• Couplings of  $\mathscr E$  to third-generation leptons:

$$-m_{\mathcal{E}\mathcal{E}}\overline{\mathcal{E}}_{L}\mathcal{E}_{R}-m_{\mathcal{E}3}\overline{\mathcal{E}}_{L}e_{R}^{3}-y_{3}H\overline{\ell}_{L}^{3}e_{R}+\mathrm{H.c}$$

Mass mixing

physical particles: au and au'

Q

• Mass eigenstate au' decays via channels

 $\tau' \rightarrow \tau Z, \ \tau' \rightarrow \tau h, \ \tau' \rightarrow \nu W$ 

ATLAS and CMS searches exclude these standard VLLs if  $m_{\tau'} \lesssim 176 \text{ GeV}$ 

ATLAS collaboration, arXiv:1506.01291

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ATLAS-CONF-2022-044

### Long-lived particles from vectorlike leptons

#### How can vectorlike leptons be connected to long-lived particles?

- Consider complex scalar  $\phi$ , responsible for VLL mass
- Most general Yukawa interaction:

$$\begin{split} &-\phi \,\overline{\mathscr{C}}_L \left( y_{\mathscr{C}} e^{i\beta_{\mathscr{C}}} \, \mathscr{C}_R + y_o \, e^{i\beta_o} \, e_R^3 \right) + \mathrm{H.c} \, . \\ & y_{\mathscr{C}}, y_o > 0 \quad \text{and} \quad 0 \leq \beta_{\mathscr{C}}, \beta_o < 2\pi \end{split}$$

• After scalar acquires vacuum expectation value:

$$\phi = \left( v_{\phi} + \frac{1}{\sqrt{2}} \varphi_{\tau} \right) \exp \left( i a_{\tau} / v_{\phi} \right)$$
pseudoscalar
vev scalar

### Long-lived particles from vectorlike leptons

• VEVs  $v_H$ ,  $v_\phi$  and explicit mass terms  $m_{\mathcal{EE}}$ ,  $m_{\mathcal{E3}}$  control mixing with SM third generation

$$(\overline{e}_{L}^{3}, \overline{\mathscr{C}}_{L}) \begin{pmatrix} y_{3} v_{H} & 0 \\ m_{\mathscr{C}3} + y_{o} e^{i\beta_{o}} v_{\phi} & m_{\mathscr{C}\mathscr{C}} + y_{\mathscr{C}} e^{i\beta_{\mathscr{C}}} v_{\phi} \end{pmatrix} \begin{pmatrix} e_{R}^{3} \\ \mathscr{C}_{R} \end{pmatrix}$$
$$\begin{pmatrix} e_{L,R}^{3} \\ \mathscr{C}_{L,R} \end{pmatrix} = \begin{pmatrix} c_{L,R} & s_{L,R} \\ -s_{L,R} & c_{L,R} \end{pmatrix} \begin{pmatrix} \tau_{L,R} \\ \tau'_{L,R} \end{pmatrix}$$

Constraint on mixing angle:  $s_L < 0.037$  from  $\frac{D(Z)}{D(Z)}$ 

$$\frac{B(Z \to \tau^+ \tau^-)}{B(Z \to \ell^+ \ell^-)}$$

• Interesting new BSM decay modes of VLL mass eigenstate:

$$au' 
ightarrow au a_{ au}$$
 ,  $au' 
ightarrow au arphi_{ au}$ 

in addition to standard SM modes  $\tau' \rightarrow \tau Z$ ,  $\tau' \rightarrow \tau h$ ,  $\tau' \rightarrow \nu W$ 

### Long-lived particles from vectorlike leptons

**Case 1:** Pseudo-Goldstone limit: entire mass of VLL is generated by  $\phi$ 



 $a_{\tau}$  is naturally light pseudo-Goldstone boson

 Ratio of BSM-to-SM branching fractions of VLL is set purely by VEVs, independent of Yukawa coupling:

$$\frac{\Gamma\left(\tau' \to \tau a_{\tau}\right)}{\Gamma\left(\tau' \to \mathrm{SM}\right)} = \frac{v_H^2}{4v_{\phi}^2} \left(1 + \frac{M_h^2}{2m_{\tau'}^2} + O(M_h^4/m_{\tau'}^4)\right)$$

explored in ongoing separate project

### In this talk focus on:

**Case 2:** With explicit mass terms  $m_{\mathscr{C}}$ ,  $m_{\mathscr{C}}$ ,  $m_{\mathscr{C}}$ 

Large parameter space with BR  $(\tau' \rightarrow \tau a_{\tau}) \approx 100 \%$ 

### Long-lived particle production



Potential LHC reach in au' mass to well above  $1 \,\, TeV$ 

## Long-lived particle decay

• If  $a_{\tau}$  is lighter than  $\tau'$ , its **dominant decay mode is**  $a_{\tau} \to \gamma \gamma$  **via heavy**  $\tau'$  **loop** 

$$a_{\tau} = -\frac{\tau'}{\tau'}$$

 Heavy VLL in the loop leads to macroscopic a<sub>τ</sub> decay length even for sizable couplings:

$$c\tau_a = 0.23 \text{ m} \times \left(\frac{0.1}{y_{\tau'}}\right)^2 \left(\frac{1 \text{ GeV}}{m_{a_{\tau}}}\right)^3 \left(\frac{m_{\tau'}}{400 \text{ GeV}}\right)^2$$

• In addition,  $a_{\tau}$  picks up substantial boost

Majority of decays can easily occur meters from the interaction point

### Search for LLPs in muon system

### The muon system is several meters away!

- Recent CMS search 2107.04838 and ATLAS search 1811.07370
- **CMS search** first search to use muon system like a calorimeter for LLPs



Sensitivity depends on LLP energy, not on mass

CMS collaboration, 2107.04838

Unique sensitivity to decay positions several meters away from interaction point

## Projected LLP sensitivity in muon system

- Motivated by VLL+LLP model, propose a search for LLPs in the muon system in association with prompt taus
- Include full muon system with barrel + endcaps, leverage prompt taus for trigger and background suppression



Vast improvement over sensitivity of existing search

## Complementarity with missing energy searches

• LLPs decaying to photons after the ECAL **appear as missing energy** 

complementary constraint from missing energy searches



LLP search superior for shorter lifetimes and higher VLL masses and scales more favorably with luminosity



Intermediate LLP decay lengths  $\sim$  tens of centimeters to meters are challenging to cover at the LHC



#### One possibility (we just saw):

Use parts of the detector in novel ways and avoid vertex reconstruction

Other possibility:

Different experiments,

in particular B factories

### Dark sector searches at the intensity frontier

**Belle II** experiment at SuperKEKB  $e^+e^-$  collider with  $\sqrt{s} \sim 10.6 \text{ GeV}$ 

- Among the world-leading probes of B mesons (along with LHCb)
- Very well suited also for searching for light LLPs beyond the SM

see e.g. Belle II physics book, arXiv:1808.10567

- Advantages over LHCb with respect to light LLPs:
  - Longer decay lengths in the lab system covered (up to  $\sim 60~{
    m cm}$ )
  - Typically lower LLP boost due to smaller c.o.m. energy



sensitive to much longer proper decay lengths

Particularly striking example: dark showers

### Strongly interacting dark sector with effective portal

• Dark showers occur in **strongly interacting dark sectors** resembling SM QCD sector:

$$\mathscr{L}_{\text{dark QCD}} \supset -\frac{1}{4} F^a_{\mu\nu} F^{\mu\nu,a} + \overline{q_d} i \mathscr{D} q_d - \overline{q_d} M_q q_d$$

• Below the dark confinement scale:

**Dark pions**  $\pi_d$  are excellent DM candidates. see e.g. Kribs & Neal arXiv:1604.04627 for review **Dark vector mesons**  $\rho_d$  are generically unstable:

can set dark pion relic abundance via "forbidden" annihilations  $\pi_d \pi_d \rightarrow \rho_d \rho_d$ gives rise to **visible signals at accelerator experiments** EB et al. arXiv:1907.04346

• Mediator with mass  $\gg \sqrt{s_{\text{Belle II}}}$  can be described by an **effective interaction** between dark quarks and SM fermions:

$$\mathscr{L}_{\rm eff} \supset \frac{1}{\Lambda^2} \sum_f q_f \bar{f} \gamma^{\mu} f \bar{q}_{\rm d} \gamma_{\mu} q_{\rm d}$$

 $\Lambda$  : scale of the effective interaction

## Dark showers at Belle II

• If dark confinement scale  $\ll \sqrt{s}$ , production of dark quarks via effective interaction leads to **dark shower** and production of dark mesons



• Signature of light dark showers:

(multiple) displaced decays number varying from event to event

 Rate of dark shower production through effective operator scales as

$$\sigma(e^+e^- \to q_{\rm d}\bar{q}_{\rm d}) \propto \frac{s}{\Lambda^4}$$

• The  $\rho_d^0$  mesons decay to visible SM particles with decay length

$$c\tau_{
ho_{\rm d}^0} \propto rac{\Lambda^4}{m_{
ho_{\rm d}}^5}$$



Parameter space of low-energy effective theory only consists of  $\Lambda$  and  $m_{\rho_{\rm d}}$ .

### Displaced vertex search at Belle II

- No dedicated search for dark showers at B factories yet, but model-independent searches for LLPs
- **Prospective search for displaced vertices at Belle II** Duerr et al., arXiv:1911.03176 and arXiv:2012.08595, and EB et al. arXiv:2203.08824

**Signature:** displaced decays to pairs of oppositely charged leptons, pions or kaons **Transverse distance** 0.2 cm < R < 60 cm **Trigger** on sum of energy depositions (E > 1 GeV), two tracks or one high- $p_T$  muon **Selection criteria:** 

electron pairs	muon pairs
$\begin{array}{c} p(e^+), p(e^-) > 0.1 \ {\rm GeV} \\ m_{e^+e^-} > 0.03 \ {\rm GeV} \\ \alpha(e^+,e^-) > 0.025 \ {\rm rad} \end{array}$	$p(\mu^+),  p(\mu^-) > 0.05  { m GeV}$ $m_{\mu^+\mu^-} < 0.48  { m GeV}  or  m_{\mu^+\mu^-} > 0.52  { m GeV}$
Displaced vertex position	
$\begin{array}{ll} 0.2\ {\rm cm} < R < 0.9\ {\rm cm}\ or\ 17\ {\rm cm} < R < 60\ {\rm cm} & 0.2\ {\rm cm} < R < 60\ {\rm cm} \\ -55\ {\rm cm} \le z \le 140\ {\rm cm} \\ 17^\circ \le \theta_{\rm lab} \le 150^\circ \end{array}$	

Negligible background expected

### Belle II LLP sensitivity for dark showers

#### Projected sensitivity for dark showers of model-agnostic Belle II search





**Displaced vertex search at Belle II** can greatly improve over BaBar in cosmologically viable parameter space even with just  $500 \text{ fb}^{-1}$  of data.

### **Complementary constraints**

#### Additional constraints on the low-energy effective theory:



- Existing searches for visible resonances at BaBar and KLOE yield complementary constraints on promptly decaying dark mesons.
- Events with invisible dark shower and ISR photon lead to single photon signal with continuous energy spectrum.

Projection weaker than for bump hunt but competitive at small LLP mass

#### How does Belle II compare to LHCb?



 Relevant LHCb search: model-agnostic search for displaced dimuon resonances

LHCb collaboration, arXiv:2007.03923

## Belle II vs LHCb

LHCb sensitivity depends on combination of mediator couplings underlying



- Belle II sensitive to much longer decay lengths ( $c\tau \gtrsim O(\text{meters})$ )
- Strong case for Belle II search even in least favorable scenario

### Conclusions

- Long-lived particles are extremely well-motivated and could still be discovered at the LHC
- Rapidly expanding search program, but many remaining gaps especially light LLPs, decay lengths ~ meters
- Wealth of interesting models; here: VLL with LLP, dark showers
- Ways to close gap:
  - Creative use of existing detectors, in particular muon system
  - Other, lower-energy experiments, in particular Belle II