## From Factorisation to **Cross Section Predictions**

The 2024 Annual Meeting of the CRC TRR 257 "Particle Physics Phenomenology after the Higgs Discovery", 12th March 2024

### Projects Bla, Blb and Ble

### M. Czakon **RWTH Aachen University**







• B1a: Bell, Czakon, Melnikov

with a resolved jet

- B1b: Czakon, Heinrich, Worek Precision top-quark physics at the LHC
- Bie: Bell, Czakon, Melnikov Power corrections in collider processes

• In this talk:

results from Aachen since last Annual Meeting including some results "in-between projects"

### N3LO QCD predictions for production cross sections of colourless systems in association



# **Power-Suppressed Effects**

## Subleading Soft-Gluon Effects at 1-Loop in QCD

Classic problem:



• Structure at next-to-leading power understood at tree-level QED by Low (1958), Burnett and Kroll (1968)

 $q \rightarrow 0$ 

- Necessity to include virtual collinear enhancements at higher orders noticed by <u>del Duca</u> (1990)
- Extension to tree-level QCD described in <u>1404.5551</u>, <u>1406.6987</u>, <u>1406.6574</u>
- Why bother (if you don't like pure theory)?

  - needed to obtain cross sections approximations at subleading power in different kinematic variables • can be used to improve numerical stability in cross section calculations

- At leading power: eikonal approximation you need to calculate a soft current,
- ? but at least the structure is understood !



## Subleading Soft-Gluon Effects at 1-Loop in QCD

- Several attempts to understand one-loop QCD amplitudes (more results for photon emission):
  - based on SCET: <u>1412.3108</u>, <u>1912.01585</u>, <u>2112.00018</u>
  - based on Feynman-diagram analysis: <u>1503.05156</u>, <u>1610.06842</u>
- Complete characterisation in Czakon, Eschment, Schellenberger, JHEP 12 (2023) 126

$$\begin{split} M_{g}^{(1)}(\{p_{i}+\delta_{i}\},q) & \Big\rangle = \mathbf{S}^{(0)}(\{p_{i}\},\{\delta_{i}\},q) \left| M^{(1)}(\{p_{i}\}) \right\rangle \\ & + \mathbf{S}^{(1)}(\{p_{i}\},\{\delta_{i}\},q) \left| M^{(0)}(\{p_{i}\}) \right\rangle + \int_{0}^{1} \mathrm{d}x \sum_{i} \mathbf{J}_{i}^{(1)}(x,p_{i},q) \left| H_{g,i}^{(0)}(x,\{p_{i}\},q) \right\rangle \\ & + \sum_{i \neq j} \sum_{\substack{\tilde{a}_{i} \neq a_{i} \\ \tilde{a}_{j} \neq a_{j}}} \mathbf{\tilde{S}}_{a_{i}a_{j}}^{(1)}(p_{i},p_{j},q) \left| M^{(0)}(\{p_{i}\}) \right|_{a_{j}}^{a_{i}} \rightarrow \tilde{a}_{i}} \right\rangle + \int_{0}^{1} \mathrm{d}x \sum_{\substack{i \neq j \\ a_{i} \neq a_{i}}} \mathbf{\tilde{J}}_{i}^{(1)}(x,p_{i},q) \left| H_{\bar{q},i}^{(0)}(x,\{p_{i}\},q) \right\rangle + \mathcal{O}(\lambda) \end{split}$$

$$\mathbf{P}_{g}(\sigma, c) \, \mathbf{S}^{(0)}(\{p_{i}\}, \{\delta_{i}\}, q) = -\sum_{i} \mathbf{T}_{i}^{c} \otimes \mathbf{S}_{i}^{(0)}(p_{i}, \delta_{i}, q, \sigma) \, \left| M^{(0)}(\{p_{i}\}) \right\rangle$$

$$\mathbf{S}_{i}^{(0)} = \frac{p_{i} \cdot \epsilon^{*}}{p_{i} \cdot q} + \frac{1}{p_{i} \cdot q} \left[ \left( \epsilon^{*} - \frac{p_{i} \cdot \epsilon^{*}}{p_{i} \cdot q} q \right) \cdot \delta_{i} + p_{i} \cdot \epsilon^{*} \sum_{j} \delta_{j} \cdot \partial_{j} + \frac{1}{2} F_{\mu\nu} \left( J_{i}^{\mu\nu} - \mathbf{K}_{i}^{\mu\nu} \right) \right]$$



## Subleading Soft-Gluon Effects at 1-Loop in QCD

Czakon, Eschment, Schellenberger, JHEP 12 (2023) 126

• Unexpected simplicity of contributions due to virtual collinear singularities

$$\begin{aligned} \mathbf{P}_{g}(\sigma,c) \mathbf{J}_{i}^{(1)}(x,p_{i},q) &= \frac{\Gamma(1+\epsilon)}{1-\epsilon} \left(-\frac{\mu^{2}}{s_{iq}}\right)^{\epsilon} \left(x(1-x)\right)^{-\epsilon} \epsilon^{*}(q,p_{i},\sigma) \cdot \epsilon(p_{i},-\sigma) \sum_{c'} \mathbf{P}_{g}(-\sigma,c') \\ &\times \left[ \left(\mathbf{T}_{i}^{c} \mathbf{T}_{i}^{c'} + \frac{1}{x} i f^{cdc'} \mathbf{T}_{i}^{d}\right) \otimes \left(-2 + x\left(1 + \mathbf{\Sigma}_{g,i}\right)\right) \right] \end{aligned}$$

• Exact evaluation of convolutions in the general case through products of tree-level amplitudes



- Tree-level subleading collinear asymptotics
- General proof of correctness via expansion-by-regions



• Numerical checks demonstrating potential gains in numerical stability in applications



Czakon, Eschment, Schellenberger, JHEP 12 (2023) 126





Jet Cross Sections at NNLO

Alvarez, Cantero, Czakon, Llorente, Mitov, Poncelet, JHEP 03 (2023) 129 ATLAS, JHEP 07 (2023) 85

• TEEC : transverse-energy-weighted distribution of the azimuthal differences between jet pairs in the final state:

$$\frac{1}{\sigma} \frac{\mathrm{d}\Sigma}{\mathrm{d}\cos\phi} \equiv \frac{1}{\sigma} \sum_{ij} \int \frac{\mathrm{d}\sigma}{\mathrm{d}x_{\mathrm{T}i} \mathrm{d}x_{\mathrm{T}j} \mathrm{d}\cos\phi} x_{\mathrm{T}i} x_{\mathrm{T}j} \mathrm{d}x_{\mathrm{T}i} \mathrm{d}x_{\mathrm{T}j}$$

• ATEEC : difference between the forward ( $\cos \phi > o$ ) and the backward (cos  $\phi$  < o) part of the TEEC function

$$\frac{1}{\sigma} \frac{\mathrm{d}\Sigma^{\mathrm{asym}}}{\mathrm{d}\cos\phi} = \frac{1}{\sigma} \frac{\mathrm{d}\Sigma}{\mathrm{d}\cos\phi} \bigg|_{\phi} - \frac{1}{\sigma} \frac{\mathrm{d}\Sigma}{\mathrm{d}\cos\phi} \bigg|_{\pi-\phi}$$

# Transverse Energy-Energy Correlations at ATLAS

 $= \frac{1}{N} \sum_{A=1}^{N} \sum_{ij} \frac{E_{Ti}^{A} E_{Tj}^{A}}{\left(\sum_{k} E_{Ti}^{A}\right)^{2}} \delta(\cos \phi - \cos \varphi_{ij})$ 









## Transverse Energy-Energy Correlations at ATLAS



### ATLAS JHEP 07 (2023) 85

Particle-level TEEC

√s = 13 TeV; 139 fb<sup>-1</sup>

anti- $k_{\rm t} R = 0.4$ 

 $p_{_{T}} > 60 \text{ GeV}$ 

|η| < 2.4

 $\mu_{R,F} = \mathbf{\hat{H}}_{T}$ 

 $\alpha_{s}(m_{z}) = 0.1180$ 

MMHT 2014 (NNLO)

- Data 



## **Transverse Energy-Energy Correlations at ATLAS**

### <u>ATLAS-CONF-2020-025</u> NLO QCD



### ATLAS, JHEP 07 (2023) 85 NNLO QCD





## Isolated $\gamma$ + di-jet at NNLO in QCD

Badger, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia, JHEP 10 (2023) 071

- Highlight: first complete  $2 \rightarrow 3$  process with full-color virtual amplitudes
- Analysis matching ATLAS measurements from JHEP 03 (2020) 179



• Unwanted dominant source of photons: hadron decays  $\Rightarrow$  remove by judicious cuts

<b>Requirements on photon</b>	$E_{\rm T}^{\gamma} > 150 { m GeV},   \eta^{\gamma}  < 2.37 \ ({ m excluding} \ 1.37 <  \eta^{\gamma}  < 1.50 { m GeV}$							
	$E_{\rm T}^{\rm iso} < 0.0042 \cdot E_{\rm T}^{\gamma} + 4.8 \text{ GeV} (\text{reconstruction level})$							
	$E_{\rm T}^{\rm iso} < 0.0042 \cdot E_{\rm T}^{\gamma} + 10 \text{ GeV} \text{ (particle level)}$							
<b>Requirements on jets</b>	at least two jets using anti- $k_t$ algorithm with $R = 0.4$							
	$ p_{ m T}^{ m jet} > 100{ m GeV}, y^{ m jet}  < 2.5,\Delta R^{\gamma- m jet} > 0.8$							
Phase space	total	fragmentation enriched	direct enriche					
		$E_{ m T}^{\gamma} < p_{ m T}^{ m jet2}$	$E_{\mathrm{T}}^{\gamma} > p_{\mathrm{T}}^{\mathrm{jet1}}$					
Number of events	755270	111666	386846					





## Isolated $\gamma$ + di-jet at NNLO in QCD

Badger, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia, JHEP 10 (2023) 071

• Scale choice:

$$\mu_R = \mu_F = H_T = E_\perp(\gamma)$$
  
 $\mu_R = \mu_F = E_\perp(\gamma)$ ,

- Very interesting comparison: ATLAS used SHERPA predictions for reference
  - NLO-matched QCD parton-shower merged with LO photon+four-jet samples

  - $E_{\perp}(\gamma)$  used for renormalisation and factorisation scale
- Why is this process interesting otherwise?
  - Non-back-to-back Born configurations

 $\rightarrow$  access to angular correlations between the photon and jets

- Access to different kinematic regimes through distinguishable photon  $\rightarrow$  enhance direct, high- or low-z fragmentation
- Background process for BSM:  $pp \rightarrow \gamma + Y(\rightarrow jj)$

 $\gamma) + p_T(j_1) + p_T(j_2)$ and

• In principle this corresponds to the double-real radiation contributions in the NNLO QCD predictions



Badger, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia, JHEP 10 (2023) 071

- Transverse photon energy
  - NNLO corrections from 1. 1% to 10%
  - Improved description of 2. data up to 1 TeV
  - Larger scale uncertainties 3. with the  $E_{\perp}(\gamma)$  scale in the inclusive phase space
  - Large experimental 4. uncertainties beyond 1 TeV
  - Default SHERPA predictions 5. with merging are a poor description with large uncertainties



## Isolated $\gamma$ + di-jet at NNLO in QCD



Badger, Czakon, Hartanto, Moodie, Peraro, Poncelet, Zoia, JHEP 10 (2023) 071

• Negligible size of the subleading color corrections



tt + X Cross Sections at NLO

### **Full Off-Shell Predictions**

- Modelling of unstable particles  $\Rightarrow ttWj$  production @  $\mathcal{O}(\alpha_s^4 \alpha^6)$ ullet
  - *Full off-shell* = *DR* + *SR* + *NR* + *interferences* + *BW propagators* \*
  - NWA = DR restricts unstable t & W to on-shell states \*

 $pp \rightarrow e^+ \nu_e \,\mu^- \bar{\nu}_\mu \,\tau^+ \nu_\tau \,b \bar{b} j + X$ 











### Bi, Kraus, Reinartz & Worek JHEP09 (2023) 026



### Additional Jet Activity in ttW Production



### Bi, Kraus, Reinartz & Worek JHEP09 (2023) 026

## Additional Jet Activity in ttW Production

- NLO cross section for both processes as function of number of resolved jets  $N_j \Rightarrow$  Resolved jets that pass all cuts
- Observed large scale dependence for exclusive NLO samples \*
- \* Gives rise to concern as jet vetoes are widely used in experimental analyses @ LHC

$(\mu_R,\mu_F)$	$\sigma_{H_T/2}^{t\bar{t}W^+j}(N_j=1)$ [ab]	$\sigma_{H_T/2}^{t\bar{t}W^+j}(N_j=2) \ [{ m ab}]$	$\sigma_{H_T/2}^{t\bar{t}W^+j}(N_j \ge 1) \text{ [ab]}$	-	$(\mu_R,\mu_F)$	$\sigma_{H_T/2}^{t\bar{t}W^+}(N_j=0)$ [ab]	$\sigma_{H_T/2}^{t\bar{t}W^+}(N_j = 1)$ [ab]	$\sigma_{H_T/2}^{t\bar{t}W^+}(N_j \ge 0) [$
$(\mu_0,\mu_0)$	$78.6^{+13\%}_{-48\%}$	$63.7^{+56\%}_{-34\%}$	$142.3^{+1.4\%}_{-8.1\%}$	-	$(\mu_0,\mu_0)$	$104.6^{+20\%}_{-44\%}$	$141.9^{+41\%}_{-27\%}$	$246.4^{+5.0\%}_{-7.0\%}$
$(\mu_R,\mu_F)$	$\delta \sigma^{t\bar{t}W^+j}_{H_T/2}(N_j = 1) \ [ab]$	$\delta \sigma_{H_T/2}^{t\bar{t}W^+j}(N_j=2) \; [ab]$	$\delta \sigma_{H_T/2}^{t\bar{t}W^+j}(N_j \ge 1)$ [ab]	-	$(\mu_R,\mu_F)$	$\delta \sigma^{t \bar{t} W^+}_{H_T/2}(N_j=0)$ [ab]	$\delta \sigma^{t\bar{t}W^+}_{H_T/2}(N_j=1)$ [ab]	$\delta \sigma^{t\bar{t}W^+}_{H_T/2}(N_j \ge 0)$
$(2\mu_0, 2\mu_0)$	+9.9(+13%)	-21.4(-34%)	-11.5(-8.1%)		$(2\mu_0,2\mu_0)$	+21.1(+20%)	-38.4(-27%)	-17.2(-7.0%
$(\mu_0/2,\mu_0/2)$	-37.4(-48%)	+35.7(+56%)	-1.6(-1.1%)		$(\mu_0/2,\mu_0/2)$	-45.7(-44%)	+58.1(+41%)	+12.4(+5.0%
$(2\mu_0,\mu_0)$	+9.0(+11%)	-18.0(-28%)	-9.0(-6.3%)		$(2\mu_0,\mu_0)$	+19.1(+18%)	-32.3(-23%)	-13.2(-5.3%
$(\mu_0/2,\mu_0)$	-28.7(-37%)	+27.8(+44%)	-0.9(-0.6%)		$(\mu_0/2,\mu_0)$	-40.0(-38%)	+46.6(+33%)	+6.6(+2.7%)
$(\mu_0,2\mu_0)$	+3.5(+4.5%)	-4.8(-7.5%)	-1.2(-0.8%)		$(\mu_0,2\mu_0)$	+4.1(+3.9%)	-7.8(-5.5%)	-3.8(-1.5%)
$(\mu_0,\mu_0/2)$	-3.5(-4.5%)	+5.5(+8.6%)	+2.0(+1.4%)		$(\mu_0,\mu_0/2)$	-3.1(-2.9%)	+8.7(+6.1%)	+5.6(+2.3%)

 $pp \rightarrow e^+ \nu_e \, \mu^- \bar{\nu}_\mu \, \tau^+ \nu_\tau \, b\bar{b}j + X$ 

NLO ttW& ttWj

\* Inclusive NLO results for  $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu \tau^+ \nu_\tau b \bar{b} j + X \& pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu \tau^+ \nu_\tau b \bar{b} + X$  have expected reduced theoretical uncertainties

$$pp \rightarrow e^+ \nu_e \,\mu^- \bar{\nu}_\mu \,\tau^+ \nu_\tau \,b\bar{b} + X$$

Bi, Kraus, Reinartz & Worek JHEP09 (2023) 026



## Photons in Production & Decays

 $pp \to \ell^+ \nu_\ell \, \ell^- \bar{\nu}_\ell \, b \bar{b} \, \gamma \gamma$ 

• Integrated fiducial cross-section level @ NLO in QCD

 $p_{T,b} > 25 \text{ GeV}, p_{T,\gamma} > 25 \text{ GeV}$ :

- \* *Mixed contribution* at the level of 44%
- \* *Prod. contribution* at the level of 40%
- \* *Decay contribution* is about half the size 16%
- Differential fiducial cross-section level
- \* Various phase-space regions with various effects

<u>Stremmer</u>, Worek JHEP08 (2023) 179

### NLO $tt\gamma\gamma$





## Photons in Production & Decays

 $pp \rightarrow \ell^- \bar{\nu}_\ell j j b \bar{b} \gamma \gamma$ 

$$d\sigma_{\rm Full} = \overbrace{d\sigma_{t\bar{t}\gamma\gamma} \times \frac{d\Gamma_t}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_t} + \overline{d\sigma_{t\bar{t}\gamma}} \times \left(\frac{d\Gamma_{t\gamma}}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_t} + \frac{d\Gamma_t}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}\gamma}}{\Gamma_t}\right)}{+ \underbrace{d\sigma_{t\bar{t}} \times \left(\frac{d\Gamma_{t\gamma\gamma}}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_t} + \frac{d\Gamma_t}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_t} + \frac{d\Gamma_t}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_t} + \frac{d\Gamma_{t\gamma\gamma}}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_t}\right)}{\sigma_{\rm Decay}}}.$$

• For *lepton* + *jet* channel additional cut needed

$$|m_W - M_{jj}| < Q_{cut} = 15 \,\text{GeV}$$

- Suppress kinematical configurations from real radiation
- \* Jets originating from *W* recombined into single jet
- \* Extra jet from real radiation gives rise to second jet

### <u>Stremmer</u>, Worek JHEP08 (2023) 179

### NLO $tt\gamma\gamma$









## Photons in Production & Decays

 $pp \rightarrow \ell^- \bar{\nu}_\ell j j b \bar{b} \gamma \gamma$ 

$$d\sigma_{\text{Full}} = \overbrace{d\sigma_{t\bar{t}\gamma\gamma} \times \overbrace{\Gamma_{t}}^{\sigma_{\text{Prod.}}} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_{t}}}{\Gamma_{t}} + \overbrace{d\sigma_{t\bar{t}\gamma} \times \left(\frac{d\Gamma_{t\gamma}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_{t}} + \frac{d\Gamma_{t}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}\gamma}}{\Gamma_{t}}\right)}{+ \underbrace{d\sigma_{t\bar{t}} \times \left(\frac{d\Gamma_{t\gamma\gamma}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_{t}} + \frac{d\Gamma_{t}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_{t}} + \frac{d\Gamma_{t\gamma\gamma}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_{t}} + \frac{d\Gamma_{t\gamma\gamma}}{\Gamma_{t}} \times \frac{d\Gamma_{\bar{t}\gamma\gamma}}{\Gamma_{t}}\right)}{\sigma_{\text{Decay}}}.$$

- Integrated fiducial cross-section level @ NLO in QCD with & without  $|m_W - M_{jj}| < Q_{cut} = 15 \text{ GeV}$  $p_{T,b} > 25 \text{ GeV}, p_{T,j} > 25 \text{ GeV}, p_{T,\gamma} > 25 \text{ GeV}$ :
- \* *Prod. contribution* at the level of  $48\% \Rightarrow 40\%$
- \* *Mixed contribution* at the level of  $40\% \Rightarrow 43\%$
- \* *Decay contribution* is about half the size  $12\% \Rightarrow 17\%$

<u>Stremmer</u>, Worek JHEP08 (2023) 179

### NLO $tt\gamma\gamma$





# The N3LO Project

### **Ambitious Plans**

- Phenomenology of di-boson final states:  $\gamma\gamma$ ,  $W^+W^-$ , ZZ,  $W^\pm\gamma$ , Z $\gamma$
- Predicitions for:  $X + \text{jet}, X \in \{\gamma, W^{\pm}, Z, H\}$
- Methodology following the project: N-jettiness slicing  $q'(\bar{q}')$  $\bar{q}'(q'$ Beam functions for N-jettiness at N3LO in perturbative QCD Baranowski, Behring, Melnikov, Tancredi, Wever, JHEP 02 (2023) 073 One-loop corrections to the double-real emission contribution to the zero-jettiness soft function at N3LO in QCD Baranowski, Delto, Melnikov, Pikelner, Wang, 2401.05245 N-jettiness soft function at next-to-next-to-leading order in perturbative QCD Agarwal, Melnikov, Pedron, <u>2403.03078</u>
- How about a proper subtraction scheme? Need limits of amplitudes Revisiting the double-soft asymptotics of one-loop amplitudes in massless QCD Czakon, Eschment, Schellenberger, JHEP 04 (2023) 065



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# Conclusions

## Achievements since last Annual Meeting

bosons, top quarks and jets in the final state at NLO and NNLO

• Theoretical progress on factorisation in QCD

be reflected in publications

• Many results for phenomenologically relevant processes with gauge

• Substantial progress towards N3LO-precise predictions that will soon

