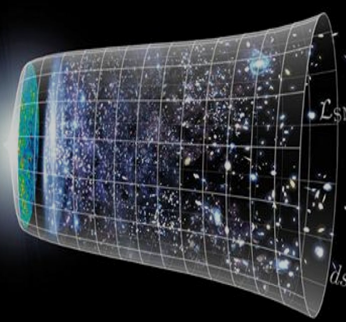


Searches for T violation in neutrino oscillations with T2HK and DUNE



$$\mathcal{L}_{\text{SM}} = -\frac{1}{4}\text{Tr}(F^{\mu\nu}F_{\mu\nu}) + \frac{\theta}{64\pi^2}\text{Tr}(G^{\mu\nu}\tilde{G}_{\mu\nu}) + |D_\mu\phi|^2 + \mu^2\phi^\dagger\phi - \lambda(\phi^\dagger\phi)^2$$

$$+ i\bar{\psi}_L D^\mu\gamma_\mu\psi_L + i\bar{\psi}_R D^\mu\gamma_\mu\psi_R - \left(\lambda_{ij}^d\bar{\psi}_{iL}\phi\psi_{jR} + \lambda_{ij}^u\bar{\psi}_{iL}\tilde{\phi}\psi_{jR}\right) + \text{h.c.}$$

$$\mathcal{L}_{\text{dim=5}} = \frac{Y_{\alpha\beta}}{\Lambda_{\text{LNV}}}\left(\bar{L}_\alpha^c\tilde{\phi}^*\right)\left(\tilde{\phi}^\dagger L_\beta\right)$$

$$ds^2 = c^2dt^2 - a(t)^2\left(\frac{dr^2}{1-kr^2} + r^2d\Omega^2\right)$$



**Sabya Sachi Chatterjee, Sudhanwa Patra, Thomas Schwetz,
Kiran Sharma***

[arXiv:2408.06419](https://arxiv.org/abs/2408.06419)

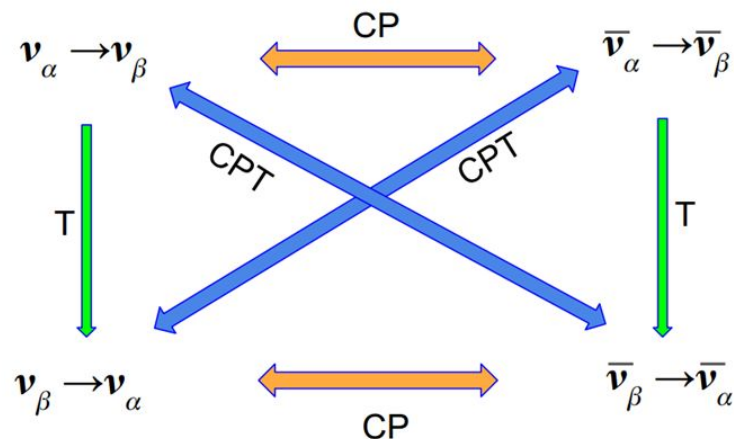
Kiran Sharma, ISAPP, KIT / Bad Liebenzell, 16–27 Sept 2024

Outline

- **CP and T violation**
- **Model Independent Approach**
- **T violation test with New Physics**

CP and T violation in neutrino oscillations

- In two flavor picture, no fundamental CP and T violation.
- In three flavor picture, fundamental CP violation due to leptonic phase.
- Normal matter (no. of particles is not equal to no. of antiparticles) induces extrinsic(fake) CP violation.
- Matter with symmetric density profiles don't induce fake T violation, asymmetric profiles do.



$$\Delta P_{\alpha\beta}^{CP} = P_{\alpha\beta} - \bar{P}_{\alpha\beta},$$

$$\Delta P_{\alpha\beta}^T = P_{\alpha\beta} - P_{\beta\alpha},$$

$$\Delta P_{\alpha\beta}^{CPT} = P_{\alpha\beta} - \bar{P}_{\beta\alpha}.$$

Model Dependent Approach in SM

- Standard approach is to perform a model dependent fit to data (combining accelerator and reactor data).
- Not possible to construct model-independent CP asymmetric observables.
- CPV signature in neutrino oscillation experiments is rather indirect.
- Observation of CPV is equivalent to establishing δ different from 0 and π at a certain confidence level.
- T violation is rather hard by exchanging the source and detector.

Model Independent Test for T violation

- Standard Approach is to perform a **model-dependent fit** to data, which assumes only 3 standard model neutrinos exist.
- However, new physics scenarios like non-unitary mixings, non-standard neutrino interactions, presence of sterile neutrinos act as **New Source of CP and T violation**.
- Develop a largely **model-independent test** covering a wide class of non-standard scenarios.

Schwetz, Segarra, On T violation in non-standard neutrino oscillation scenarios [[2112.08801](#)]

Schwetz, Segarra, Model-independent test of T violation in neutrino oscillations [[2106.16099](#)]

Assumptions of Model

- The evolution of the flavor state is described :

$$i\partial_t|\psi\rangle = H(E_\nu)|\psi\rangle.$$

- Allow for non-unitary mixing among energy eigenstates and flavor states at detection and production:

$$|\nu_\alpha^{s,d}\rangle = \sum_{i=1}^3 (N_{\alpha i}^{s,d})^* |\nu_i\rangle$$

- Medium effects are defined by constant matter density approximately.
- The eigenvalues and their energy dependence resembles approximately the one following from the effective neutrino mass squared differences in matter in the SM.

Appearance Probability

- The appearance probability is defined as :

$$P = \left| \sum_{i=1}^3 c_i e^{-i\lambda_i L} \right|^2, \quad c_i \equiv N_{\mu i}^{s*} N_{ei}^d,$$

- Expanding it out, in terms of new variable, ϵ that describes deviation from unitarity and leads to a “zero distance effect”.

$$P = \left| c_2(e^{-i(\lambda_2-\lambda_1)L} - 1) + c_3(e^{-i(\lambda_3-\lambda_1)L} - 1) + \epsilon \right|^2,$$

$$\epsilon \equiv \sum_{i=1}^3 c_i.$$

$$P^{\text{ND}} \equiv P(L \rightarrow 0) = |\epsilon|^2.$$

T violation Test for 2 experiments

$$P_{\text{even}} = \gamma_2 c_2 (c_2 - \epsilon) + \gamma_3 c_3 (c_3 - \epsilon) + \gamma_{23} c_2 c_3 + \epsilon^2$$

With the abbreviations

$$\begin{aligned} \gamma_i &= 4 \sin^2 \phi_{i1} \quad (i = 2, 3), \\ \gamma_{23} &= 8 \sin \phi_{21} \sin \phi_{31} \cos(\phi_{31} - \phi_{21}). \end{aligned} \quad \left\{ \begin{array}{l} \phi_{ij} \approx \frac{\Delta m_{ij,\text{eff}}^2(E_\nu)L}{2E_\nu} \end{array} \right.$$

- There is always a fit for two experiments plus a near detector, which provide three data points.
- Under certain conditions the quadratic nature of the parameter dependence does not provide a solution for three data points.

Define a model-independent observable X_T , built out of the observed probabilities $P_{\nu_\mu \rightarrow \nu_e}(L)$ at two baselines L_1, L_2 and at a near detector.

$$X_T \equiv P_{\text{even}}(L_2) - P_{\text{even}}(L_1) - \epsilon^2 \delta_0 = \delta_2 c_2^2 + \delta_3 c_3^2 + \delta_{23} c_2 c_3$$

With,

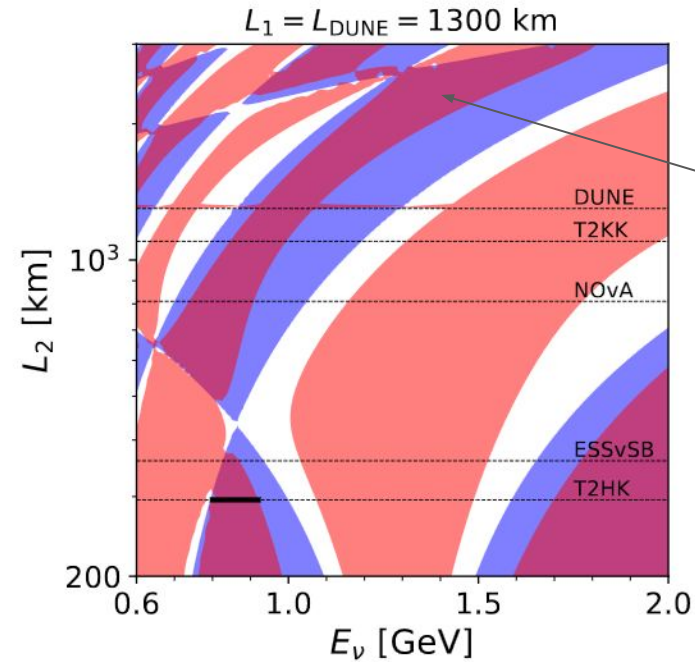
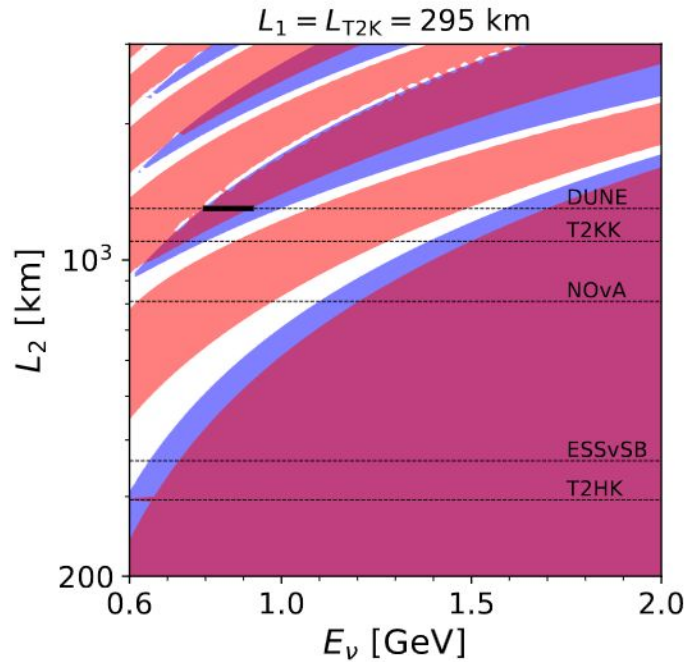
$$\delta_0 = \frac{\delta_2 + \delta_3 - \delta_{23}}{\delta_{23}^2 / (\delta_2 \delta_3) - 4}.$$

The right-hand side of eq. is a non-negative function of c_2 and c_3 if

$$\begin{aligned} \delta_3 > 0 \quad \text{and} \quad \delta_2 > 0, \quad \text{and} \\ |\alpha| < 2 \quad \text{with} \quad \alpha \equiv \frac{\delta_{23}}{\sqrt{\delta_2 \delta_3}}. \end{aligned}$$

$$\chi_T^{\text{obs}} = P_{\nu_\mu \rightarrow \nu_e}^{\text{obs}}(L_2) - P_{\nu_\mu \rightarrow \nu_e}^{\text{obs}}(L_1) - \delta_0 P_{\nu_\mu \rightarrow \nu_e}^{\text{ND,obs}}$$

If it can be established within experimental uncertainties that $\chi_T^{\text{obs}} < 0$ and the conditions are fulfilled then T has to be violated in Nature.

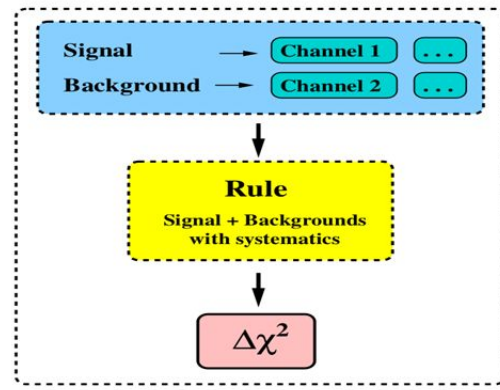
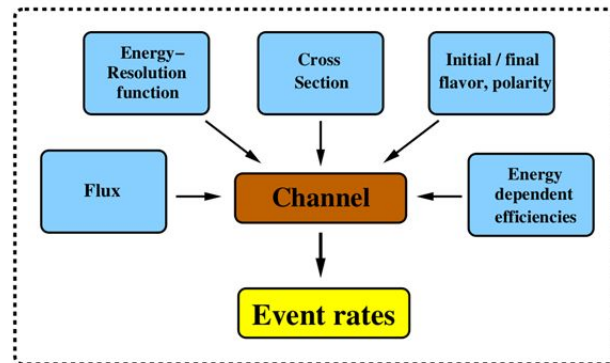


$E_\nu \in [0.80, 0.92] \text{ GeV}$ (neutrinos/NO, anti-neutrinos/IO)
 $E_\nu \in [0.86, 0.99] \text{ GeV}$ (neutrinos/IO, anti-neutrinos/NO)

Simulation Details

- Modified Probability engine for new physics
- Flux Spectrum
- Cross-Sections
- Event-Rates
- Bin-based energy smearing
- Detector Resolution
- Systematics

GLOBES Toolkit



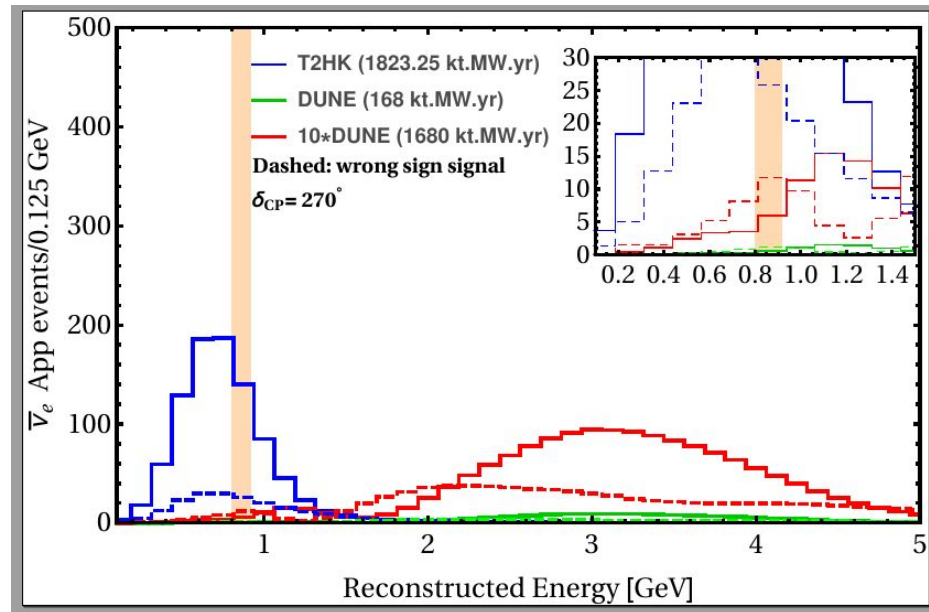
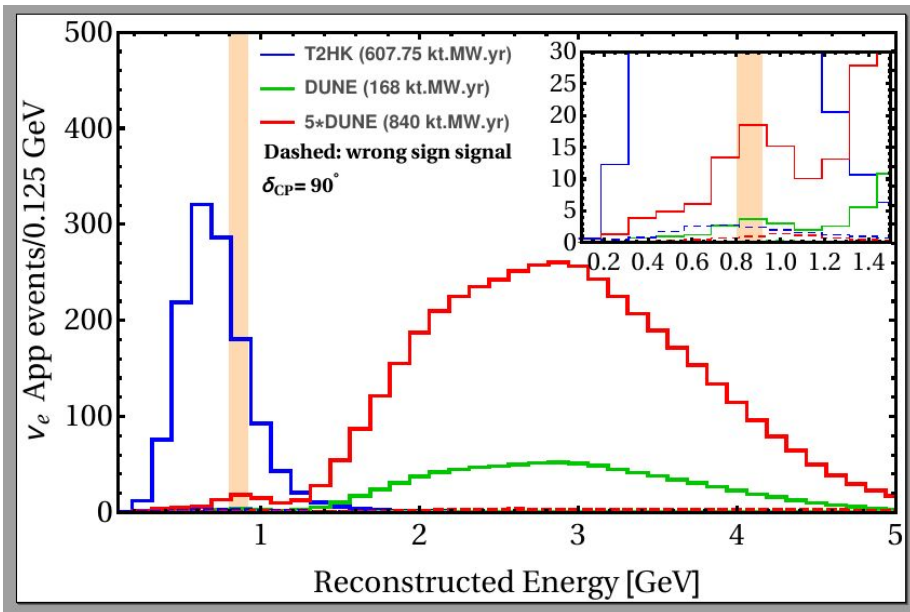
Hypothesis testing

$$\chi_{\text{even}}^2(E; \theta) = \sum_{a=1}^{N_L} \left[\frac{P_{\mu e}^{\text{even}}(L_a, E; \theta) - p_a^{\text{app}}}{\sigma_a^{\text{app}}} \right]^2$$

where, P_a^{app} is calculated in the standard three flavor scenario.

The value of $\chi^2_{\text{min}}(E)$ is an indication of T violation by data.

Event Spectra



Hypothesis testing

Default Configuration:

T2HK[L=295 km]:

Runtime(neutrino mode): 2.5 yrs

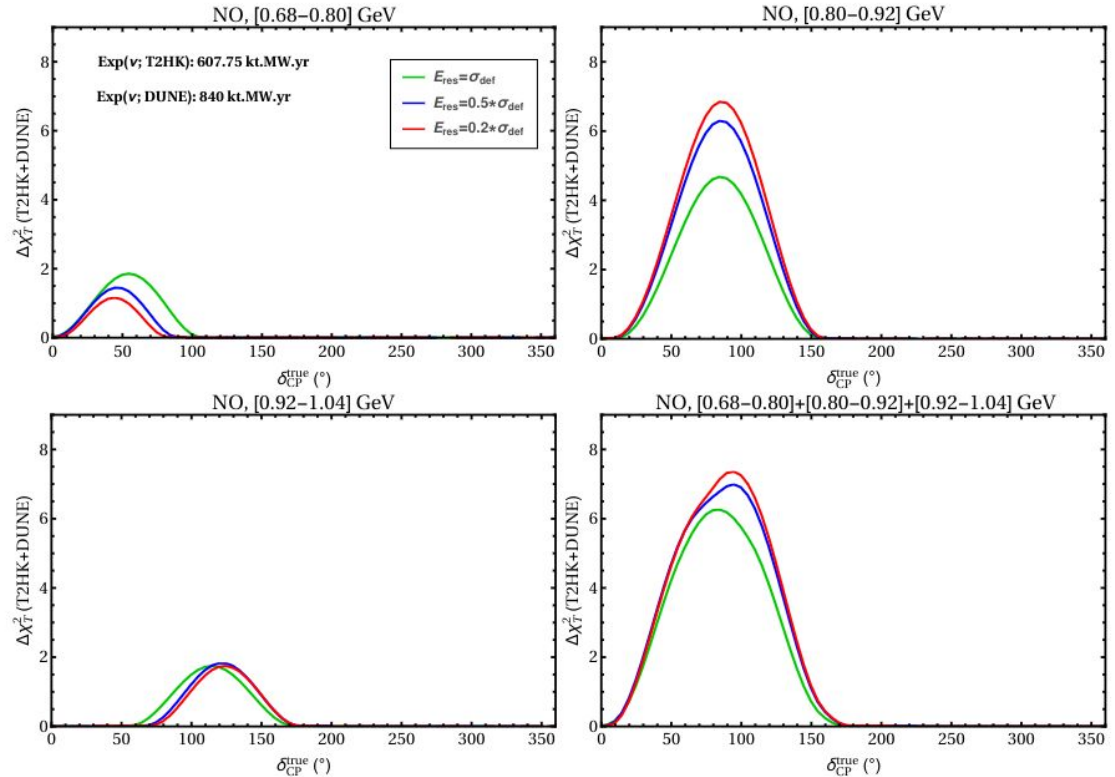
Detector mass: 187 kton

DUNE[L=1300 km]:

Runtime(neutrino mode): 3.5 yrs

Detector mass: 40 kton

*5 times of actual runtime for Dune.



Exposure and Resolution Effect

Standard Resolution:

T2HK~16%

DUNE~8.5% (better than CDR and TDR files)

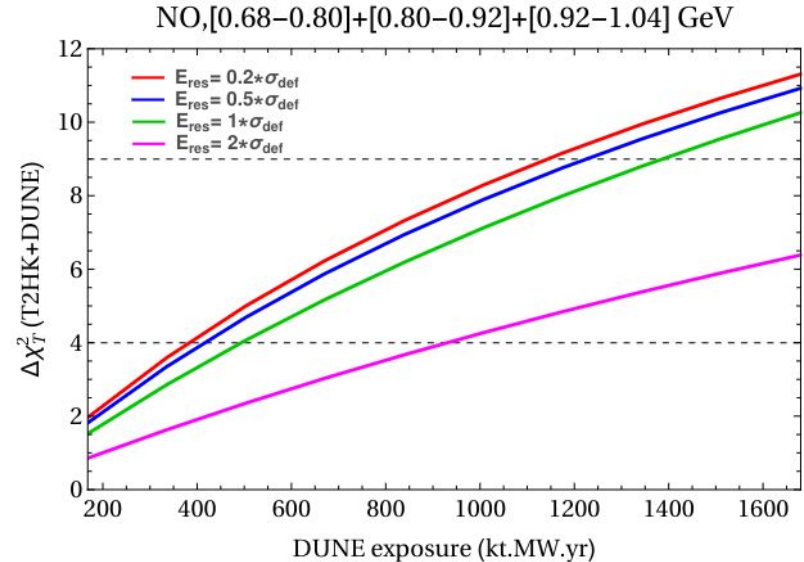
$$E_{\text{res}} = \alpha \cdot E + \beta \cdot \sqrt{E} + \gamma$$

$$\sigma_{\text{def}}(\text{T2HK}, \nu): \{\alpha, \beta, \gamma\} = \{0.12, 0.07, 0.0\}$$

$$\sigma_{\text{def}}(\text{T2HK}, \bar{\nu}): \{\alpha, \beta, \gamma\} = \{0.12, 0.0, 0.09\}$$

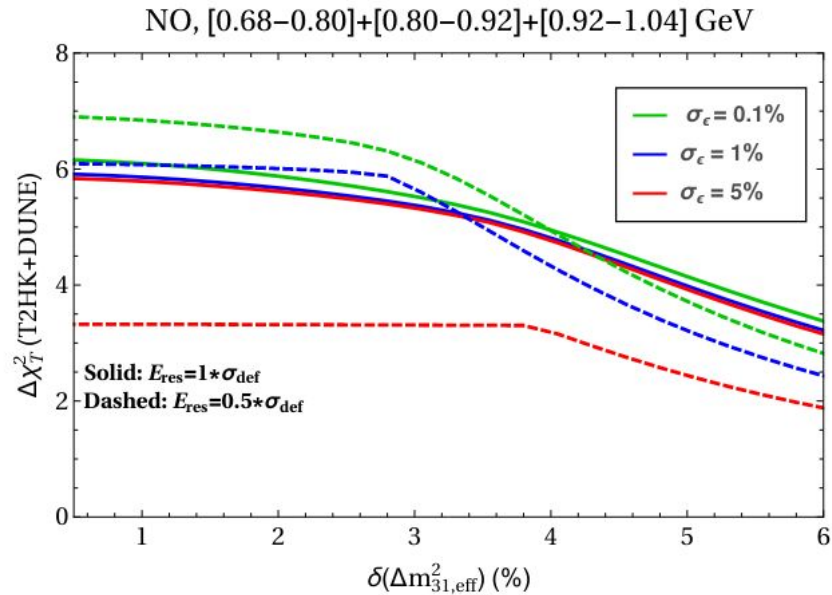
$$\sigma_{\text{def}}(\text{DUNE}, \nu): \{\alpha, \beta, \gamma\} = \{0.045, 0.001, 0.048\}$$

$$\sigma_{\text{def}}(\text{DUNE}, \bar{\nu}): \{\alpha, \beta, \gamma\} = \{0.026, 0.001, 0.085\}$$



Chatterjee, et.al , Impact of Improved Energy Resolution on DUNE sensitivity to Neutrino Non-Standard Interactions [[2106.04597](#)]

Zero-distance effect and prior on oscillation frequencies



Conclusions

- The variable X_T depending solely on oscillation probabilities provides an efficient way to probe T violation signature experimentally.
- We find the potential region for studying T violation with T2HK and DUNE at low energies.
- The improved statistics and better detector resolution, particularly for **DUNE**, plays a crucial role in improving sensitivities.
- There is a possibility of finding the potential region at higher energies and longer baselines.

Thank You 😊