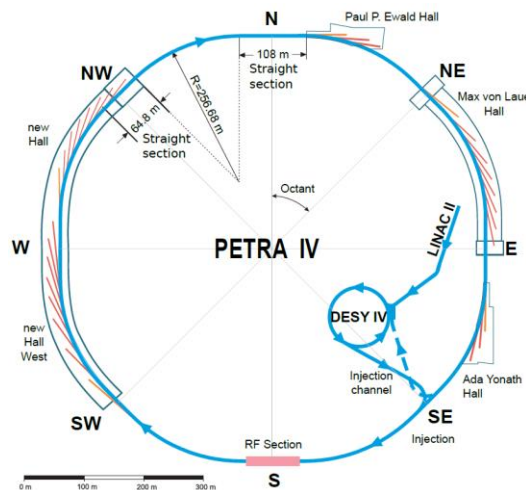
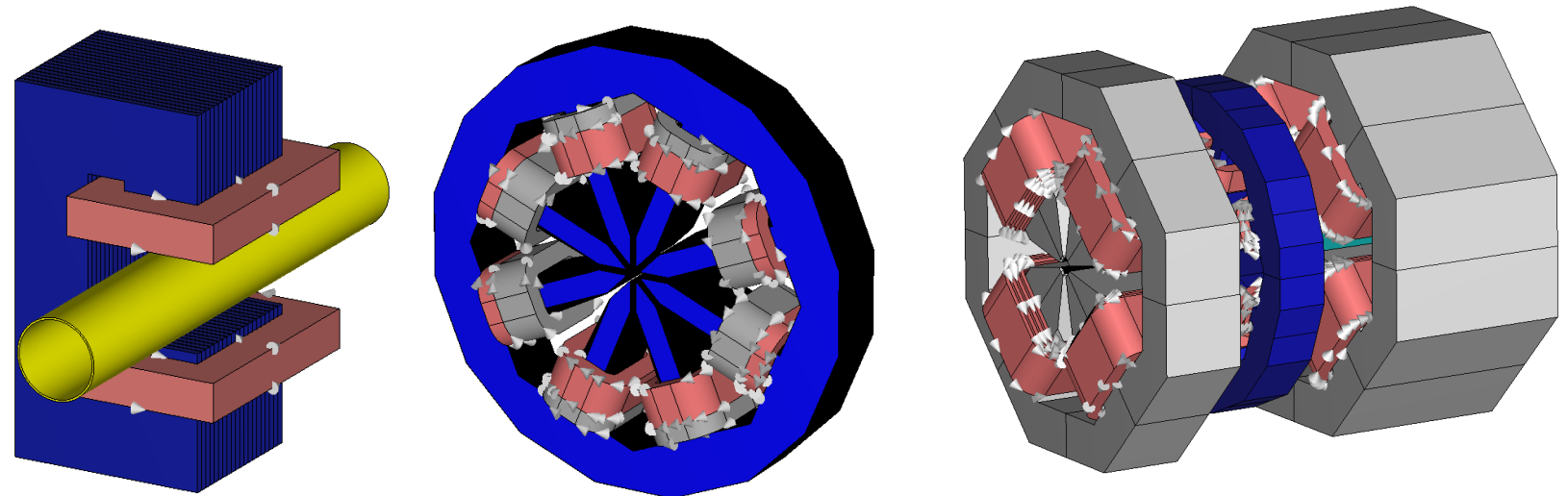


# FINITE ELEMENT SIMULATION OF FAST CORRECTOR MAGNETS FOR PETRA IV

Jan-Magnus Christmann<sup>1</sup>, Moritz von Tresckow<sup>1</sup>, Herbert De Gersem<sup>1</sup>,  
Alexander Aloev<sup>2</sup>, Sajjad H. Mirza<sup>2</sup>, Sven Pfeiffer<sup>2</sup>, and Holger Schlarb<sup>2</sup>



<sup>1</sup>TEMF, TU Darmstadt, Germany  
<sup>2</sup>DESY, Hamburg, Germany



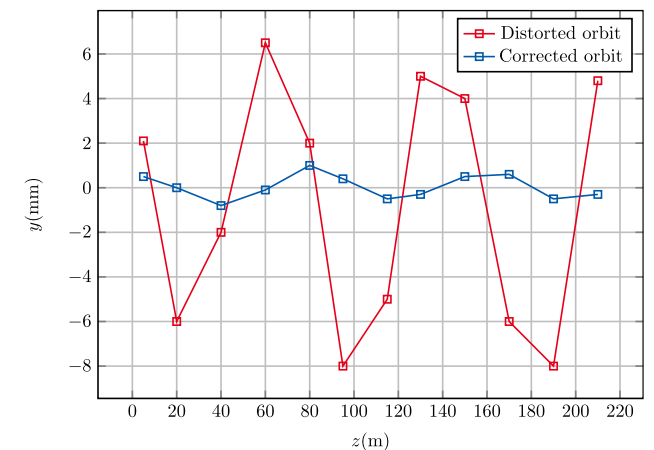
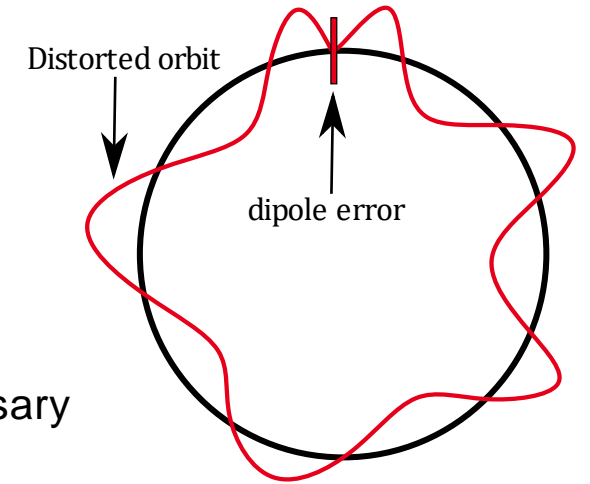
PETRA IV Conceptual Design Report

# CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring Quadrupoles
- 6** Conclusion/Outlook

# INTRODUCTION

- Circular accelerators need dipole magnets to correct orbit distortions
- **PETRA IV**: ultra-low emittance synchrotron radiation source
- ➔ fast orbit feedback system, **corrector magnets with frequencies in kHz range** necessary
- **Strong eddy currents** ➔ power losses, time delay, and field distortion
- **Simulation challenging** due to small skin depths and laminated yoke
- ➔ **Need for technique to simplify simulations**



# CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring Quadrupoles
- 6** Conclusion/Outlook

# HOMOGENIZATION TECHNIQUE

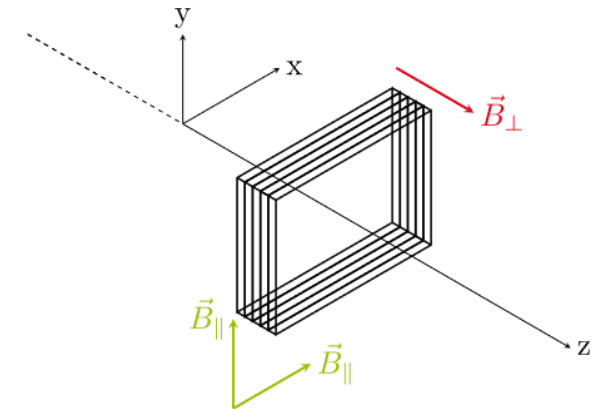
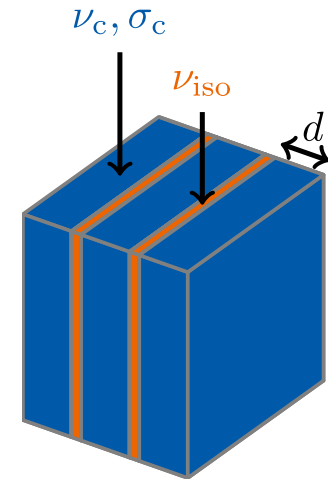
- Magnetoquasistatic PDE:  $\nabla \times (\nu(\vec{r}) \nabla \times \vec{A}(\vec{r})) + j\omega\sigma(\vec{r})\vec{A}(\vec{r}) = \vec{J}_s(\vec{r})$
- Replace reluctivity  $\nu(\vec{r})$  and conductivity  $\sigma(\vec{r})$  in the laminated yoke with spatially constant tensors

$$\nu(\vec{r}) \rightarrow \underline{\bar{\nu}} = \frac{1}{8} \sigma_c d \delta \omega (1 + j) \frac{\sinh((1 + j)\delta^{-1}d)}{\sinh^2((1 + j)\delta^{-1}d/2)} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \nu_c \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\sigma(\vec{r}) \rightarrow \underline{\bar{\sigma}} = \gamma \sigma_c \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Skin depth  $\delta = \sqrt{2/\omega\sigma_c\mu_c}$   
Stacking factor  $\gamma = \frac{V_c}{V_{\text{Yoke}}}$

P. Dular et al., 2003  
L. Krähenbühl et al., 2004  
H. De Gersem et al., 2012

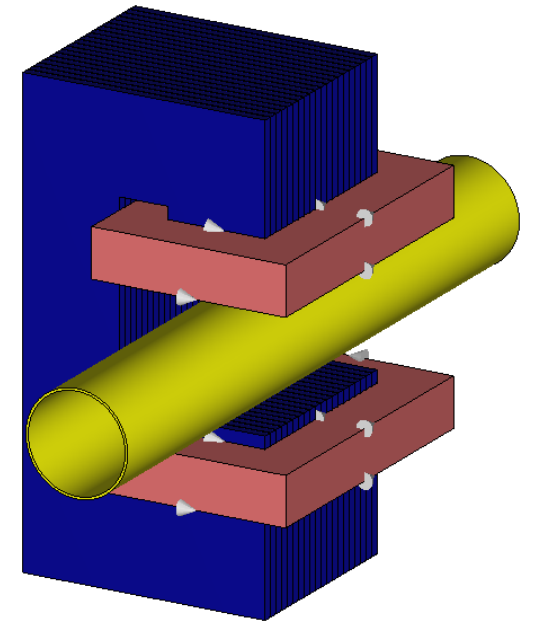
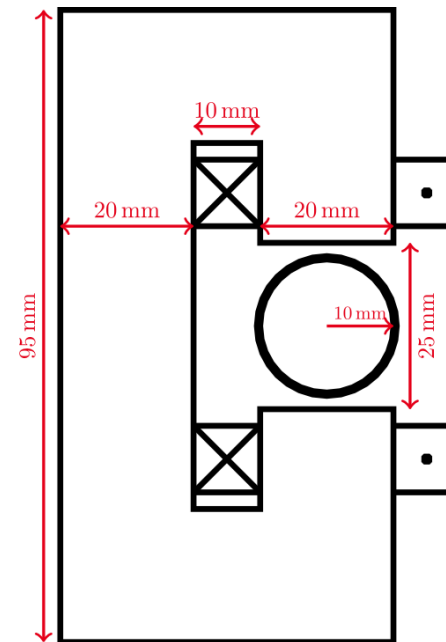


# CONTENTS

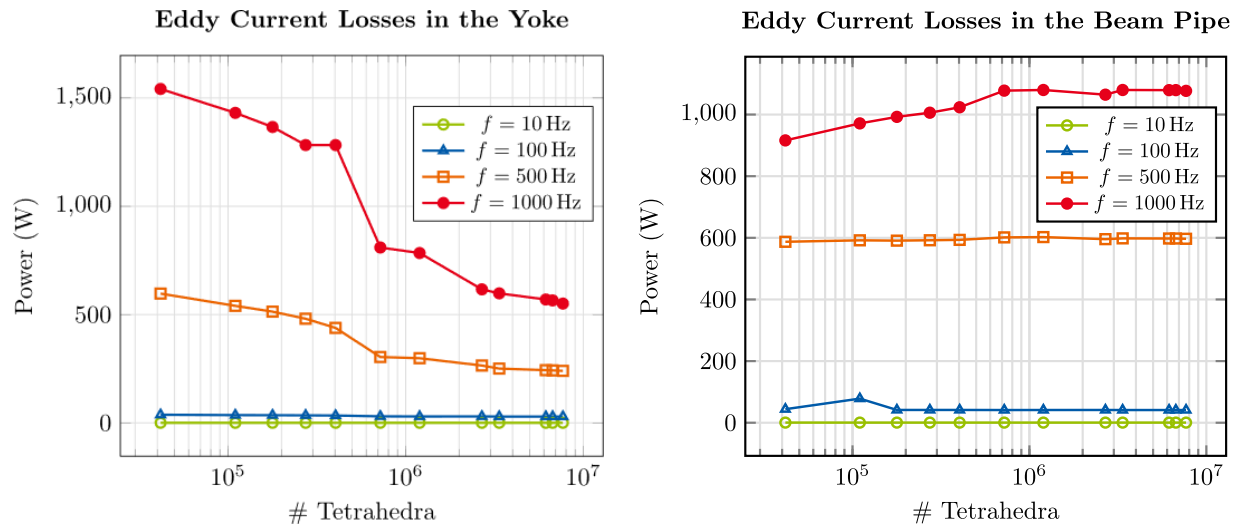
- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring Quadrupoles
- 6** Conclusion/Outlook

# MODEL DESCRIPTION

- **Iron yoke:** length = 40 mm, lamination thickness = 1.83 mm
- **Copper beam pipe:** thickness = 0.5 mm, length = 140 mm
- **Coils:** current = 10 A (peak), # turns = 250
- **Frequency domain simulation via CST Studio Suite®**



# SIMULATION OF THE FULL MODEL



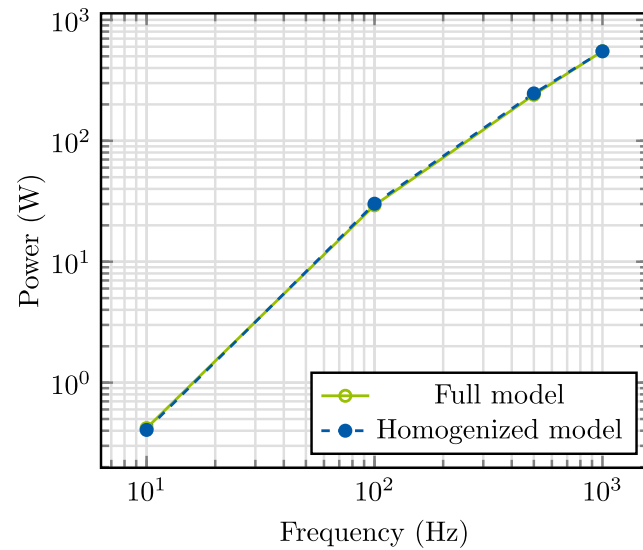
- Strong **mesh dependence of power losses** at higher frequencies
  - ➔ Obtaining reliable results is difficult
  - ➔ **Need for simplified model**

<b># Tetrahedra</b>	$4.2 \cdot 10^4$	$4.0 \cdot 10^5$	$1.2 \cdot 10^6$	$3.4 \cdot 10^6$	$7.7 \cdot 10^6$
<b>Simulation time</b>	2 min	20 min	1 h	7.5 h	21.5 h

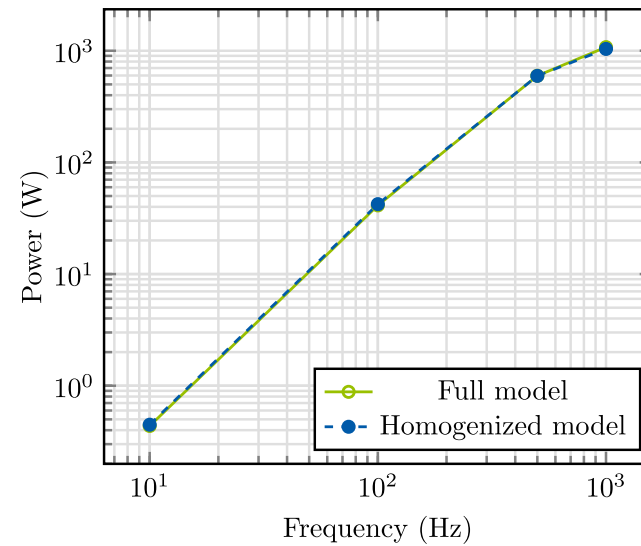


# HOMOGENIZED VS. FULL MODEL

Eddy Current Losses in the Yoke

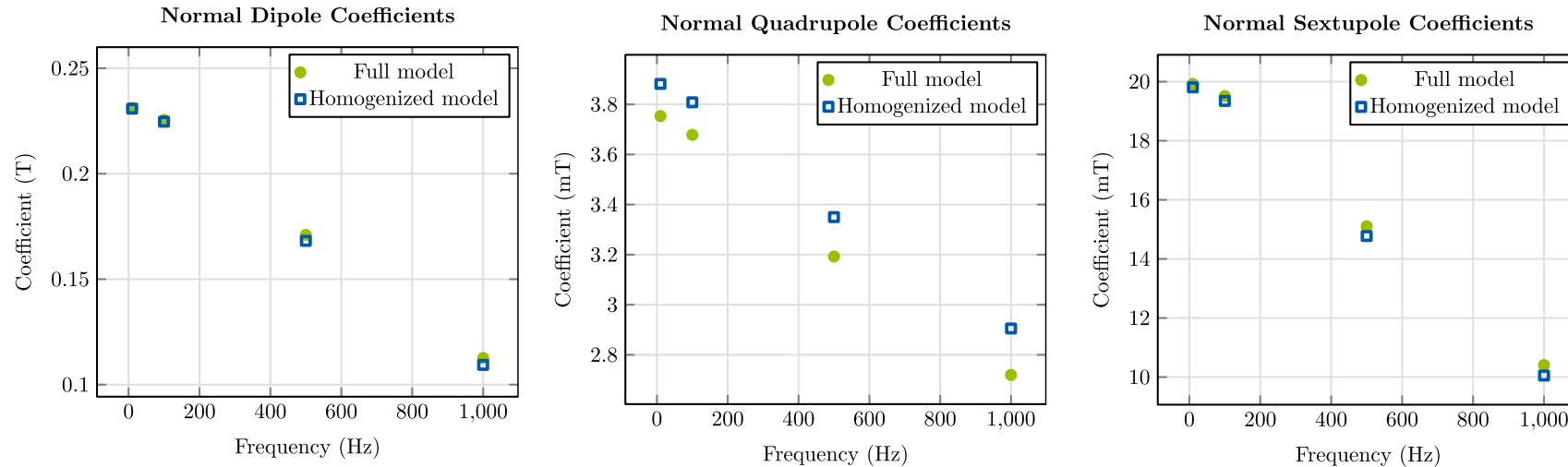


Eddy Current Losses in the Beam Pipe



- **Good approximation** of losses in yoke & beam pipe (max. relative error 4 %)
- **Simulation time reduced** from several hours to 4 min

# HOMOGENIZED VS. FULL MODEL



- Homogenization technique yields accurate approximation of multipole coefficients  
 → Aperture field accurately represented

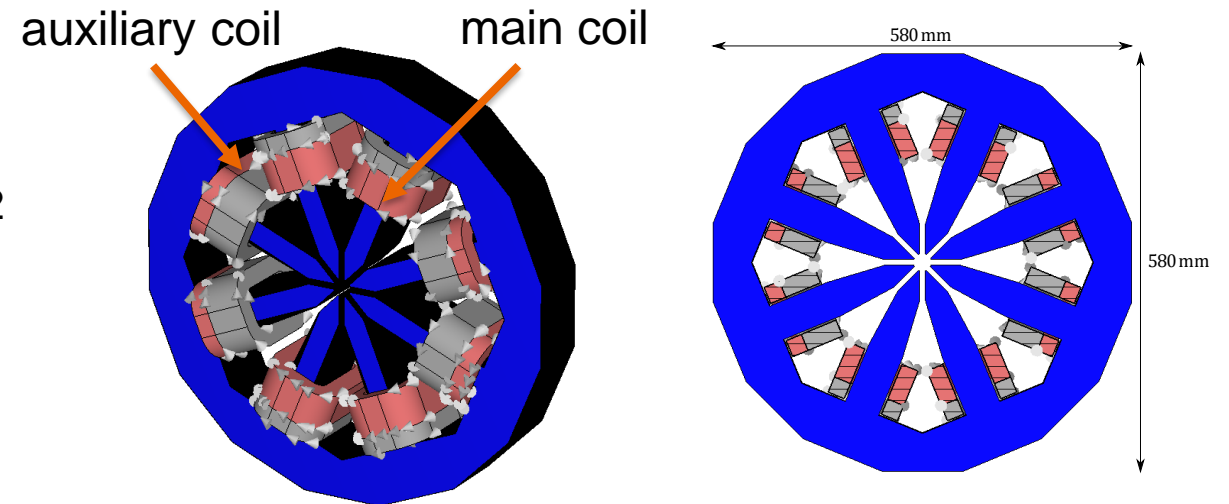
Multipole coefficient	Average rel. error
Dipole	1 %
Quadrupole	5 %
Sextupole	2 %

# CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring  
Quadrupoles
- 6** Conclusion/Outlook

# MODEL DESCRIPTION

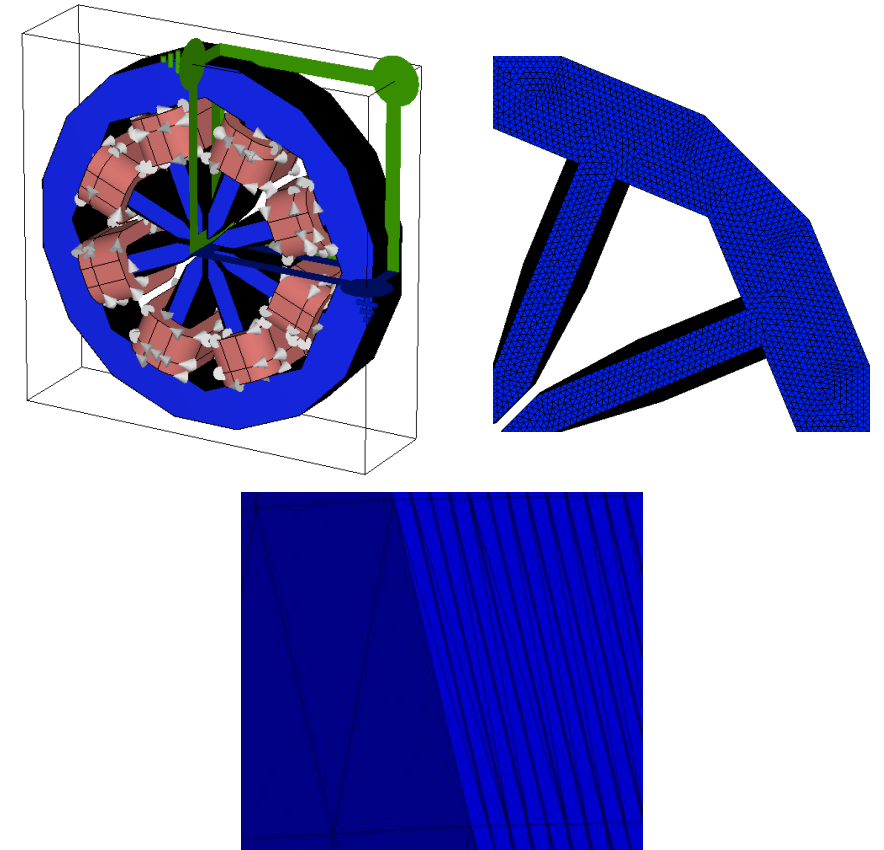
- **Dipole** corrector with **octupole-like design**
- **Coils:**
  - 4 main coils: current = 27.4 A (peak), # turns = 53
  - 4 auxiliary coils: current = 27.4 A (peak), # turns = 22
- **Iron yoke:**
  - Diameter = 580 mm, length = 90 mm
  - Lamination thickness = 0.5 mm
- At first **no beam pipe**



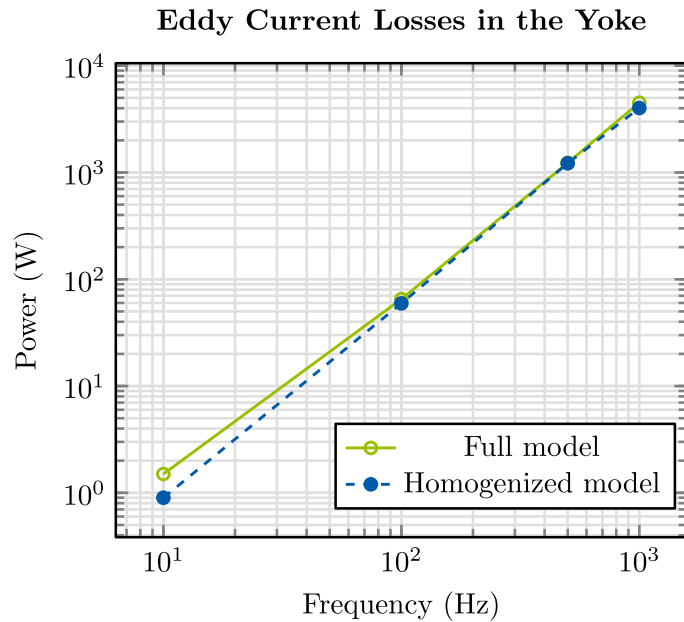
Design by A. Aloev (DESY),  
inspired by APS

# SIMULATION OF THE FULL MODEL

- **Frequency domain simulation** via **CST Studio Suite<sup>®</sup>**
- Three symmetry planes, test frequencies  $f = 10 \text{ Hz}, 100 \text{ Hz}, 500 \text{ Hz}, 1 \text{ kHz}$
- Long simulation times even for relatively coarse meshes
- Finest mesh: # tetrahedra =  $2.3 \cdot 10^6$  → **simulation time = 26 h**
- Skin depth cannot be resolved → **power loss still mesh-dependent**



# HOMOGENIZED VS. FULL MODEL

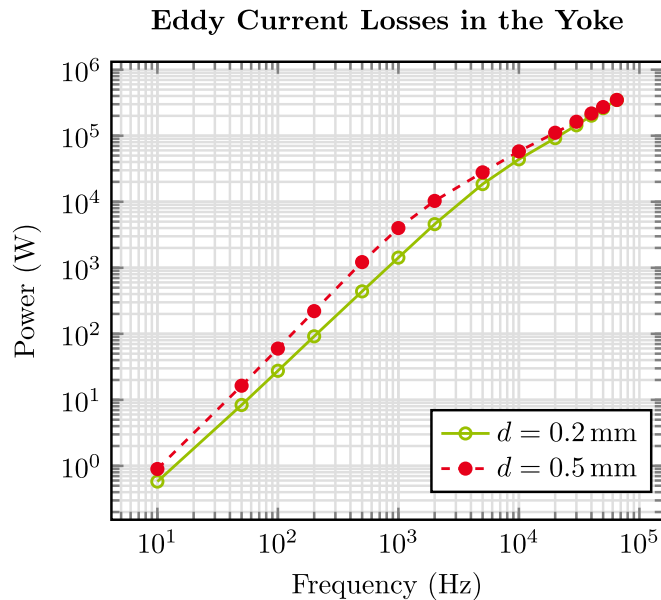


Multipole coefficient	Average rel. deviation
Dipole	1 %
14-pole	1 %
18-pole	3 %

- Similar power losses
  - Good agreement in multipole coefficients
  - Simulation time reduces from 26 h to 5 min
- ➔ Homogenized model can be used for further studies

Keep in mind:  
Power losses in full model  
are still mesh-dependent !

# LOSSES FOR DIFFERENT LAMINATION THICKNESSES

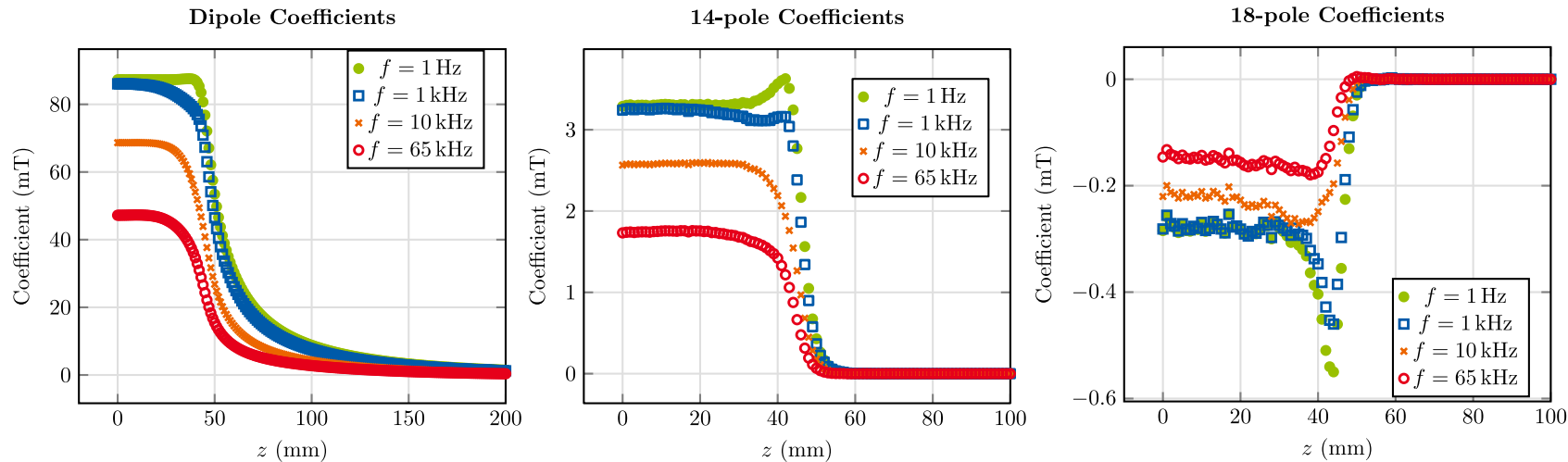


$f(\text{Hz})$	Eddy current losses (W)			
	$d = 0.2 \text{ mm}$	$d = 0.3 \text{ mm}$	$d = 0.4 \text{ mm}$	$d = 0.5 \text{ mm}$
10	$5.8 \cdot 10^{-1}$	$6.5 \cdot 10^{-1}$	$7.6 \cdot 10^{-1}$	$9.0 \cdot 10^{-1}$
100	$2.8 \cdot 10^1$	$3.4 \cdot 10^1$	$4.6 \cdot 10^1$	$6.0 \cdot 10^1$
500	$4.4 \cdot 10^2$	$6.2 \cdot 10^2$	$9.0 \cdot 10^2$	$1.2 \cdot 10^3$
1000	$1.4 \cdot 10^3$	$2.1 \cdot 10^3$	$3.1 \cdot 10^3$	$4.0 \cdot 10^3$
10000	$4.4 \cdot 10^4$	$4.9 \cdot 10^4$	$5.5 \cdot 10^4$	$5.8 \cdot 10^4$
30000	$1.4 \cdot 10^5$	$1.6 \cdot 10^5$	$1.6 \cdot 10^5$	$1.6 \cdot 10^5$
65000	$3.5 \cdot 10^5$	$3.6 \cdot 10^5$	$3.6 \cdot 10^5$	$3.5 \cdot 10^5$

Simulation uses the same current for all frequencies !

- Use homogenization to investigate losses up to 65 kHz
- Vary  $d = 0.2 - 0.5 \text{ mm}$ , keep  $\gamma \approx 0.91$  constant
- At **low frequencies**, the **lamination thickness has strong influence** on the losses
- At **very high frequencies**, the **lamination thickness has no influence** on the losses

# LONGITUDINAL MULTIPOLE DISTRIBUTION

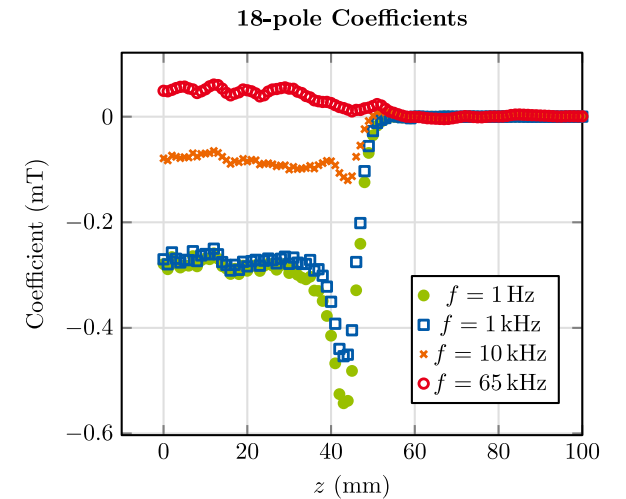
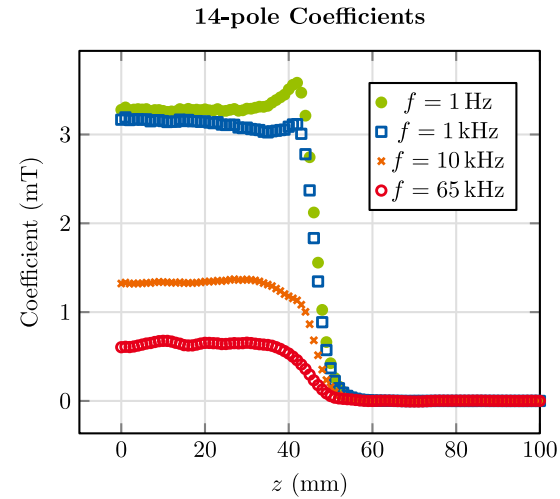
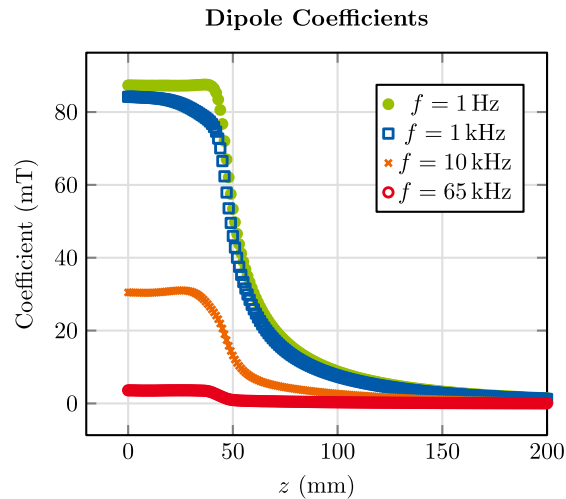
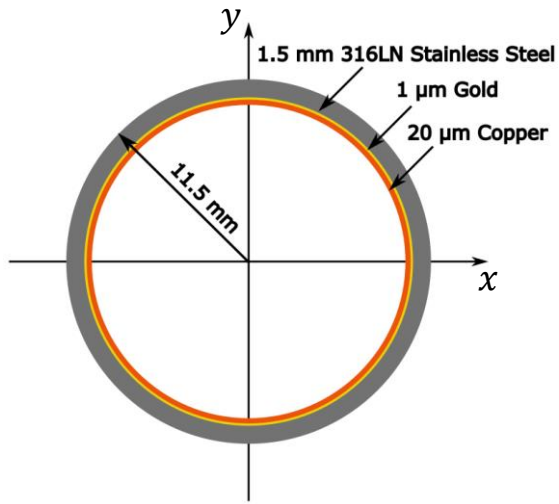


- Updated # turns & current:
  - ➔ Main coils: 65 turns, 15 A
  - ➔ Aux. coils: 27 turns, 15 A
- 65 kHz vs. 1 Hz:
  - ➔ Int. dipoles: -57 %
  - ➔ Int. 14-poles: -52 %
  - ➔ Int. 18-poles: -54 %

$f$ (Hz)	Int. dipole (mT m)	Int. 14-pole ( $\mu$ T m)	Int. 18-pole ( $\mu$ T m)
1	11.6	316.4	-30.3
1000	10.7	300.4	-28.6
10000	7.6	229.0	-21.5
65000	5.0	150.3	-13.9



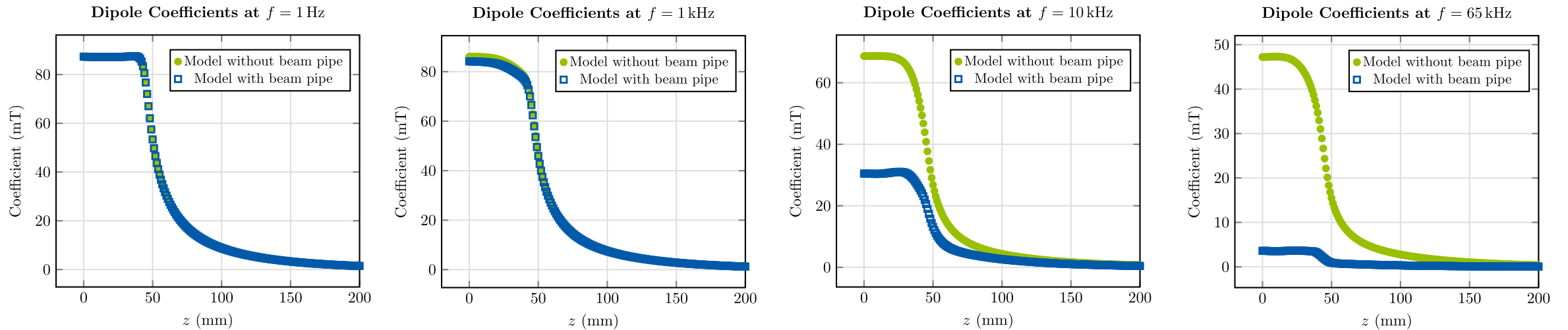
# INCLUSION OF BEAM PIPE



$f$ (Hz)	Int. dipole (mT m)	Int. 14-pole ( $\mu\text{T m}$ )	Int. 18-pole ( $\mu\text{T m}$ )
1	11.5	313.3	-30.6
1000	10.5	292.7	-28.1
10000	3.6	122.6	-8.3
65000	0.4	57.4	4.3

- General shape similar to model without beam pipe
- 65 kHz vs. 1 Hz:
  - ➔ Int. dipoles: -97 % (-57 %)
  - ➔ Int. 14-poles: -82 % (-52 %)
  - ➔ Int. 18-poles change sign (-54 %)

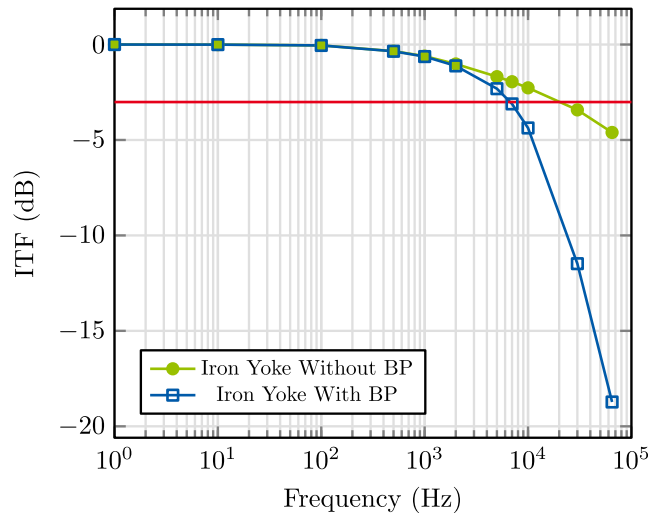
# INCLUSION OF BEAM PIPE



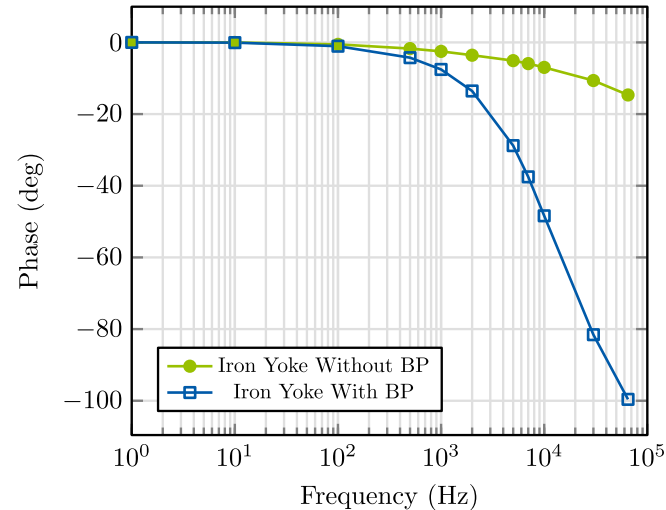
- **Up to  $f \approx 1$  kHz:** Only minor differences between the two models
- **For  $f \gg 1$  kHz:** Strong attenuation of dipole field due to eddy currents in beam pipe
- At higher frequencies, beam pipe leads to greater effective length of the magnet

# INTEGRATED TRANSFER FUNCTION AND FIELD LAG

Integrated Transfer Function



Field Lag w.r.t. Current



$$ITF(f) = \frac{\int_l B_1(z, f) dz}{\int_l B_1(z, f = 1\text{Hz}) dz}$$

Yoke material	3 dB bandwidth	Phase shift at bandwidth
Iron	7 kHz	38°
M-19 Steel	10 kHz	46°
1010 Steel	7 kHz	38°

Yoke material	Average relative permeability*	Conductivity (MS/m)
Iron	5690	10.4
M-19 Steel	4166	1.9
1010 Steel	2780	6.993

- **Beam pipe is made out of 316 LN SS** ( $\sigma = 1.351 \cdot 10^6$ ,  $\mu_r = 1.01$ ) and has an **outer radius of 11 mm** and a **thickness of 1 mm**

\* Values are computed from results of static simulations with non-linear BH-curve

# CONTENTS

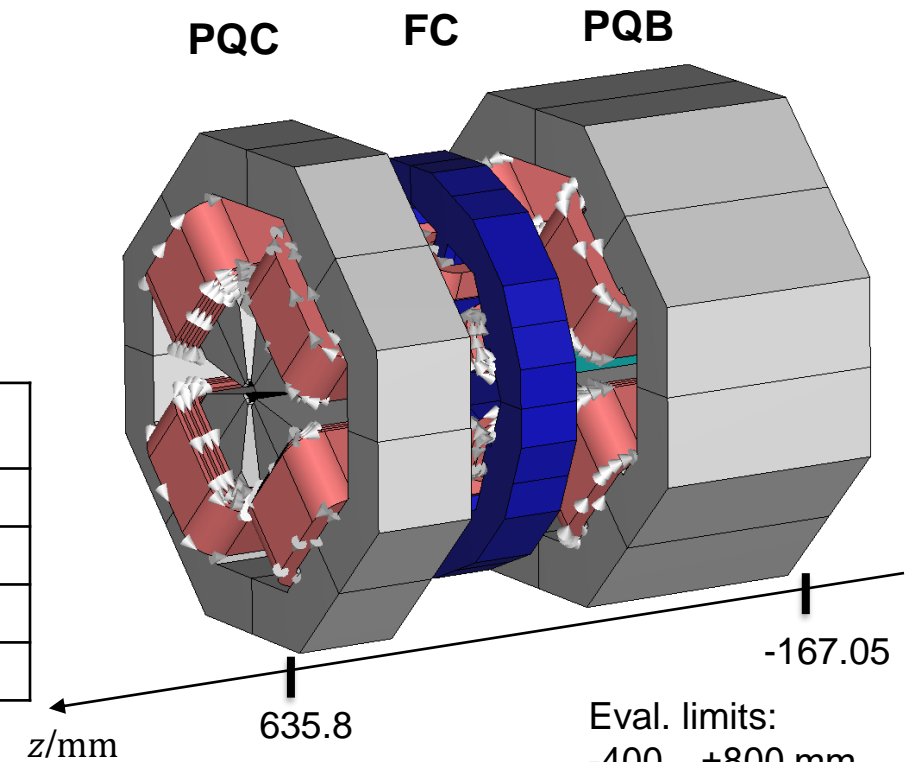
- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring  
Quadrupoles
- 6** Conclusion/Outlook

# MODEL DESCRIPTION

- Corrector magnet (FC) with two neighboring quadrupole magnets (PQB & PQC)
- AC currents in corrector coils, DC currents in quadrupole coils
- All yokes are 1010 steel, PQB quadrupoles have Vacoflux-50 poles
- Quadrupole yokes are solid, corrector yoke is laminated
- Beam pipe made out of 316LN SS with outer radius of 11 mm and thickness of 1 mm
- Distance between corrector yoke and quadrupole yokes ~ 11.5 cm

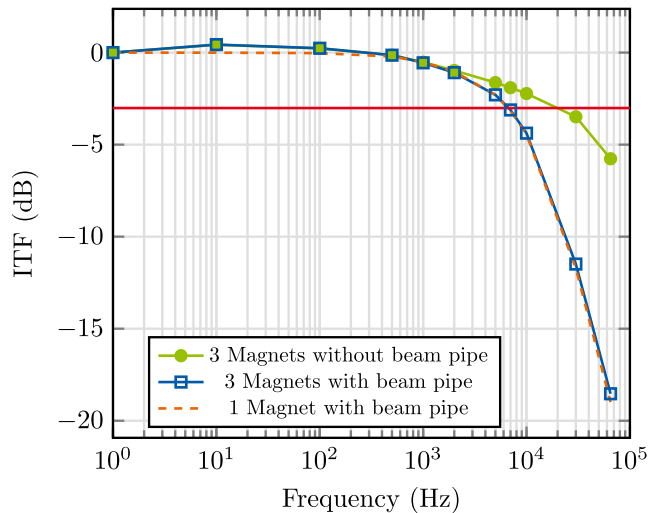
Material	Average Relative Permeability*	Conductivity (MS/m)	Coils	Ampere turns
1010 Steel (PQC)	1450	6.993	PQB	5728.1 At
1010 Steel (PQB)	1810	6.993	FC (main)	975 At
Vacoflux-50 (PQB)	5000	2.38	FC (aux.)	405 At
1010 Steel (FC)	2780	6.993	PQC	5659.5 At

\* Values are computed from results of static simulations with non-linear BH-curve

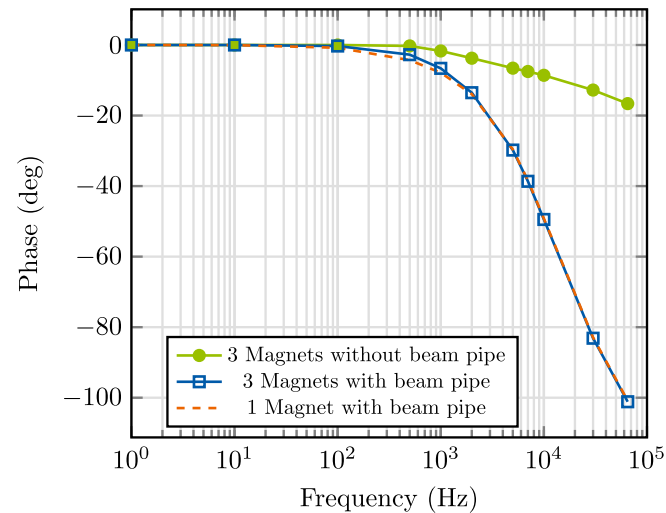


# INTEGRATED TRANSFER FUNCTION AND FIELD LAG

Integrated Transfer Function



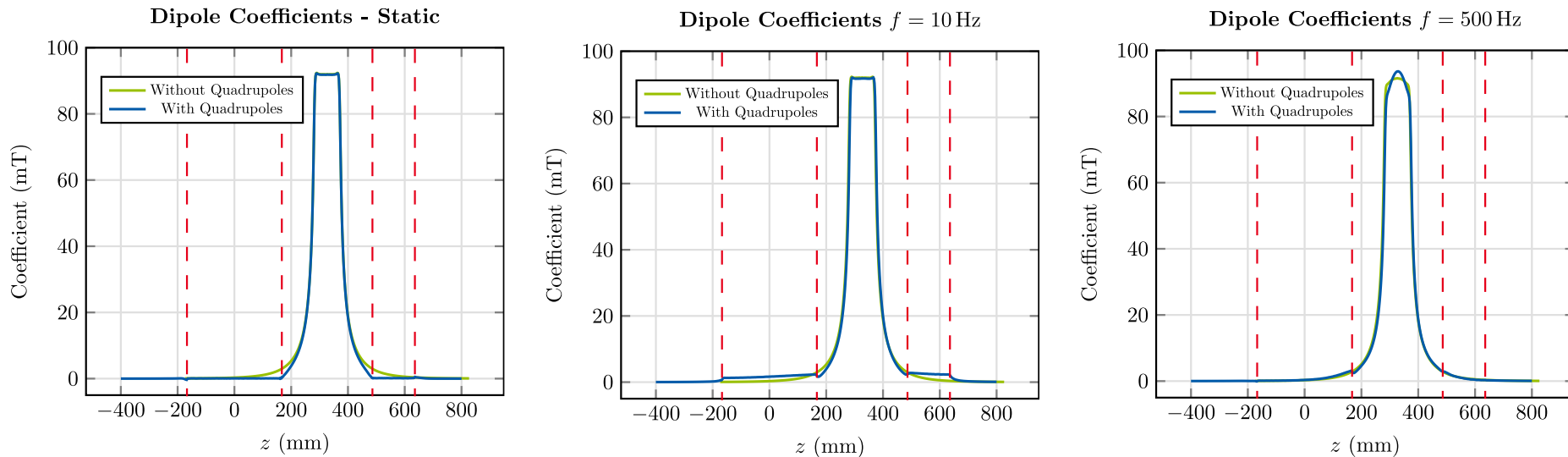
Field Lag w.r.t. Current



	Model without beam pipe	Model with beam pipe
<b>3 dB bandwidth</b>	20 kHz	7 kHz
<b>Phase shift at bandwidth</b>	11°	39°

- Very similar results as for the model without neighboring quadrupoles
- **Main difference:** at low frequencies, a  $\sim 0.5$  dB peak is occurring in the ITF of the model with the neighboring quadrupoles

# DIPOLE COEFFICIENTS ALONG THE AXIS



Simulation with  
beam pipe

- **At low frequencies ( $f \leq 100$  Hz), we observe a **parasitic dipole** component **inside the quadrupole magnets****
- This dipole component is due to eddy currents induced in the quadrupole yokes by the AC corrector field
- ➔ Peak in ITF at low frequencies
- ➔ Shift of the center of mass (  $\sim 0.5$  cm at most)

# CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-alone Corrector Magnet
- 5** Corrector Magnet with Neighboring  
Quadrupoles

- 6** Conclusion/Outlook



# CONCLUSION/OUTLOOK

- Validation of homogenization technique using **toy model**
  - **Good approximation** of multipoles and power losses
  - Simulation time reduced from **several hours to a few minutes**

- Application to **corrector magnet model**

- Power losses for different lamination thicknesses
- Longitudinal multipole distributions
- Integrated transfer function and field lag
- Cross-talk with neighboring magnets



Homogenization enables us to study this over the frequency range of interest from **DC up to 65 kHz**

- **Ongoing investigations:**

- Simulations with different variations of the beam pipe and cooling channels
- Approximate treatment of non-linear material properties

# REFERENCES

- [1] PETRA IV Conceptual Design Report.
- [2] K. Wille, *Physik der Teilchenbeschleuniger und Synchrotronstrahlungsquellen*. Stuttgart, Germany: Teubner, 1992.
- [3] P. Dular et al., “A 3-D Magnetic Vector Potential Formulation Taking Eddy Currents in Lamination Stacks Into Account,” *IEEE Trans. Magn.*, vol. 39, no. 3, pp. 1424-1427, May 2003.
- [4] L. Krähenbühl et al., “Homogenization of Lamination Stacks in Linear Magnetodynamics,” *IEEE Trans. Magn.*, vol. 40, no. 2, pp. 912 - 915 Mar. 2004.
- [5] H. De Gersem, S. Vanaverbeke, and G. Samaey, “Three-Dimensional-Two-Dimensional Coupled Model for Eddy Currents in Laminated Iron Cores,” *IEEE. Trans. Magn.*, vol. 48, no. 2, pp.815 – 818, Feb. 2012.