



COSMOLOGICAL EFFECTS OF DM- γ SCATTERING Céline Bœhm, <u>Julia Stadler</u>

DM- γ scattering rate: [Bœhm et al. (2001, 2004)]	
$u_{\gamma \rm DM} = \frac{\sigma_{\gamma \rm DM}}{\sigma_{\rm Th}} \left(\frac{m_{\gamma \rm DM}}{100 {\rm GeV}}\right)^{-1} \Rightarrow \begin{array}{c} \dot{\mu} = a n_{\gamma \rm DM} \sigma_{\gamma \rm DM} \\ \propto u_{\gamma \rm DM} \rho_{\gamma \rm DM} \end{array}$	
Modified Boltzmann Equations:	
$\dot{\theta}_{\gamma \rm DM} = k^2 \psi - \mathcal{H} \theta_{\gamma \rm DM} + c_{\gamma \rm DM}^2 k^2 \delta_{\gamma \rm DM} - \frac{4\rho_{\gamma}}{3\rho_{\gamma \rm DM}} \dot{\mu} \left(\theta_{\gamma \rm DM} - \theta_{\gamma}\right)$	
	$\begin{aligned} DM-\gamma \text{ scattering rate:} & [B \alpha hm et al. (2001, 2004)] \\ u_{\gamma DM} &= \frac{\sigma_{\gamma DM}}{\sigma_{Th}} \left(\frac{m_{\gamma DM}}{100 GeV} \right)^{-1} \Rightarrow & \dot{\mu} &= a n_{\gamma DM} \sigma_{\gamma DM} \\ & \propto u_{\gamma DM} \rho_{\gamma DM} \end{aligned}$ $\begin{aligned} \mathbf{Modified Boltzmann Equations:} \\ \dot{\theta}_{\gamma DM} &= k^2 \psi - \mathcal{H} \theta_{\gamma DM} + c_{\gamma DM}^2 k^2 \delta_{\gamma DM} - \frac{4\rho_{\gamma}}{3\rho_{\gamma DM}} \dot{\mu} \left(\theta_{\gamma DM} - \theta_{\gamma} \right) \end{aligned}$



 $\dot{\theta}_{\gamma} = k^2 \psi + k^2 \left(\frac{1}{4}\delta_{\gamma} - \sigma_{\gamma}\right) + \dot{\kappa} \left(\theta_b - \theta_{\gamma}\right) + \dot{\mu} \left(\theta_{\gamma \text{DM}} - \theta_{\gamma}\right)$

+ for higher order photon multipoles : $\dot{\kappa} \rightarrow (\dot{\kappa} + \dot{\mu})$

Improvements in this work

- Accurate treatment of the tight coupling regime. 0
- Parameter constraints including Planck polarization data. 0
- Inclusion of the dark matter sound speed. 0
 - Extension to multicomponent DM ("mixed DM"). 0

Effects of γ -DM scattering on the CMB spectrum



Constraints from Planck (arXiv[1802.0658])

We analyse the Planck 2015 temperature & polarization data to obtain limits on $u_{\gamma DM}$:

- Computation of CMB spectra with CLASS Ο
- MCMC sampling with MontePython Ο
- Parameter Space: Ο

[Blas, Lesgourgues, Tram (2011)]

- [Audren, Lesgourgues et al. (2012)]
- $H_0 \mid \Omega_b h^2 \mid \Omega_{\mathrm{DM}} h^2 \mid \ln \left(10^{10} A_s \right) \mid n_s \mid \tau_{\mathrm{reio}} \mid u_{\gamma \mathrm{DM}}$

"Planck TT + lowTEB": $\sigma_{\gamma \text{DM}} < 1.5 \times (m_{\gamma \text{DM}}/\text{GeV}) \text{ fm}^2$

DM sound speed (arXiv[1802.0658])

At early times DM- γ scattering maintains kinetic equilibrium between the two species. Consequently, the DM temperature is larger, and the sound speed can become important.



size of the sound speed. $u_{\gamma \text{DM}}$ sets the time of decoupling and

For $m_{\gamma \rm DM} \gtrsim 10 \, \rm keV$ the DM



Mixed Dark Matter (arXiv[1807.10034])

There are two components of dark matter: on scatters elastically off photons (γ DM), the other one is collisionles. The fraction of interacting DM is $f_{\gamma \text{DM}} = \Omega_{\gamma \text{DM}} / (\Omega_{\gamma \text{DM}} + \Omega_{\text{CDM}})$.





Background Image: ESA and the Planck Collaboration[©], Star formation and magnetic turbulence in the Orion Molecular Cloud