

Quasi-local energy and compactification

Raquel Santos García

INVISIBLES₁₈ WORKSHOP

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Work in collaboration with
Enrique Alvarez, Jesus Anero and Guillermo Milans del Bosch

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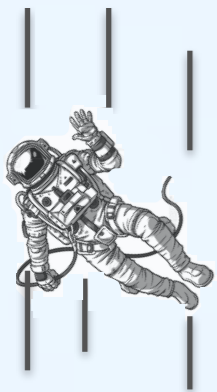


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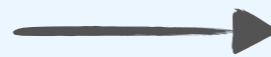
Energy in General Relativity

The four-momentum cannot be a **four-vector** because it can always be made to **vanish locally** in a free falling frame.



Free falling
observer

Equivalence
principle

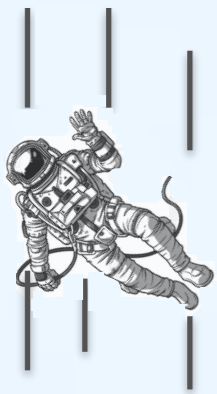


Zero gravitational
energy



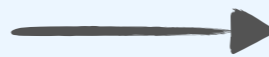
Energy in General Relativity

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Zero gravitational
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The energy momentum tensor is **covariantly** conserved

$$\nabla_{\mu} T^{\mu\nu} = 0$$

$$\underbrace{\partial_{\mu} T^{\mu\nu}}_{\text{Conservation in the usual sense}} = - \underbrace{\Gamma^{\mu}_{\mu\lambda} T^{\lambda\nu} - \Gamma^{\nu}_{\mu\lambda} T^{\lambda\mu}}_{\text{The purely gravitational contribution is not a tensor}}$$

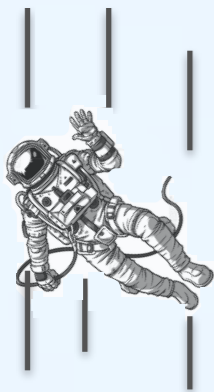
Conservation in the
usual sense

The purely **gravitational
contribution** is not a
tensor



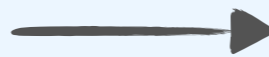
Energy in General Relativity

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Free falling
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**Ground state of a Quantum
theory of Gravity?**



We need a **definition of energy**
for the gravitational field



Quasi-local gravitational energy

We want to compute a quantity which can be ‘locally’ sensitive to the strength of the gravitational field.

Quasi-local gravitational energy

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We associate to a given hypersurface of a spacetime, the **integral of the trace of the extrinsic curvature**.

The **quasi-local energy** (QLE) is defined as

$$Q(\Sigma) \equiv \int_{\Sigma} K$$

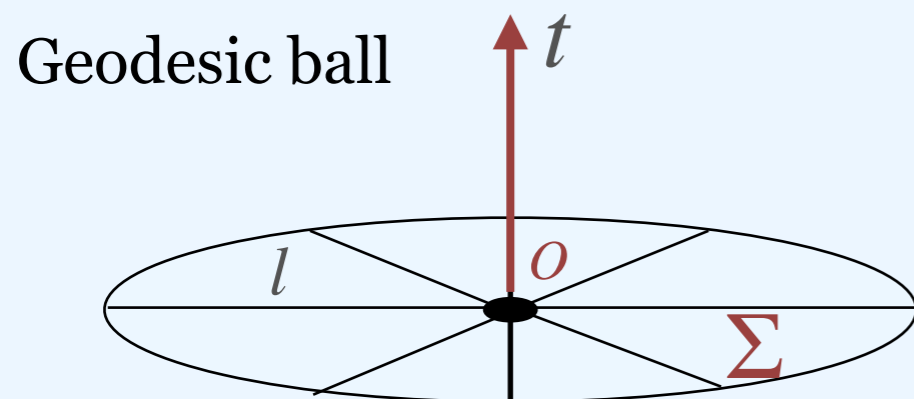
Dependent on the **tangent and normal vectors** of the hypersurface

Quasi-local gravitational energy

We want to compute a quantity which can be ‘locally’ sensitive to the strength of the gravitational field.

We associate to a given hypersurface of a spacetime, the **integral of the trace of the extrinsic curvature**.

We can construct the surface as



We take a point and send geodesics of fixed length normal to the time direction.

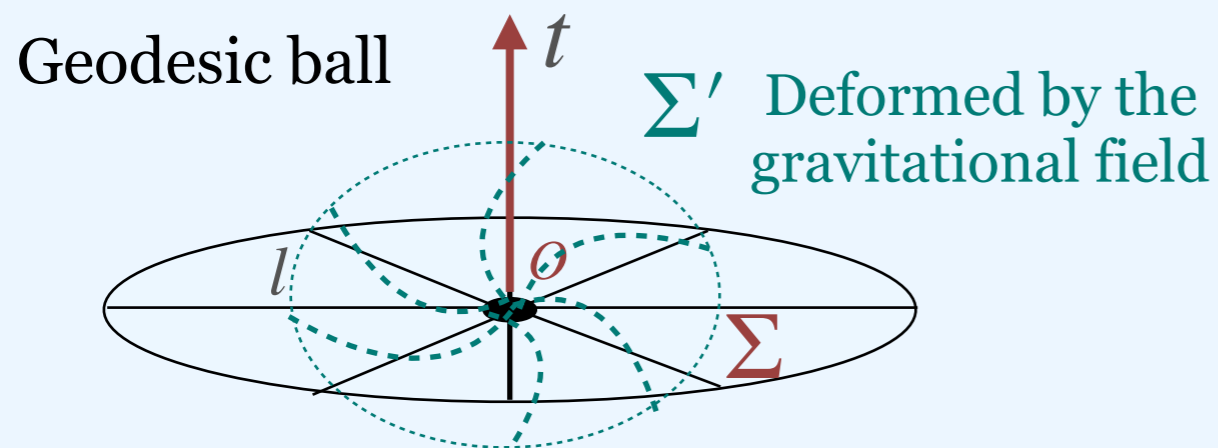
This defines a spacelike hypersurface Σ

Quasi-local gravitational energy

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The extrinsic curvature of the hypersurface is **sensitive** to the strength of the **gravitational field**

QLE and compactification

Idea

The QLE definition of Brown and York is able to discriminate between the **uncompactified** Minkowski spacetime and the spacetimes with **compact dimensions**.

QLE and compactification

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The QLE definition of Brown and York is able to discriminate between the **uncompactified** Minkowski spacetime and the spacetimes with **compact dimensions**.

Strategy

Compute the QLE for 3-spheres in M_5

Compute the QLE for '3-spheres' in $M_4 \times \underbrace{S_1}$

One compact dimension

$$-\pi l \leq y \leq \pi l$$

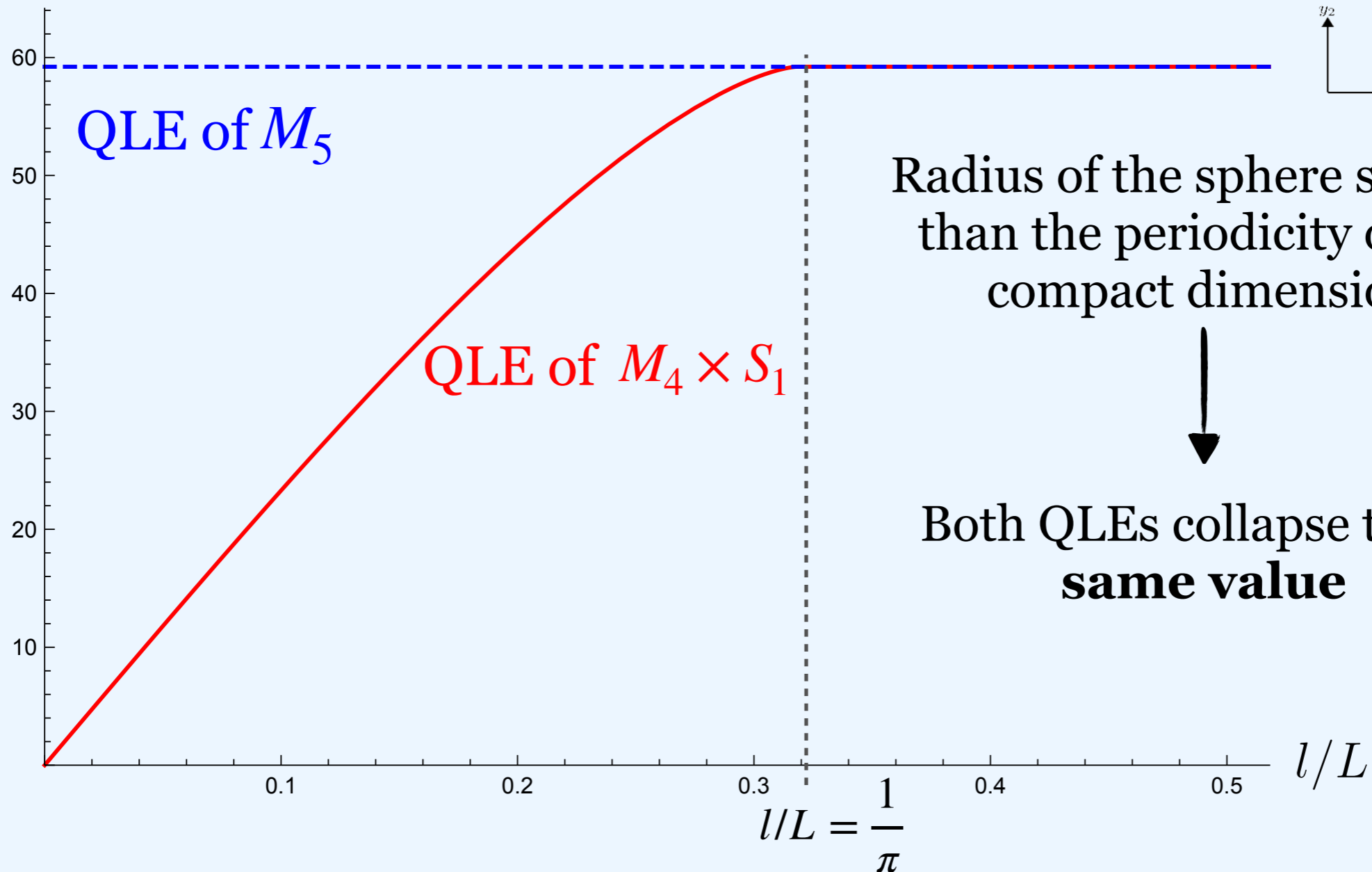
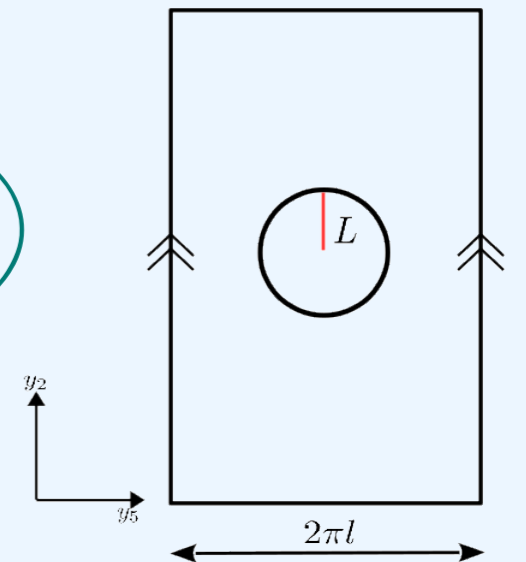
QLE and compactification

l Radius of the compact dimension

L Radius of the hypersurface

$$l/L \geq \frac{1}{\pi}$$

Q/L^2



Radius of the sphere smaller than the periodicity of the compact dimension



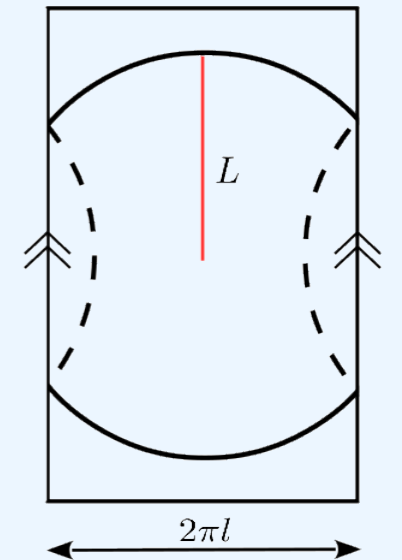
Both QLEs collapse to the **same value**

QLE and compactification

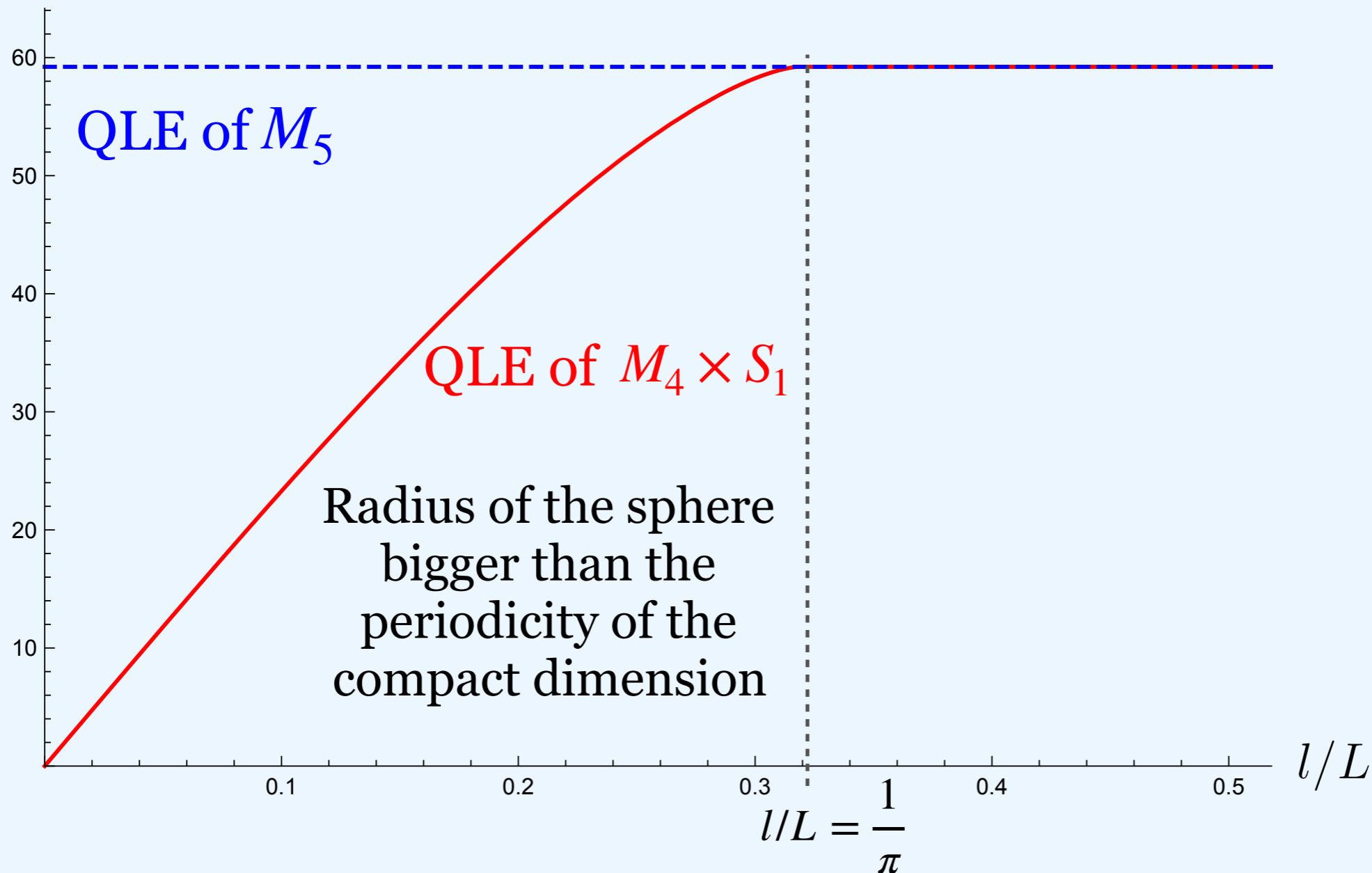
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Q/L^2



Summary and Outlook

The **quasi-local energy** could provide an **energetical argument** in favour of compactified or uncompactified spacetimes.

This is a proof of concept, but more general setups need to be studied, such as the introduction of fluxes in order to stabilise the compact dimension.

Work is ongoing regarding the use of the **quasi-local energy** as a tool to compute the **total energy** of full spacetimes in a **more covariant way**.

BACKUP

QLE and compactification

Codim-2 spheres M_5

Codim-2 spheres $M_4 \times S_1$

For both, the embedding

$$y_1 = T, \quad \sum_{i=2}^{i=5} (y_i)^2 \equiv L^2.$$

The two **normal** vectors and the three **tangent** vectors read

$$n_A \equiv \left(\frac{\partial}{\partial t}, \frac{y^i}{L} \frac{\partial}{\partial y^i} \right) \quad A = 1, 2$$

$$t_{\theta_1} = L (0, c \theta_1 s \theta_2 s \theta_3, c \theta_1 s \theta_2 c \theta_3, c \theta_1 c \theta_2, -s \theta_1)$$

$$t_{\theta_2} = L (0, s \theta_1 c \theta_2 s \theta_3, s \theta_1 c \theta_2 c \theta_3, -s \theta_1 s \theta_2, 0)$$

$$t_{\theta_3} = L (0, s \theta_1 s \theta_2 c \theta_3, -s \theta_1 s \theta_2 s \theta_3, 0, 0)$$

(Spherical coordinates)

QLE and compactification

Codim-2 spheres M_5

Codim-2 spheres $M_4 \times S_1$

Hence, the **extrinsic curvature tensor** yields $K_{ab}^A = \left(0, -\frac{1}{L} \delta_{\alpha\beta} t_a^\alpha t_b^\beta \right)$

QLE and compactification

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Hence, the **extrinsic curvature tensor** yields $K_{ab}^A = \left(0, -\frac{1}{L} \delta_{\alpha\beta} t_a^\alpha t_b^\beta \right)$

The trace of the extrinsic curvature is given by $K^\alpha = K^A n_A^\alpha = h^{ab} K_{ab}^A n_A^\alpha$

And the integration measure $\sqrt{h} n_\alpha dS$

Induced metric

QLE and compactification

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So we are interested in the integral

$$Q = \int_{\Sigma} \sqrt{h} K^\alpha n_\alpha dS$$

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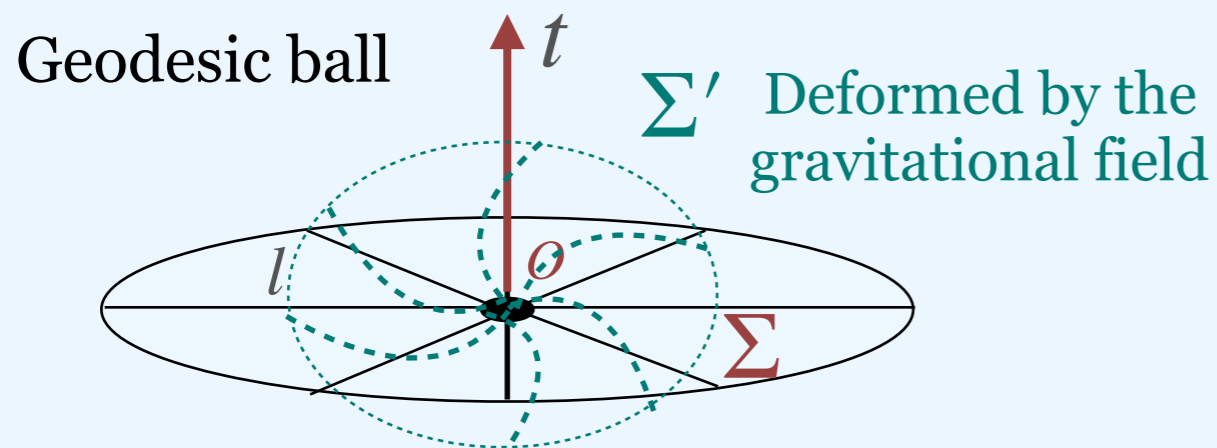
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**Difference between
the two spacetimes**

Quasi-local gravitational energy

We can define the surface as



The extrinsic curvature of the hypersurface is **sensitive** to the strength of the **gravitational field**

The **extrinsic curvature tensor** is given by

$$K_{ab}^A \equiv -t_b^\lambda t_a^\alpha \nabla_\lambda n_\alpha^A$$

t_a^α Tangent vectors $a = 1, \dots, m$

$p \equiv (n - m)$ Codimension of the hypersurface

n_α^A Normal vectors $A = 1, \dots, p$

QLE and compactification

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QLE and compactification

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Codim-2 spheres $M_4 \times S_1$

No restrictions in the range of the variables appearing in the integral

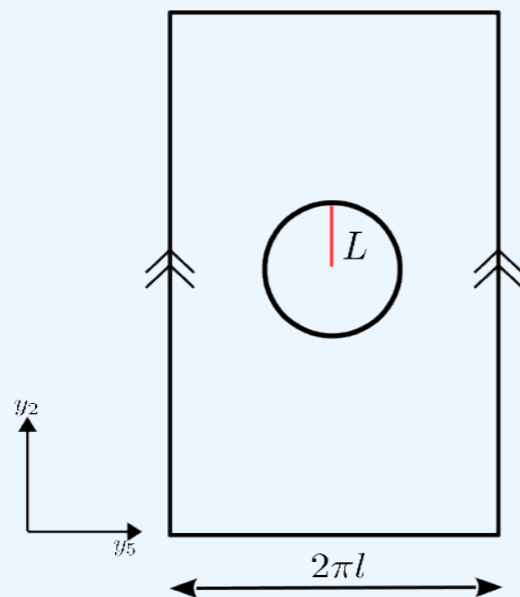
$$Q_{M_5} = \int_{\Sigma} \sqrt{h} K^{\alpha} n_{\alpha} dS = 6\pi^2 L^2$$

QLE and compactification

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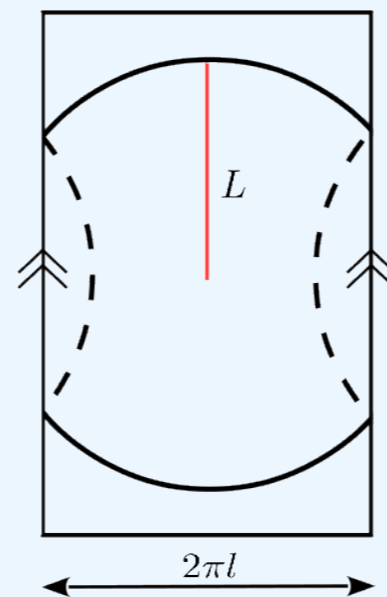
Codim-2 spheres $M_4 \times S_1$

For **small 3-spheres** that completely lie within the compact dimension,

$$-\frac{L}{2} \leq y_5 \leq \frac{L}{2}$$

When $L < 2\pi l$, the **range is restricted**

$$-l\pi \leq y_5 \leq l\pi$$



QLE and compactification

Surfaces of constant time and radius

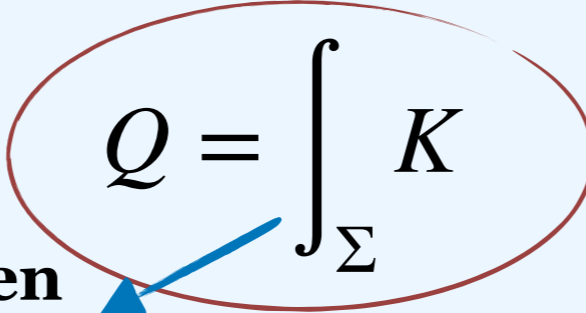
$$y_1 = T, \quad \sum_{i=2}^{i=5} (y_i)^2 \equiv L^2.$$

The two **normal** vectors and the three **tangent** vectors are the same for the two spacetimes

→ Same value for the extrinsic curvature K

But for the QLE the **integration range** is different

Difference between the two spacetimes


$$Q = \int_{\Sigma} K$$

QLE and compactification

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$$Q_{M_4 \times S_1} = 6\pi^2 L^2 \quad \text{for } L \leq l\pi$$

$$Q_{M_4 \times S_1} = 12\pi^2 l \sqrt{L^2 - \pi^2 l^2} + 12\pi L^2 \tan^{-1} \left(\frac{\pi l}{\sqrt{L^2 - \pi^2 l^2}} \right) \quad \text{for } L > l\pi$$