



Charged lepton flavour violation: theory overview

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Flavour in the **Standard Model**: interactions between *fermion* families (and the *Higgs*)

Y_{ij}^u, Y_{ij}^d and $Y_{ij}^\ell \rightsquigarrow$ encode flavour dynamics (masses, mixings & CP violation)
flavour-universal gauge interactions

SM quark sector: 6 massive states

flavour violated in charged current interactions $V_{CKM}^{ij} W^\pm \bar{q}_i q'_j$

conservation of total **baryon number** in SM interactions

CP violation sources: δ_{CKM} and θ_{QCD} 
(strongly constrained by tiny neutron EDM)

not enough to explain observed **BAU** from baryogenesis

Extensive probes of the “**CKM paradigm**”: meson oscillation and decays, CP violation...

... and a roller-coaster ride for hints of New Physics in recent years!

SM lepton sector: (strictly) massless neutrinos

conservation of total **lepton number** and **lepton flavours**

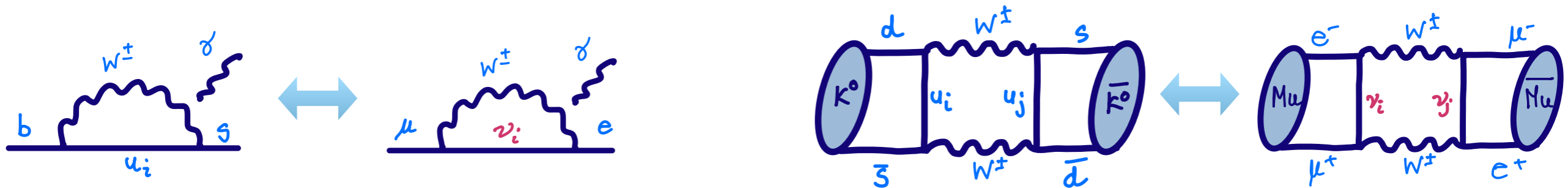
tiny leptonic EDMs (4-loop... $d_e^{CKM} \leq 10^{-38} e \text{ cm}$)

Neutrino oscillations: SM description insufficient! First laboratory discovery of New Physics!

Lepton flavours: from ν oscillations...

Neutrino oscillations: SM description insufficient! Added complexity to the flavour problem...

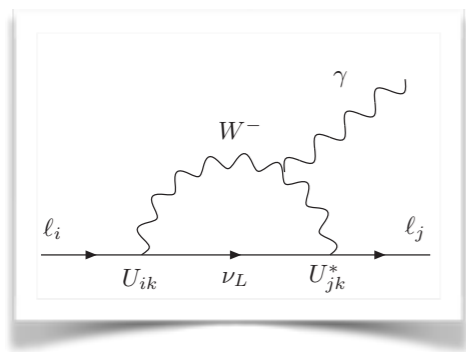
Violation of lepton flavour in neutral lepton sector opens a wide door to **flavour violation in the charged lepton sector!**



How general is this once we extend the SM to accommodate $\nu_\alpha \leftrightarrow \nu_\beta$?

In the most minimal extension SM_{m_ν}

[SM_{m_ν} = “ad-hoc” m_ν (Dirac), U_{PMNS}]



$$\Gamma(\mu \rightarrow e\gamma) = \frac{m_\mu^5}{16\pi} (|A_L|^2 + |A_R|^2)$$

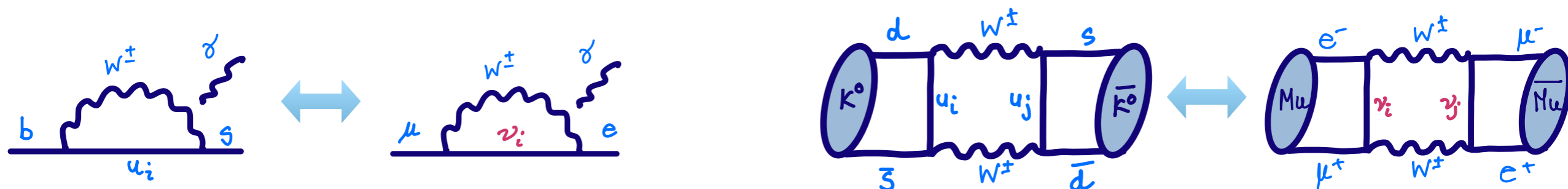
$$\sum_i \frac{U_{ei}^* U_{\mu i}}{(k^2 - m_i^2)} = \sum_i \frac{U_{ei}^* U_{\mu i}}{k^2} + \sum_i \frac{U_{ei}^* U_{\mu i}}{k^2} \left(\frac{m_i^2}{k^2} \right) + \mathcal{O} \left(\frac{m_i^4}{k^4} \right) \quad \& \quad \sum_i U_{ei}^* U_{\mu i} \frac{m_i^2}{M_W^2} = U_{e2}^* U_{e2} \frac{\Delta m_{21}^2}{M_W^2} + U_{e3}^* U_{e3} \frac{\Delta m_{31}^2}{M_W^2}$$

$$\text{BR}(\mu \rightarrow e\gamma) = \frac{3\alpha_e}{32\pi} \left| \sum_i U_{ei}^* U_{\mu i} \frac{m_i^2}{M_W^2} \right|^2 \Rightarrow \text{BR}(\mu \rightarrow e\gamma) \approx 10^{-54 \div -55}$$

Lepton flavours: from ν oscillations...

Neutrino oscillations: SM description insufficient! Added complexity to the flavour problem...

Violation of lepton flavour in neutral lepton sector opens a wide door to **flavour violation in the charged lepton sector!**



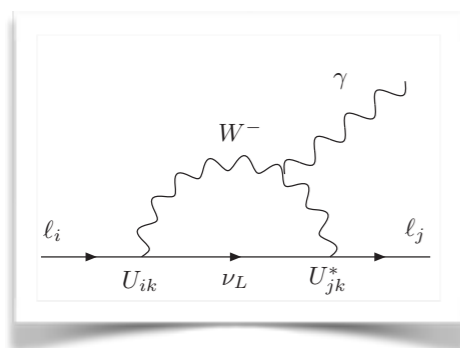
How general is this once we extend the SM to accommodate $\nu_\alpha \leftrightarrow \nu_\beta$?

In the most minimal extension SM_{m_ν} [SM_{m_ν} = "ad-hoc" m_ν (Dirac), U_{PMNS}]

total lepton number still conserved (LNC)

lepton EDMs still beyond observation (2-loop contributions from δ_{CP})

cLFV possible... but **not observable!!** $BR(\mu \rightarrow e\gamma) \sim 10^{-54}$

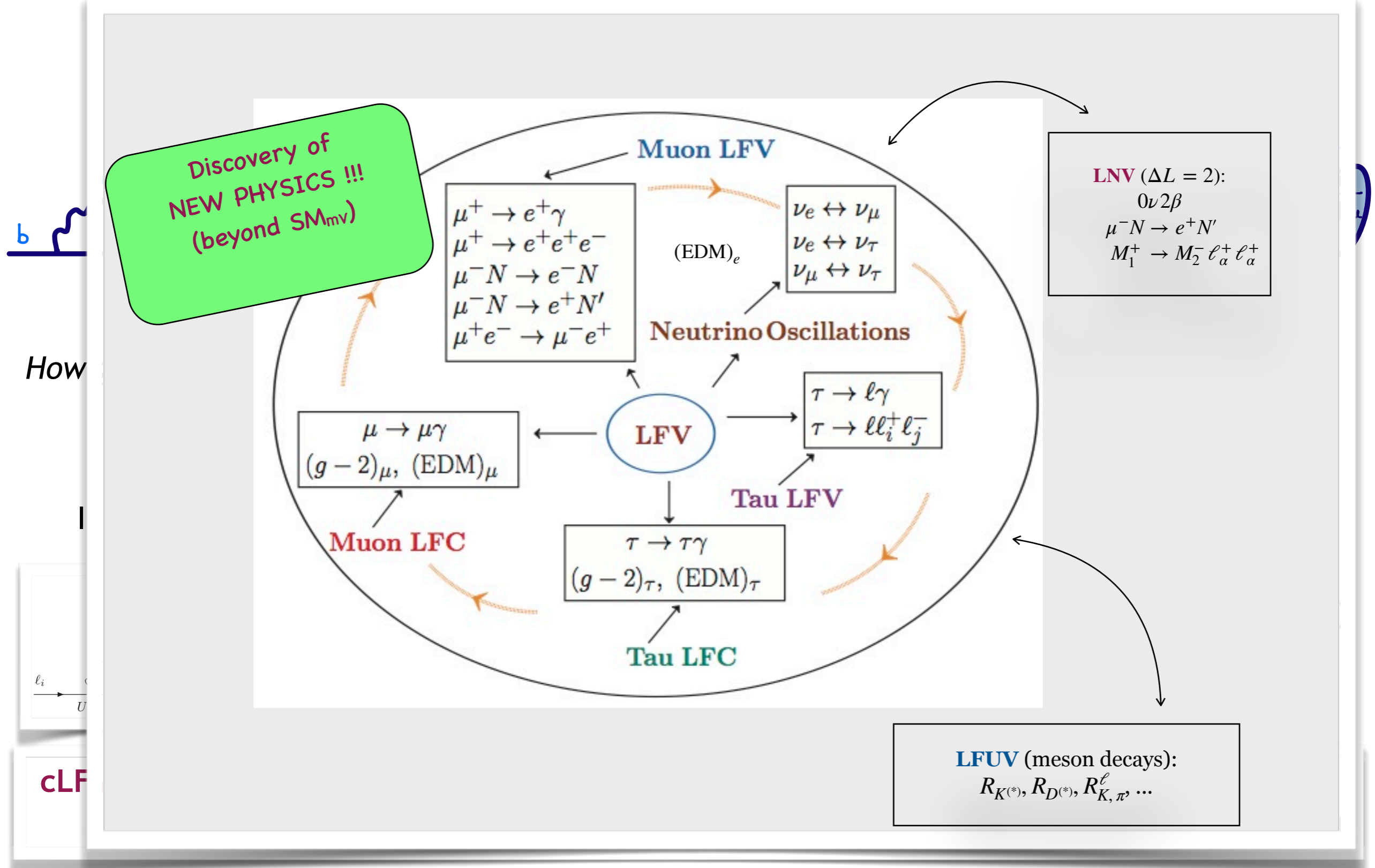


cLFV, LNV, lepton EDMs, ...: observation of **SM-forbidden leptonic modes**

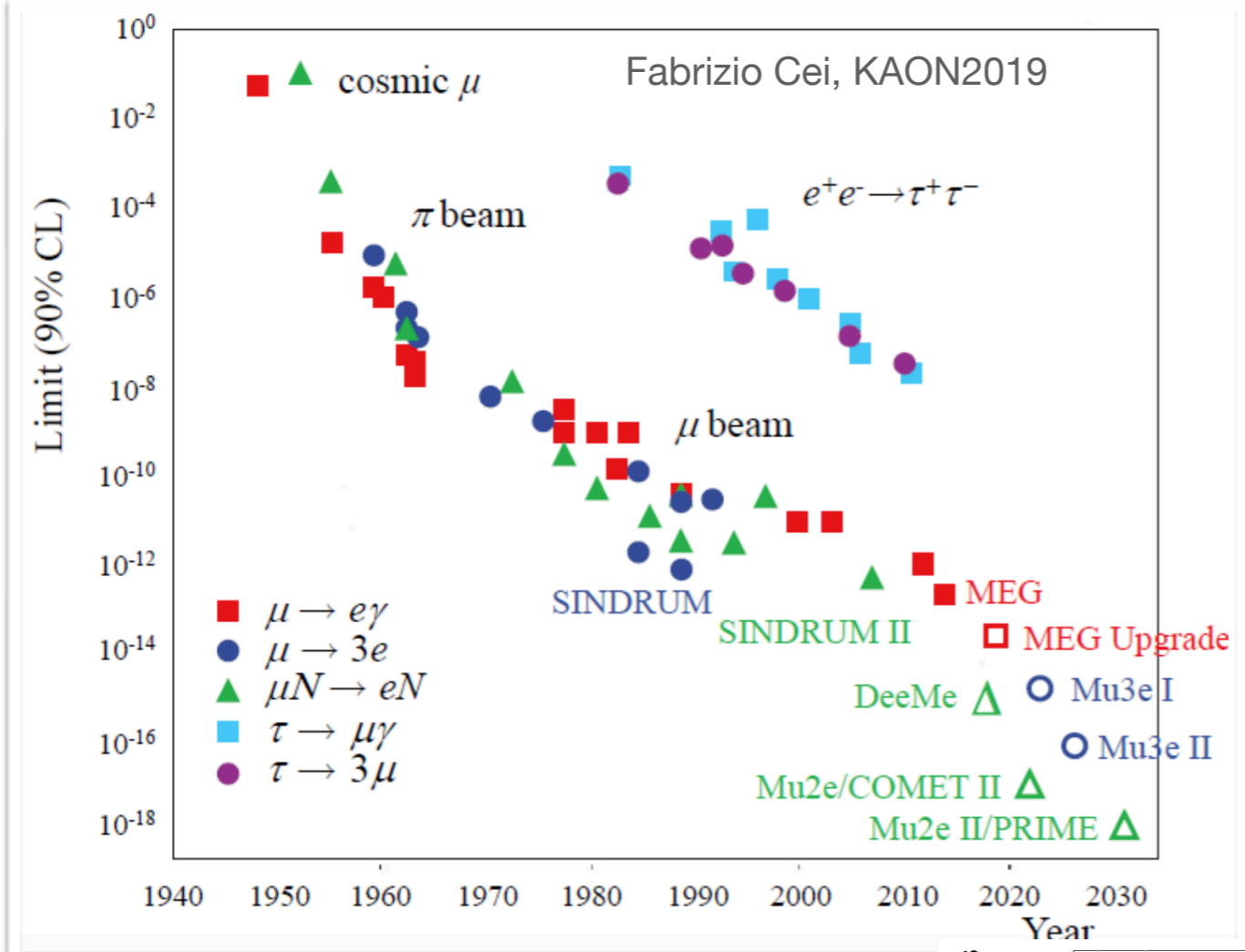
\Rightarrow **Discovery of New Physics!** (possibly before direct signal @ LHC)

Lepton flavours: from ν oscillations...

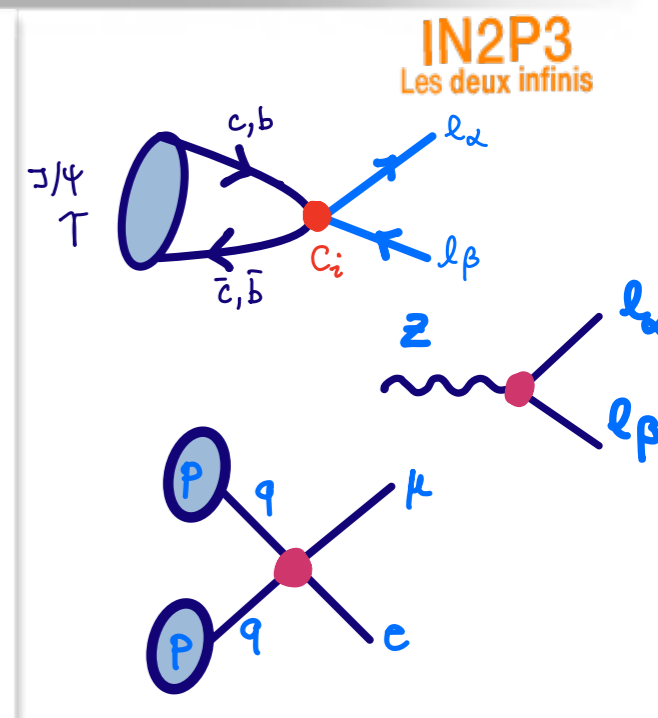
Neutrino oscillations: SM description insufficient! Added complexity to the flavour problem...



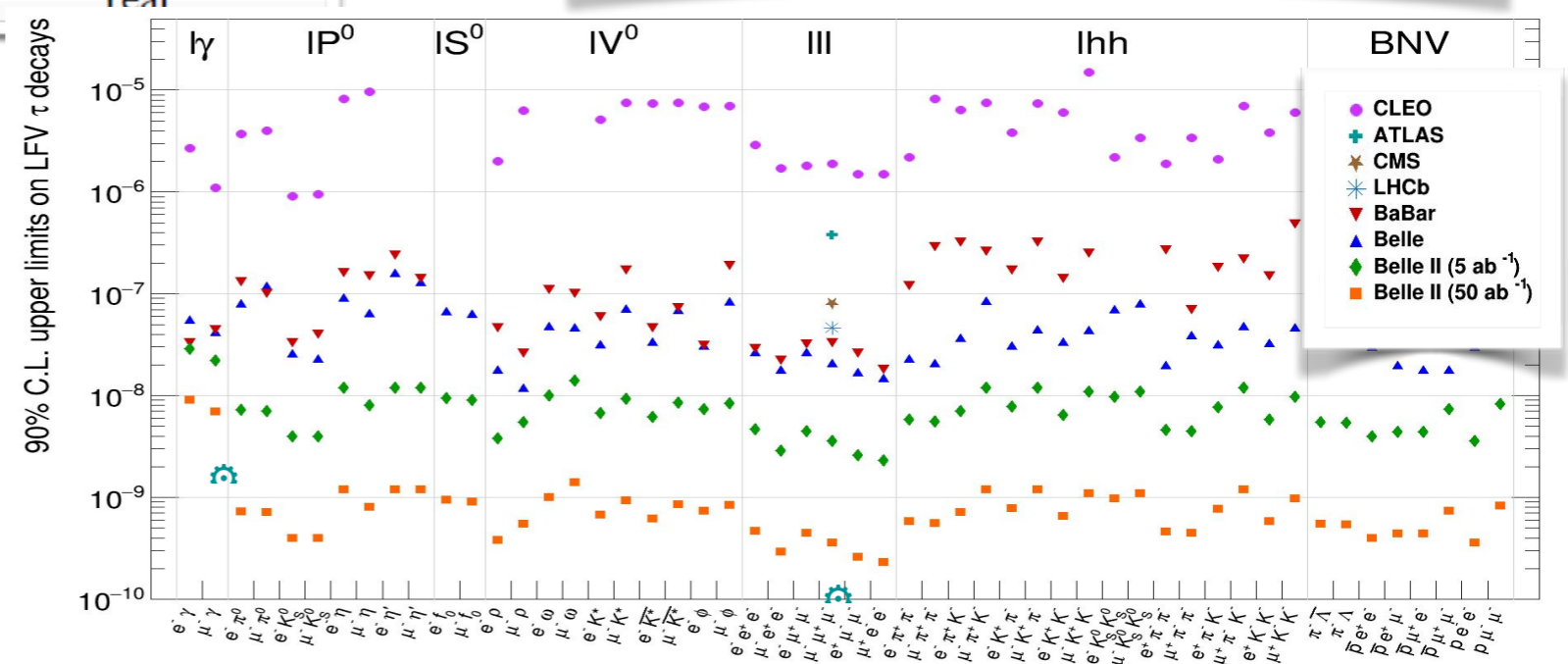
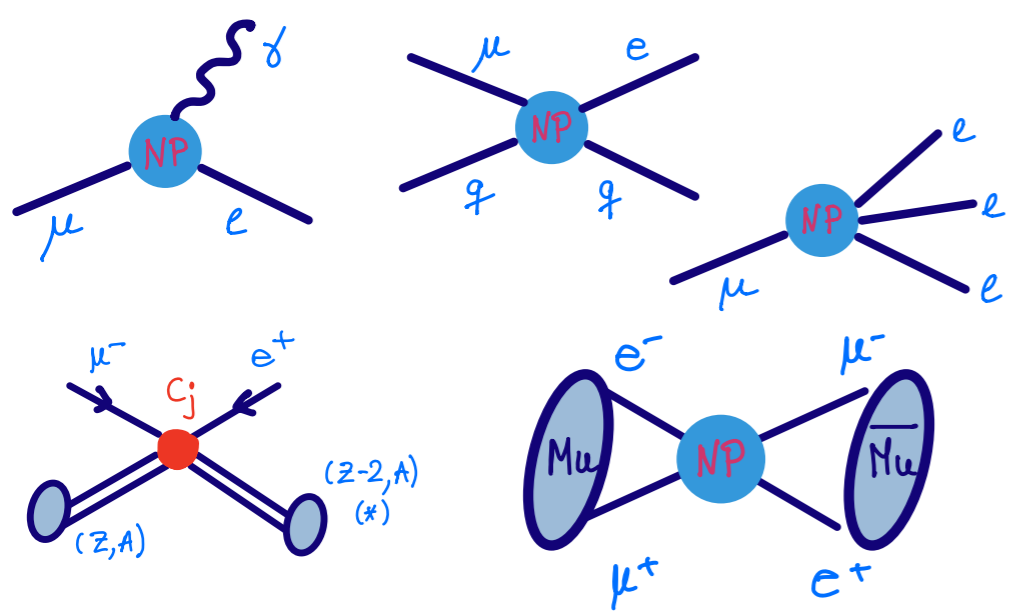
Amazing prospects for NP searches - cLFV!



$\pi^0 \rightarrow \mu e$	$< 3.6 \times 10^{-10}$
$K_L^0 \rightarrow \mu e$	$< 4.7 \times 10^{-12}$
$K_L^0 \rightarrow \pi^0 \mu^+ e^-$	$< 7.6 \times 10^{-11}$
$K^+ \rightarrow \pi^+ \mu^+ e^-$	$< 1.3 \times 10^{-11}$
$J/\psi \rightarrow \mu e$	$< 1.5 \times 10^{-7}$
$J/\psi \rightarrow \tau e$	$< 8.3 \times 10^{-6}$
$J/\psi \rightarrow \tau \mu$	$< 2.0 \times 10^{-6}$
$B^0 \rightarrow \mu e$	$< 2.8 \times 10^{-9}$
$B^0 \rightarrow \tau e$	$< 2.8 \times 10^{-5}$
$B^0 \rightarrow \tau \mu$	$< 2.2 \times 10^{-5}$
$B \rightarrow K \mu e^{(b)}$	$< 3.8 \times 10^{-8}$
$B \rightarrow K^* \mu e^{(b)}$	$< 5.1 \times 10^{-7}$
$B^+ \rightarrow K^+ \tau \mu$	$< 4.8 \times 10^{-5}$
$B^+ \rightarrow K^+ \tau e$	$< 3.0 \times 10^{-5}$
$B_s^0 \rightarrow \mu e$	$< 1.1 \times 10^{-8}$
$\Upsilon(1s) \rightarrow \tau \mu$	$< 6.0 \times 10^{-6}$



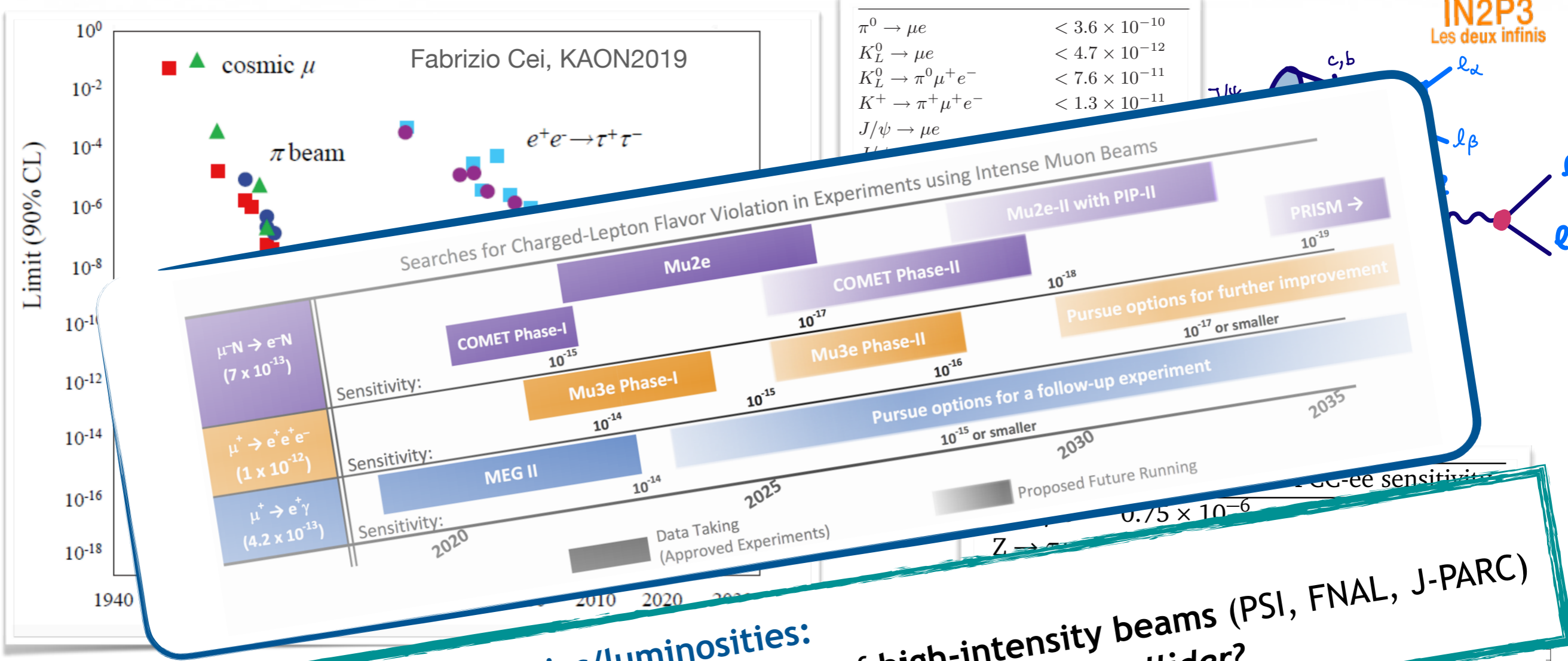
Decay	Present bound	FCC-ee sensitivity
$Z \rightarrow \mu e$	0.75×10^{-6}	$10^{-10} - 10^{-8}$
$Z \rightarrow \tau \mu$	12×10^{-6}	10^{-9}
$Z \rightarrow \tau e$	9.8×10^{-6}	10^{-9}



Amazing prospects for NP searches - cLFV!

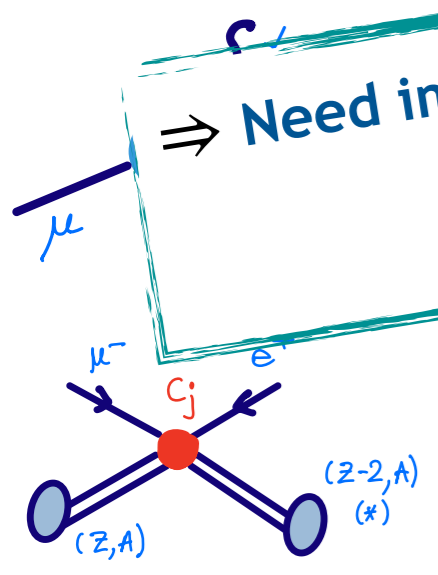


IN2P3
Les deux infinis



⇒ Need impressive intensities/luminosities:
 Amazing prospects with advent of high-intensity beams (PSI, FNAL, J-PARC)
 and beyond?... FCC-ee? Muon facility? Muon collider?

See presentations by G. Dal Maso, S. Müller,
 Y. Kwon, R. Fantechi



Strong arguments in favour of New Physics!

► Neutrino oscillations: 1st laboratory ("flavoured") evidence of NP

- ⇒ massive neutrinos and leptonic mixings U_{PMNS}^{ai}
- ⇒ New (Majorana) fields? New sources of CP violation?
 $\Delta L \neq 0$ and leptogenesis... (?)
- ⇒ **Open door for cLFV transitions and decays!**

Observations unaccounted for in the SM:

baryon asymmetry of the Universe, viable dark matter candidate,
neutrino oscillations (and some "tensions"...))

And a number of **theoretical** caveats...

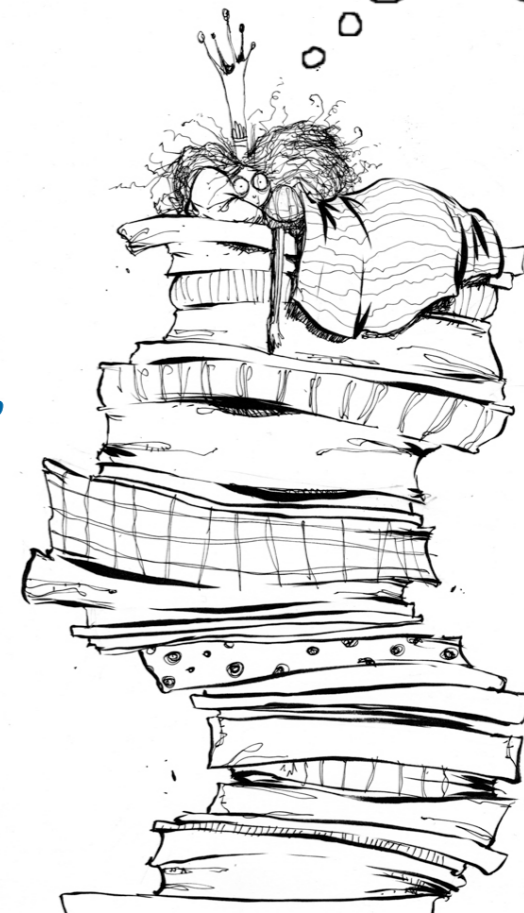
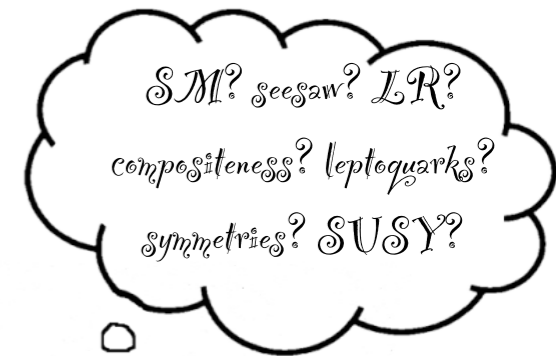
Many hints and a **clear necessity** of **New Physics...**

Which NP model? Realised at which scale Λ_{NP} ?

⇒ **Unique opportunities to search for NP in the lepton sector via cLFV**

first characterisation of New Physics (scale, interactions) - EFT approach;

exploring connections to ν mass generation! (among many other possible BSM!)



New Physics EFT quests with (muon) cLFV



EFT approach to New Physics

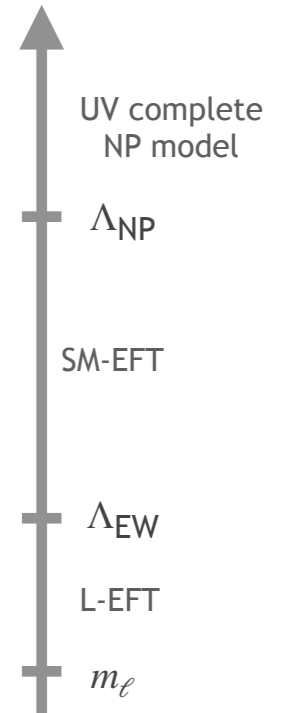
SM interpreted as a **low-energy limit** of a (complete, yet unknown) NP model
 \Rightarrow **Model-independent, effective field theory approach (EFT)**

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \sum_{n \geq 5} \frac{1}{\Lambda^{n-4}} \mathcal{C}^n(g, Y, \dots) \mathcal{O}^n(\ell, q, H, \gamma, \dots)$$

effective operators

(unknown) NP scale effective coefficients

$\mathcal{O}^5 \rightsquigarrow$ Weinberg operator (m_ν)
 $\mathcal{O}^6 \rightsquigarrow$ flavoured contributions
 (among many others!)



Derive the new "effective" interactions (vertices, ...), and compute **contributions to observables**
 Agnostic approach, allowing to generically parametrise NP effects
 on observables **forbidden** in SM and/or observables **suggesting deviations** from SM

$$\mathcal{A} \sim \mathcal{A} \left(\frac{\mathcal{C}^6}{\Lambda_{\text{NP}}^2} \right) + \dots$$

$$\mathcal{A} \sim \mathcal{A}^{\text{SM}} + \mathcal{A} \left(\frac{\mathcal{C}^6}{\Lambda_{\text{NP}}^2} \right) + \dots$$

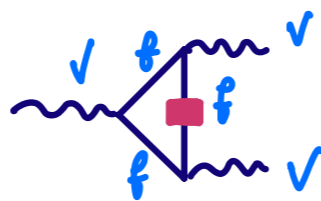
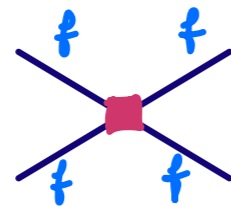
\Rightarrow master SM prediction!

EFT approach to New Physics

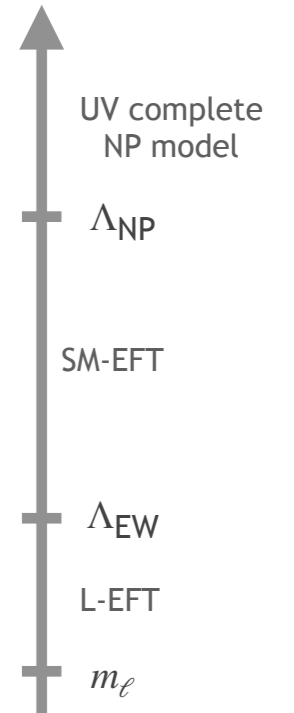
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(unknown) NP scale
effective coefficients
effective operators



$\mathcal{O}^5 \sim$ Weinberg operator (m_ν)
 $\mathcal{O}^6 \sim$ flavoured contributions
 (among many others!)



Cast **current data** (limits, ...) in terms of \mathcal{C}_{ij}^6 and Λ_{NP}^2

and attempt at **inferring info** on the **dominant operator**, and **scale of NP**

\Rightarrow Beyond $(V - A)$ structure? New vector/axial, (pseudo)scalar or tensor currents?
 Flavour violation beyond SM flavour paradigm?

\Rightarrow But **many unknowns**: minimal assumptions must be made, e.g.

"natural" $\Lambda_{\text{NP}} \rightarrow$ constrain \mathcal{C}_{ij}^6

"natural" $\mathcal{C}_{ij}^6 \approx 1 \rightarrow$ hint on Λ_{NP}

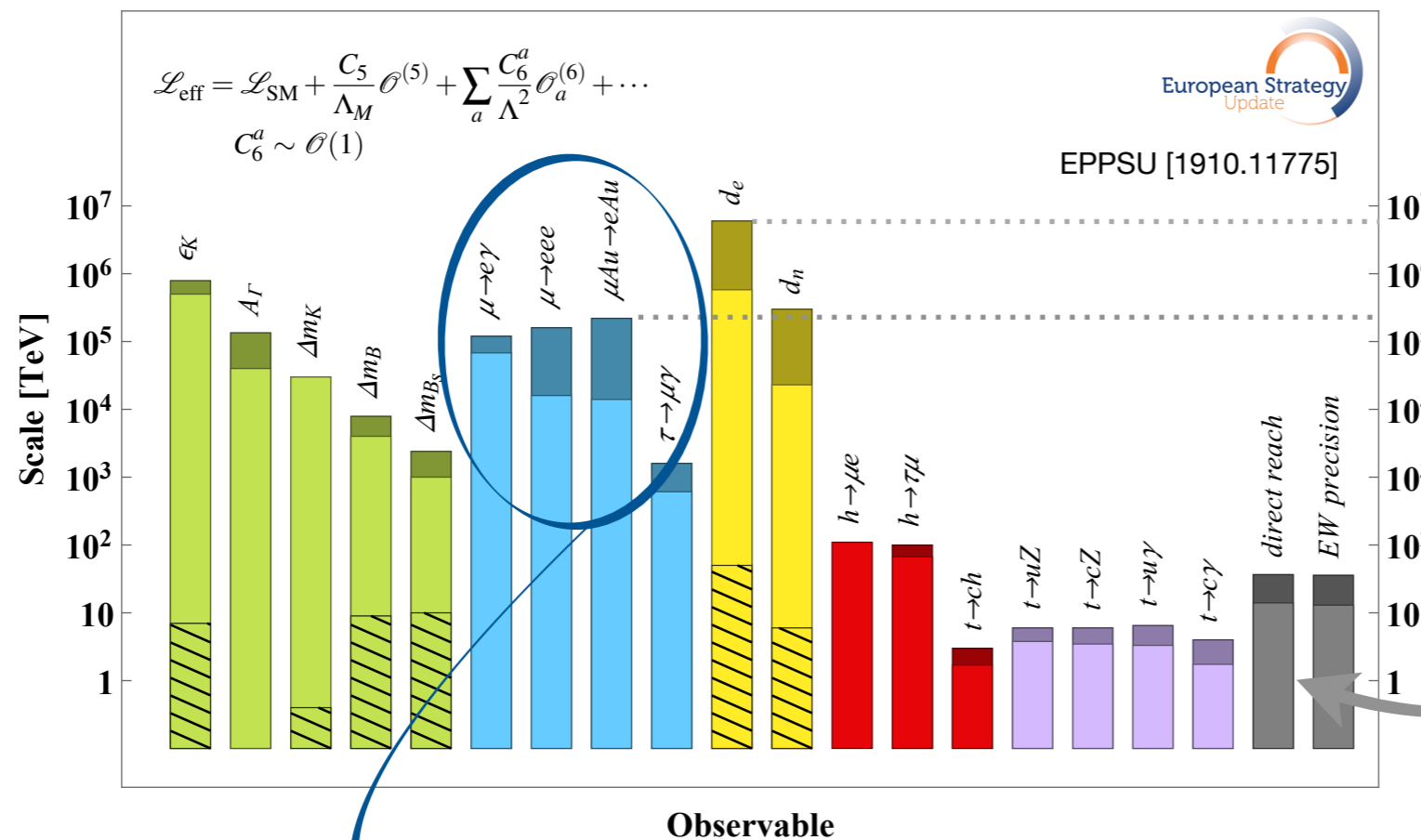
The probing power of flavour & CPV

SM interpreted as a **low-energy limit** of a (complete, yet unknown) NP model

⇒ **Model-independent, effective approach (EFT)**

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \sum_{n \geq 5} \frac{1}{\Lambda^{n-4}} \mathcal{C}^n(g, Y, \dots) \mathcal{O}^n(\ell, q, H, \gamma, \dots)$$

Cast **current data** in terms of \mathcal{C}_{ij}^6 and Λ_{NP} : $\mathcal{C}_{ij}^6 \approx 1 \Rightarrow$ bounds on Λ_{NP}



Flavour observables:
probes sensitive to very high NP scales

$$\Lambda_{\text{NP}} \sim \mathcal{O}(10^5 \text{ TeV})$$

well beyond collider's reach!

charged lepton flavour violating observables!

► Generic New Physics observables in the **lepton sector**:

- **Lepton number violation** (e.g. neutrino masses, $0\nu 2\beta$ decays, ...)
- **Electric and (anomalous) magnetic moments** - $d_\ell, (g - 2)_\ell$
- **charged lepton flavour violation**

Cast observables in terms of \mathcal{C}_{ij} and Λ_{NP} ; Apply **current data** (limits, ...)

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \frac{\mathcal{C}_5 \mathcal{O}^5}{\Lambda_{\text{LNV}}} (m_\nu) + \frac{\mathcal{C}_6 \mathcal{O}^6}{\Lambda_{\text{CLFV}}^2} (\ell_\alpha \leftrightarrow \ell_\beta) + \dots + \frac{\mathcal{C}_7 \mathcal{O}^7}{\Lambda_{\text{LNV}}^3} (0\nu 2\beta) + \dots$$

Majorana ν masses

Kinetic corrections, ...

EW precision, top physics, ...

Electric dipole & anomalous magnetic moments, ...

cLFV (dipole, 3 body, matter assisted, ...)

Lepton number violation,
cLFV & **LNV**,
...

⇒ **cLFV data to constrain \mathcal{C}_{ij} and/or infer sensitivity of process to large sets of \mathcal{C}_{ij}**

⇒ **Hints on Λ_{NP} (and on properties of new states & nature of couplings)**

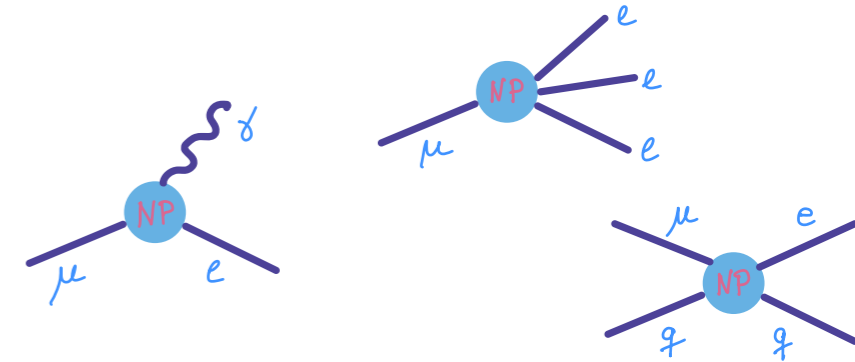
Deceptively simple task... different new physics scales, numerous operators!

Technically very involved, even if no "SM background"...

Muon cLFV: EFT approach to New Physics

Cast **current data** (limits, ...) in terms of \mathcal{C}_{ij} and Λ_{NP} : cLFV operators (\mathcal{O}^6)

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \frac{\mathcal{C}_5 \mathcal{O}^5}{\Lambda_{\text{LNV}}} (m_\nu) + \frac{\mathcal{C}_6 \mathcal{O}^6}{\Lambda_{\text{CLFV}}^2} (\ell_\alpha \leftrightarrow \ell_\beta) + \dots$$



Simple **"one-at-a-time"** limits:

	Br($\mu^+ \rightarrow e^+ \gamma$)		Br($\mu^+ \rightarrow e^+ e^- e^+$)		Br $_{\mu \rightarrow e}^{\text{Au/Al}}$	
	$4.2 \cdot 10^{-13}$	$4.0 \cdot 10^{-14}$	$1.0 \cdot 10^{-12}$	$5.0 \cdot 10^{-15}$	$7.0 \cdot 10^{-13}$	$1.0 \cdot 10^{-16}$
C_L^D	$1.0 \cdot 10^{-8}$	$3.1 \cdot 10^{-9}$	$2.0 \cdot 10^{-7}$	$1.4 \cdot 10^{-8}$	$2.0 \cdot 10^{-7}$	$2.9 \cdot 10^{-9}$
$C_{ee}^{S LL}$	$4.8 \cdot 10^{-5}$	$1.5 \cdot 10^{-5}$	$8.1 \cdot 10^{-7}$	$5.8 \cdot 10^{-8}$	$1.4 \cdot 10^{-3}$	$2.1 \cdot 10^{-5}$
$C_{\mu\mu}^{S LL}$	$2.3 \cdot 10^{-7}$	$7.2 \cdot 10^{-8}$	$4.6 \cdot 10^{-6}$	$3.3 \cdot 10^{-7}$	$7.1 \cdot 10^{-6}$	$1.0 \cdot 10^{-7}$
$C_{\tau\tau}^{S LL}$	$1.2 \cdot 10^{-6}$	$3.7 \cdot 10^{-7}$	$2.4 \cdot 10^{-5}$	$1.7 \cdot 10^{-6}$	$2.4 \cdot 10^{-5}$	$3.5 \cdot 10^{-7}$
$C_{\tau\tau}^{T LL}$	$2.9 \cdot 10^{-9}$	$9.0 \cdot 10^{-10}$	$5.7 \cdot 10^{-8}$	$4.1 \cdot 10^{-9}$	$5.9 \cdot 10^{-8}$	$8.5 \cdot 10^{-10}$
$C_{bb}^{S LL}$	$2.8 \cdot 10^{-6}$	$8.6 \cdot 10^{-7}$	$5.4 \cdot 10^{-5}$	$3.8 \cdot 10^{-6}$	$9.0 \cdot 10^{-7}$	$1.2 \cdot 10^{-8}$
$C_{bb}^{T LL}$	$2.1 \cdot 10^{-9}$	$6.4 \cdot 10^{-10}$	$4.1 \cdot 10^{-8}$	$2.9 \cdot 10^{-9}$	$4.2 \cdot 10^{-8}$	$6.0 \cdot 10^{-10}$
$C_{ee}^{V RR}$	$3.0 \cdot 10^{-5}$	$9.4 \cdot 10^{-6}$	$2.1 \cdot 10^{-7}$	$1.5 \cdot 10^{-8}$	$2.1 \cdot 10^{-6}$	$3.5 \cdot 10^{-8}$
$C_{\mu\mu}^{V RR}$	$3.0 \cdot 10^{-5}$	$9.4 \cdot 10^{-6}$	$1.6 \cdot 10^{-5}$	$1.1 \cdot 10^{-6}$	$2.1 \cdot 10^{-6}$	$3.5 \cdot 10^{-8}$
$C_{\tau\tau}^{V RR}$	$1.0 \cdot 10^{-4}$	$3.2 \cdot 10^{-5}$	$5.3 \cdot 10^{-5}$	$3.8 \cdot 10^{-6}$	$4.8 \cdot 10^{-6}$	$7.9 \cdot 10^{-8}$
$C_{bb}^{V RR}$	$3.5 \cdot 10^{-4}$	$1.1 \cdot 10^{-4}$	$6.7 \cdot 10^{-5}$	$4.8 \cdot 10^{-6}$	$6.0 \cdot 10^{-6}$	$1.0 \cdot 10^{-7}$
C_{bb}^{RA}	$4.2 \cdot 10^{-4}$	$1.3 \cdot 10^{-4}$	$6.5 \cdot 10^{-3}$	$4.6 \cdot 10^{-4}$	$1.3 \cdot 10^{-3}$	$2.2 \cdot 10^{-5}$
C_{bb}^{RV}	$2.1 \cdot 10^{-3}$	$6.4 \cdot 10^{-4}$	$6.7 \cdot 10^{-5}$	$4.7 \cdot 10^{-6}$	$6.0 \cdot 10^{-6}$	$1.0 \cdot 10^{-7}$

\Rightarrow BR($\mu \rightarrow e \gamma$) depends on dipole C_D
(but mixing effects from RGE running and loop contributions render it also sensitive to scalar/tensor/vector contributions, even for $q\bar{q}$ operators)

Unexpected findings!

► Include as many **observables & operators as possible!**

(e.g. $\mu e \gamma \gamma$ contact interactions, angular observables in polarised $\mu \rightarrow 3e$ decays, ...)

[Davidson et al, 2007.09612]

[Bolton, Petcov, 2204.03468]

Muon cLFV: EFT approach to New Physics

Recent (novel) **EFT** approach to **muon transitions**:

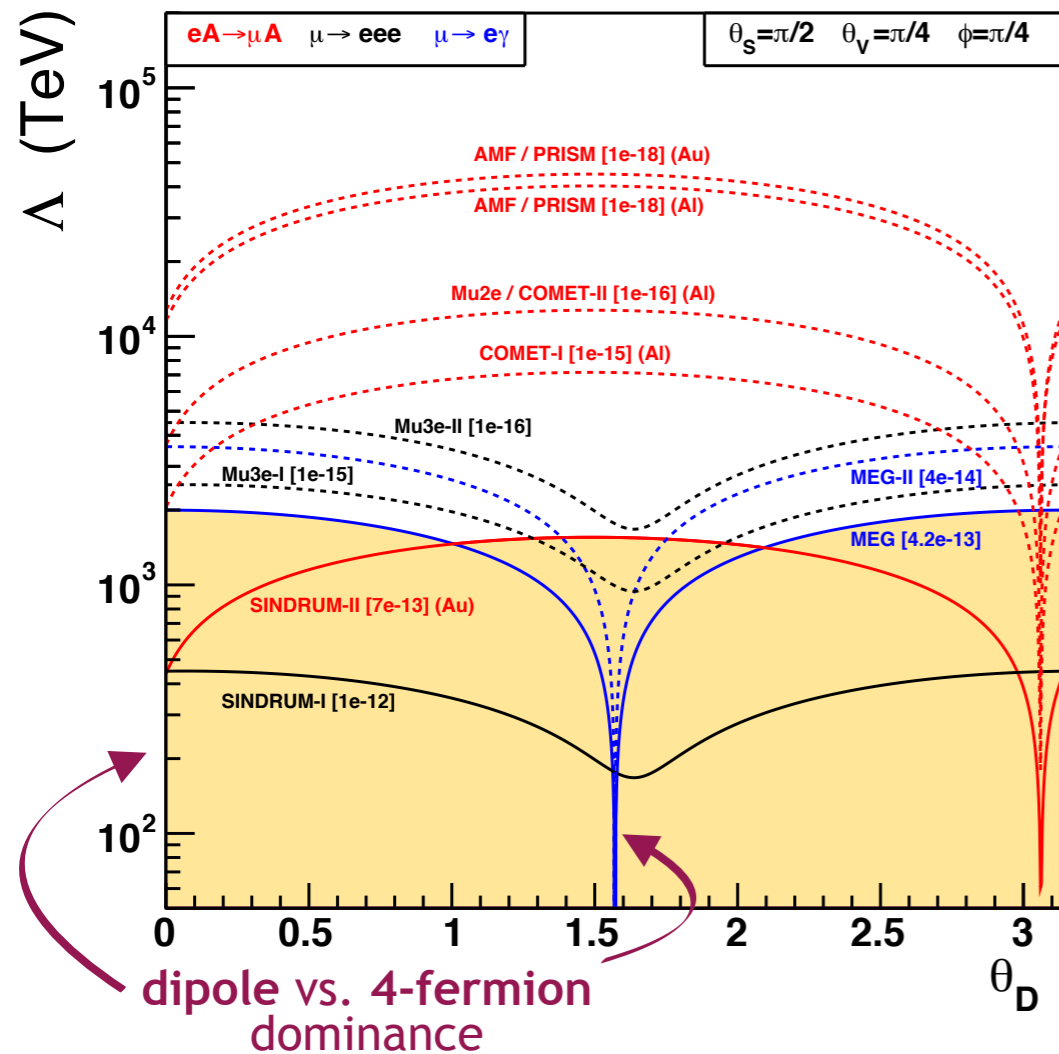
$$\mathcal{L}_{\text{NP, cLFV}}^{\text{eff}} = \frac{1}{\Lambda^2} \left[C_D (\bar{e} \sigma^{\nu\rho} P_R \mu) F_{\nu\rho} + C_S (\bar{e} P_R \mu) (\bar{e} P_R e) + C_{VR} (\bar{e} \gamma^\nu P_L \mu) (\bar{e} \gamma_\nu P_R e) + C_{VL} (\bar{e} \gamma^\nu P_L \mu) (\bar{e} \gamma_\nu P_L e) + C_{\text{N-light}} \mathcal{O}_{\text{N-light}} + C_{\text{N-heavy}\perp} \mathcal{O}_{\text{N-heavy}\perp} \right]$$

$$\vec{C} = \{C_D, C_S, C_{VR}, C_{VR}, C_{VL}, C_{\text{N-light}}, C_{\text{N-heavy}\perp}\}$$

$$\Rightarrow \text{BR}(\mu \rightarrow e\gamma) \simeq 384\pi^2 \frac{v^4}{\Lambda^4} |\vec{C} \cdot \hat{e}_{DR}|^2 \rightsquigarrow \leq \text{BR}^{\text{exp}}(\text{future})$$

and likewise for other observables

$\text{BR}(\mu \rightarrow 3e)$, $\text{CR}(\mu - e, N)$, Muonium oscillations...



$$2\sqrt{2} C_D \approx \frac{\cos \theta_D}{\Lambda^2}$$

Sensitivity to NP scales (current & future):

MEG ($\mu \rightarrow e\gamma$) $\leftrightarrow \Lambda_{\text{cLFV}} \sim \mathcal{O}(10^3 \text{ TeV})$ [dipole]
steadily improved by Mu3e $\sim \mathcal{O}(5 \times 10^3 \text{ TeV})$

SINDRUM II ($\mu - e$, Au) $\leftrightarrow \Lambda_{\text{cLFV}} \sim \mathcal{O}(10^3 \text{ TeV})$ [4f]

Mu2e/COMET II ($\mu - e$, Al) $\leftrightarrow \Lambda_{\text{cLFV}} \lesssim \mathcal{O}(10^4 \text{ TeV})$
[either dipole or 4f]

⇒ **cLFV** data to constrain \mathcal{O}^6 (and infer sensitivity of a process to operator \mathcal{O}^6)

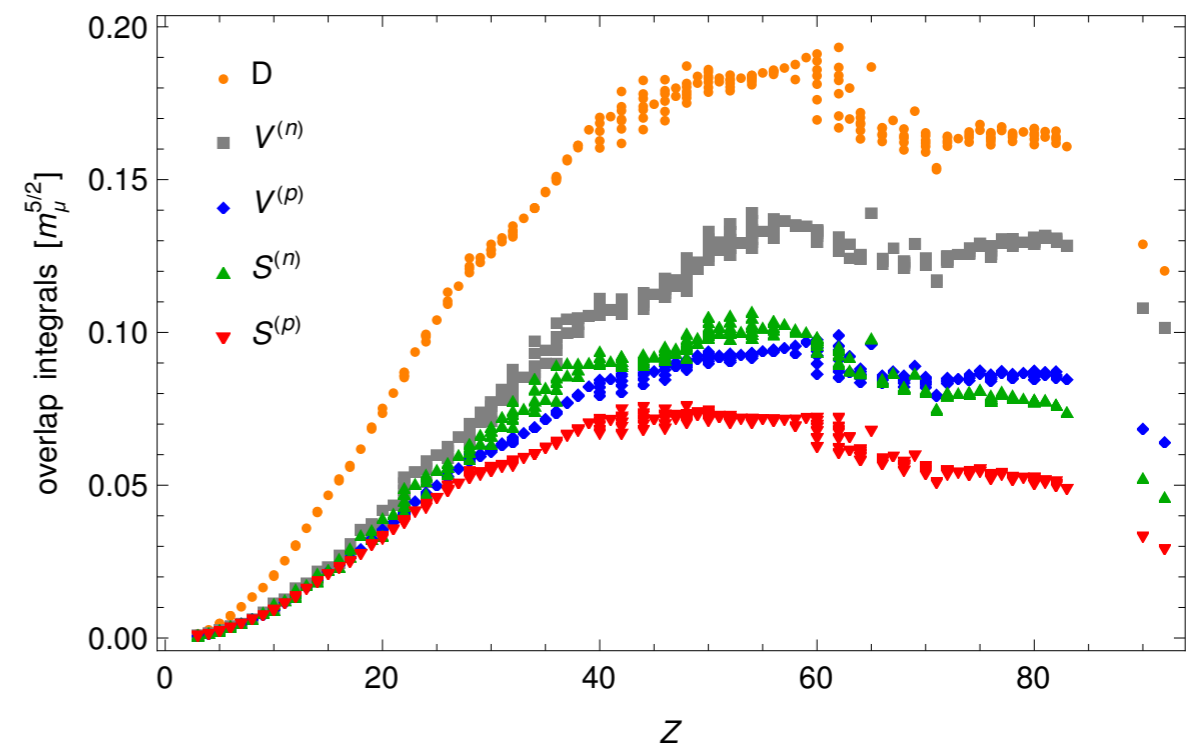
► Fully exploring the potential of atomic (elastic) **muon-electron conversion**, $CR(\mu - e, N)$:

Comparatively more involved theoretical approach!

Explore target-nucleus dependence to distinguish dominant operator (hint on NP model!)

[extensive contributions since Kitano et al, 0203110! see Davidson et al, 1810.01884; Heeck et al, 2203.00702, ...; Haxton et al, 2406.12818]

$$BR_{SI}(\mu A \rightarrow eA) = \frac{32G_F^2}{\Gamma_{\text{capture}}} \left[\left| C_{V,R}^{pp} V^{(p)} + C_{S,L}^{pp'} S^{(p)} + C_{V,R}^{nn} V^{(n)} + C_{S,L}^{nn'} S^{(n)} + C_{D,L} \frac{D}{4} \right|^2 + \{L \leftrightarrow R\} \right].$$



[Heeck et al, 2203.00702]

Overlap integrals:

more distinguishable at **large Z** !

Better disentangle dominant NP contributions... but not "experimentally" feasible...

Muon cLFV: EFT approach & conversion in nuclei



IN2P3
Les deux infinis

⇒ cLFV data to constrain \mathcal{O}^6 (and infer sensitivity of a process to operator \mathcal{O}^6)

► Fully exploring the potential of atomic (elastic) muon-electron conversion, $CR(\mu - e, N)$:

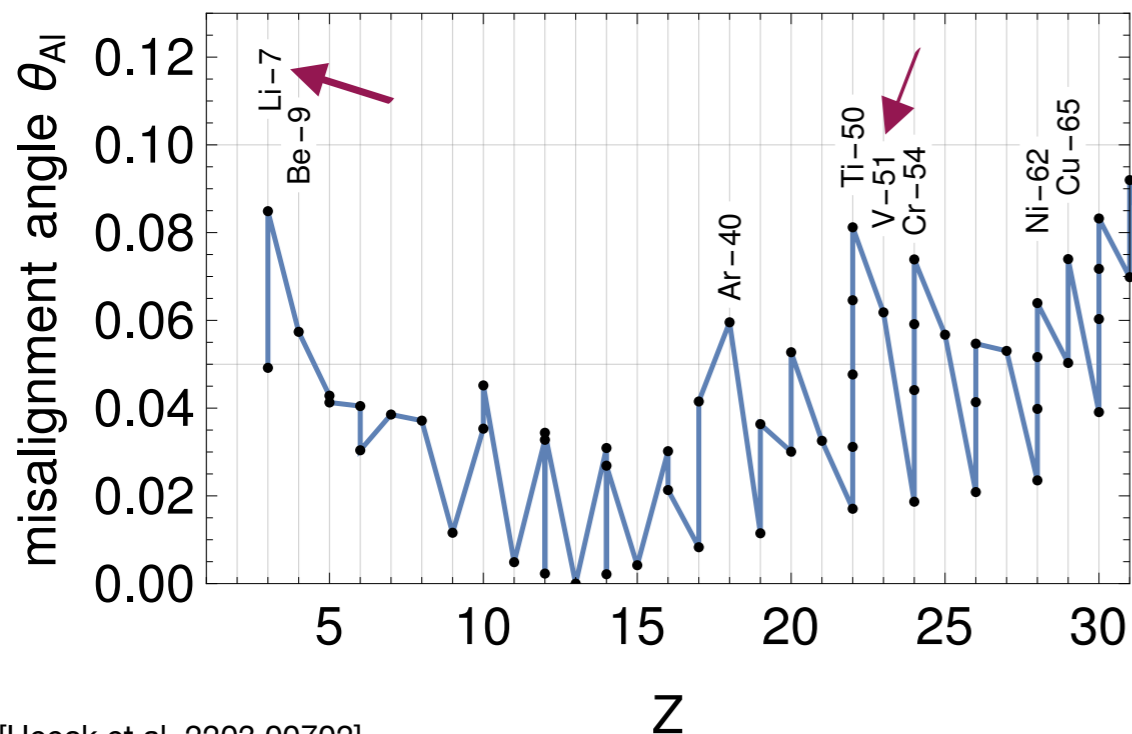
Comparatively more involved theoretical approach!

Explore target-nucleus dependence to distinguish dominant operator (hint on NP model!)

[extensive contributions since Kitano et al, 0203110! see Davidson et al, 1810.01884; Heeck et al, 2203.00702, ...; Haxton et al, 2406.12818]

In the advent of an observation (@ Mu2e, COMET \leadsto using Aluminium targets)

prepare choice of future targets! Largest complementarity with respect to Al? θ_{Al}



[Heeck et al, 2203.00702]

- Heavier nuclei (Au, Pb)! ... not feasible... (pulsed beams)
- Among experimental-friendly $Z \leq 25$ targets several (theoretically good) candidates
Li-7, Ti-50, Ti-49, Cr-54, .., V-51

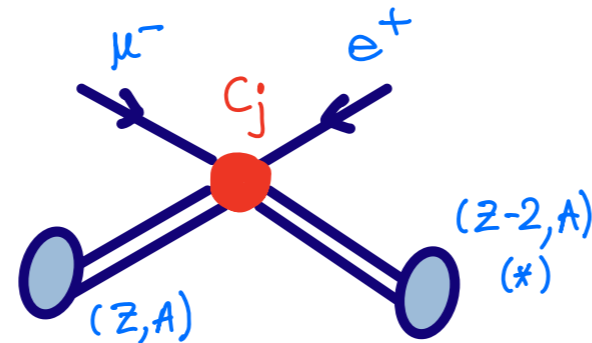
⇒ Li-7 and/or V-51 : preferable "second" targets post $CR(\mu - e, Al)$ observation

⇒ cLFV data to constrain \mathcal{O}^6 (and infer sensitivity of a process to operator \mathcal{O}^6)

► Fully exploring the potential of atomic (elastic) muon-electron conversion, $CR(\mu - e, N)$:

And of its lepton number violating counterpart, $\mu^- + (A, Z) \rightarrow e^+ + (A, Z - 2)^{(*)}$

A unique connection between LNV (in association with Majorana nature and possibly, neutrino mass generation) and cLFV



$\mu^- - e^-$ conversion: coherent process, single nucleon, nuclear ground states
 $\mu^- - e^+$ conversion: 2 nucleons ($\Delta Q = 2$), possibly excited final state

From a theoretical point of view, not straightforward!

- Higher-dimension operators in \mathcal{L}^{eff} (dim 6, 10, 14...)
- Nuclear matrix elements extremely hard to compute!

$$\Gamma_{\mu e}^{\text{LNV}} \approx \frac{G_F^4 g_A^4}{32\pi^2} |\epsilon_{C_j}^2| \frac{m_e^2 m_\mu^2}{R^2} |F(Z-2, E_e)| \langle \phi_\mu \rangle^2 |\mathcal{M}^{(\mu^-, e^+)}|^2$$

(only two $\mathcal{M}^{(\mu^-, e^+)}$ known, for Ti-48...)

[Domin et al, 0409033; Simkovic et al, 0103029]

⇒ Very hard to draw implications... **Must tackle NME!**

Tau cLFV: (semi-) leptonic modes

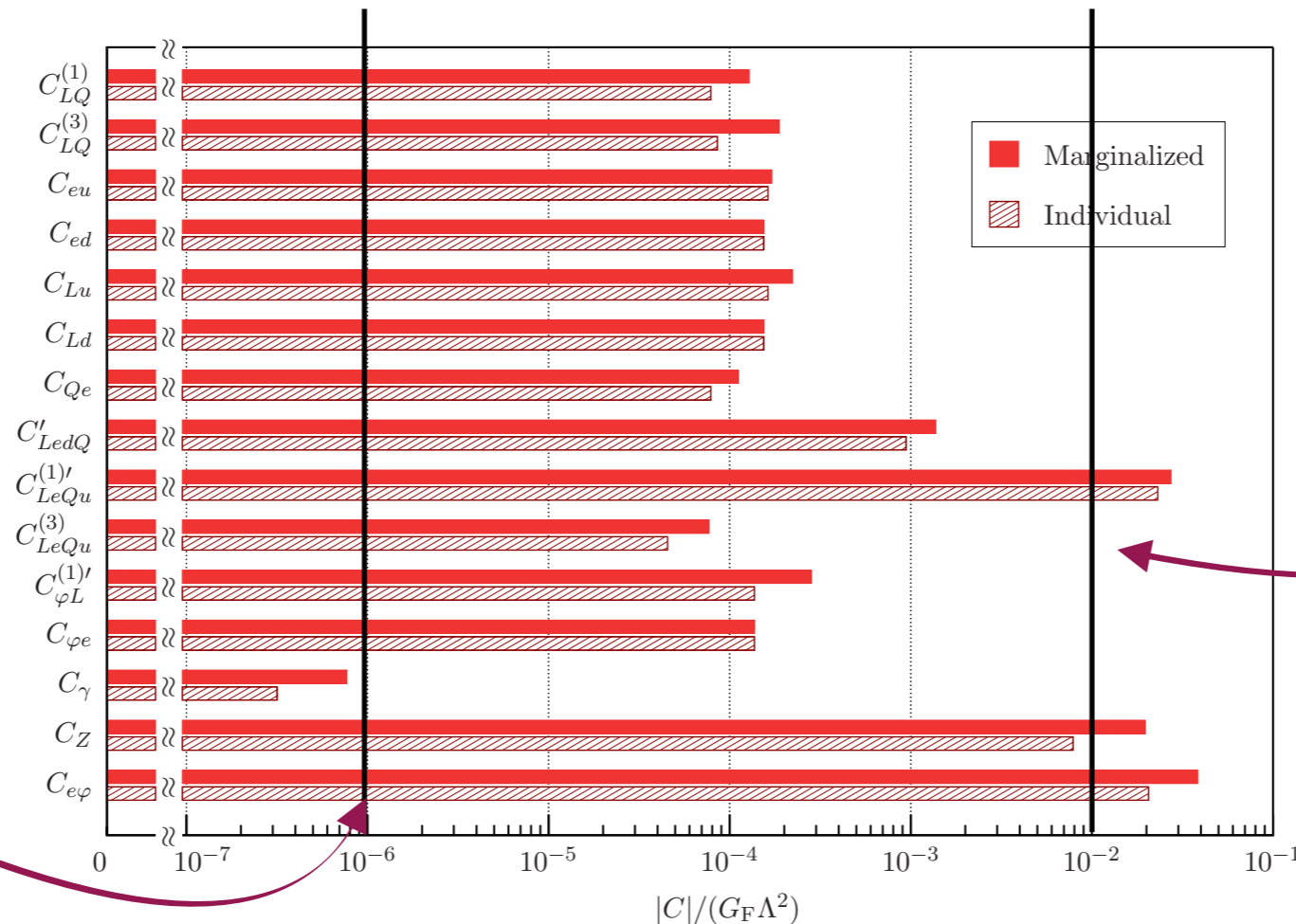
Flavour violating tau decays: comparatively large number of modes

Leptonic (radiative, three-body) as well as semi-leptonic (light mesons, 2- and 3-body)

⇒ theoretically much more involved (scales, hadronisation, ...)

⇒ larger set of (tree-level) contributing operators (e.g. numerous $qq\ell\ell$, gluon, ...)!

⇒ more challenging to disentangle operator dominance... (even @ tree level!)



For $C \approx 1$,
 $\Lambda_{NP} \sim 300 \text{ TeV}$

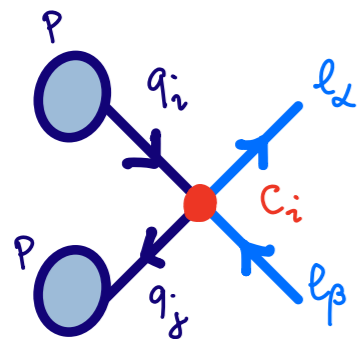
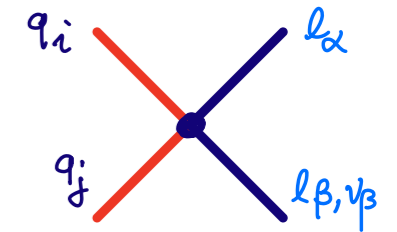
For $C \approx 1$,
 $\Lambda_{NP} \sim 3 \text{ TeV}$

[Banerjee et al, 2203.14919]

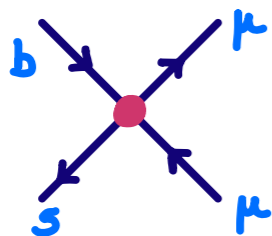
Overview of Belle II limits on relevant coefficients (and NP scales) for cLFV tau decays

- Albeit leading to formally different transitions, the same leptonic and semi-leptonic operators can be at the origin of **flavour violating transitions** in very distinct contexts

Consider a 4-fermion quark-lepton operator $(q_i q_j \ell_\alpha \ell_\beta)$, with $i = j, \alpha \neq \beta$
 One operator can source **rare LHC cLFV decays** (rich "flavour" content!),



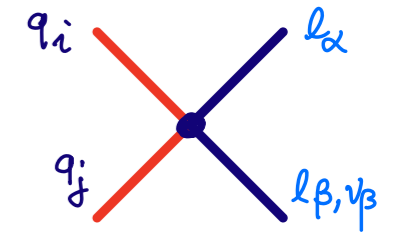
$$pp \rightarrow \ell_\alpha \ell_\beta$$



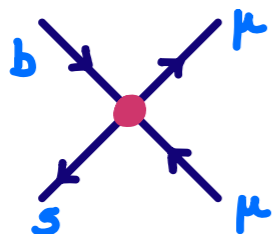
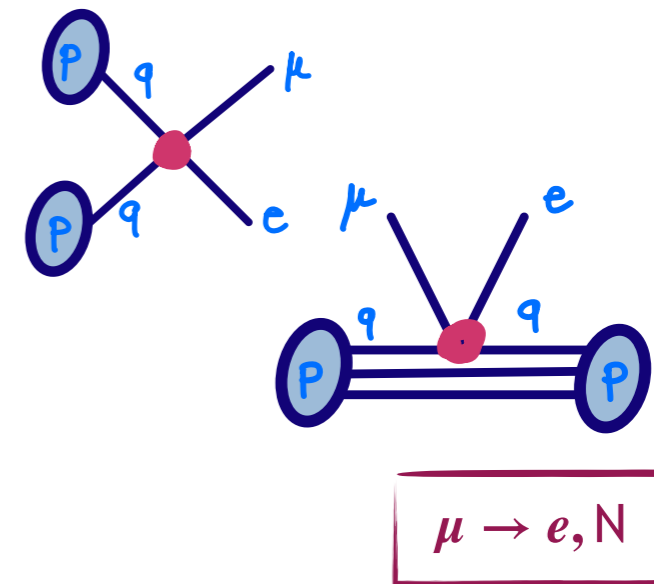
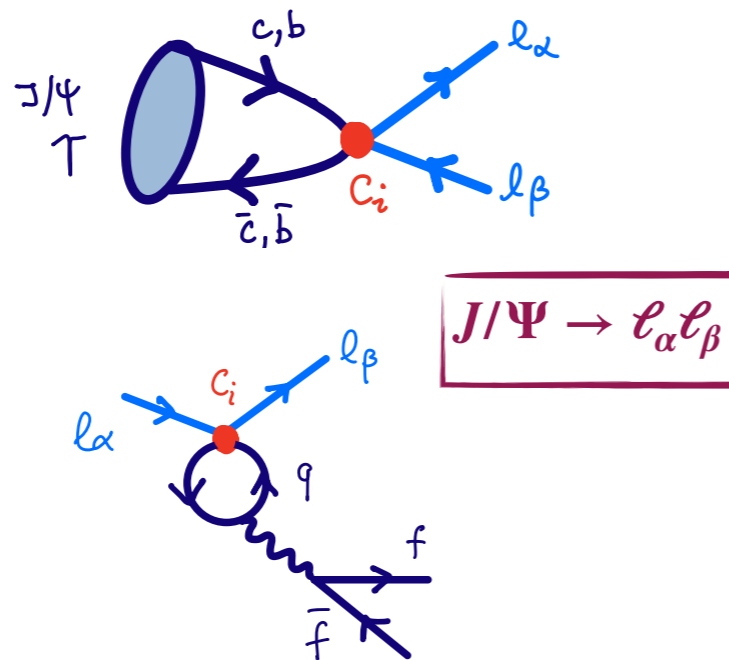
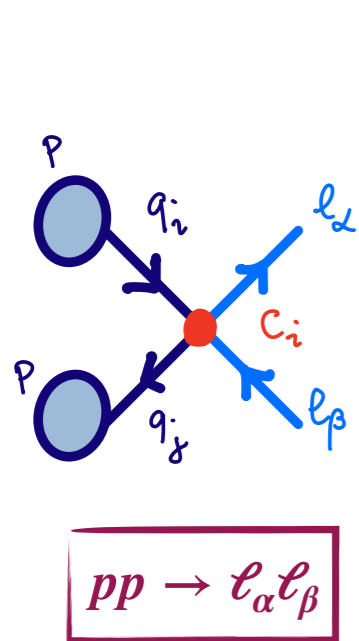
cLFV at higher-energies: spinning operators

- ▶ Albeit leading to formally different transitions, the same leptonic and semi-leptonic operators can be at the origin of **flavour violating transitions** in very distinct contexts

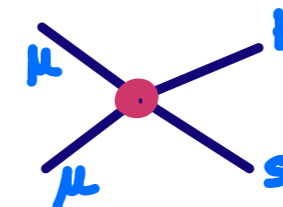
Consider a 4-fermion quark-lepton operator $(q_i q_j \ell_\alpha \ell_\beta)$, with $i = j, \alpha \neq \beta$
 One operator can source **rare LHC cLFV decays** (rich "flavour" content!),
cLFV semileptonic decays, muon-electron conversion, ...



[recent review, see Fernandez-Martinez et al, 2403.09772]



$\Rightarrow b \rightarrow s \ell \ell$ at a $\mu^+ \mu^-$ collider



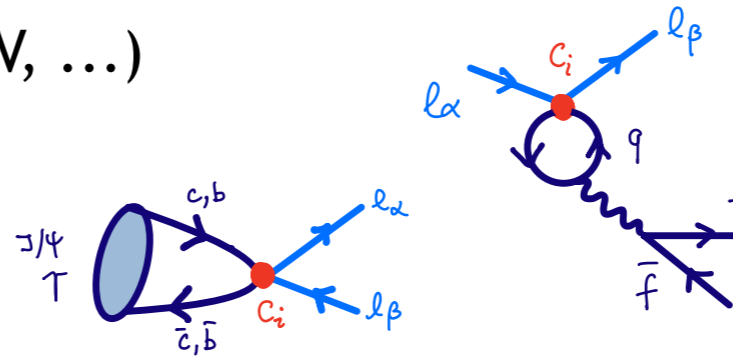
Meson decays: cLFV and more

► **Meson decays:** excellent hunting grounds for "leptophilic" New Physics

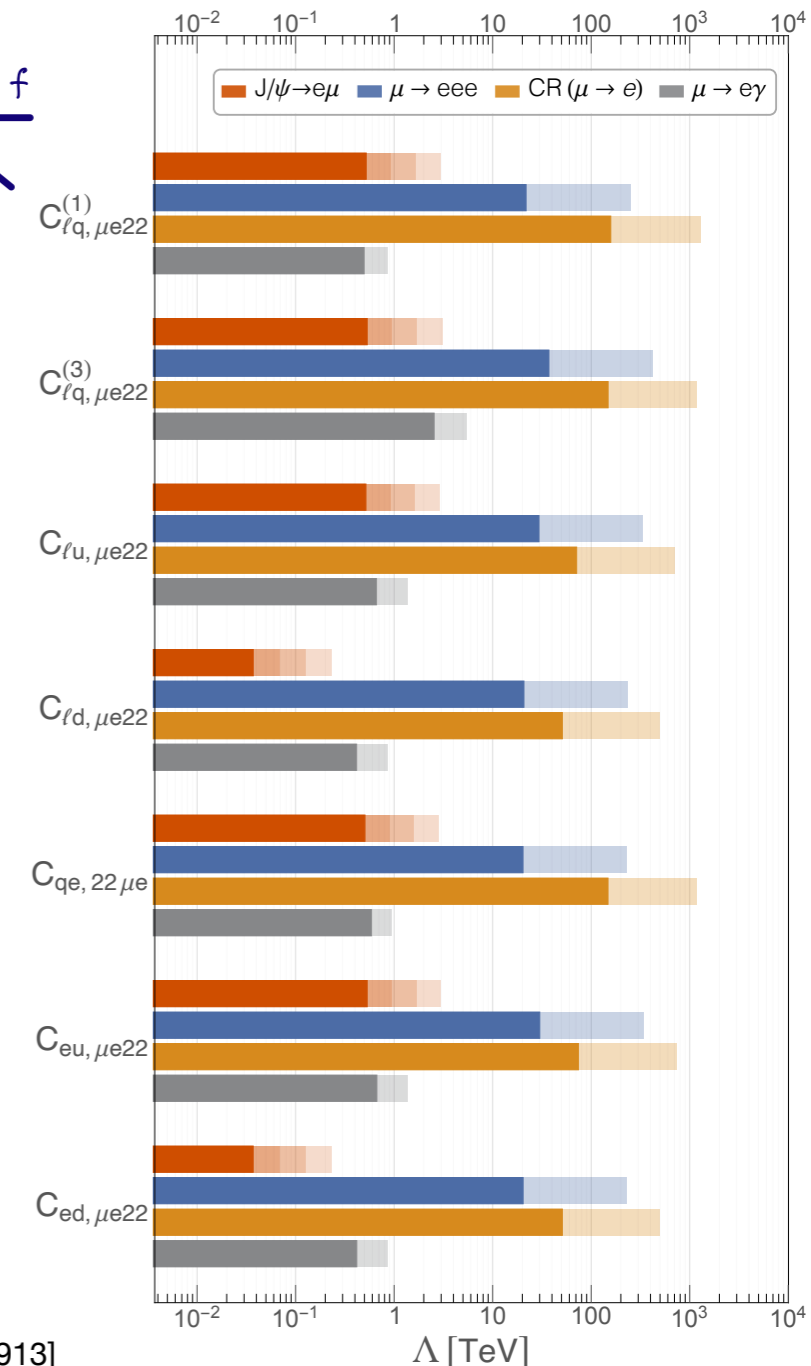
⇒ deviations from SM (lepton flavour universality violation, angular distributions, ...)

⇒ new phenomena (cLFV, LNV, ...)

► **cLFV semileptonic meson decays!**



2q2l operators			
$\mathcal{O}_{lq,prst}^{(1)}$	$(\bar{L}_p \gamma_\mu L_r)(\bar{Q}_s \gamma^\mu Q_t)$	$\mathcal{O}_{lq,prst}^{(3)}$	$(\bar{L}_p \gamma_\mu \tau^I L_r)(\bar{Q}_s \gamma^\mu \tau^I Q_t)$
$\mathcal{O}_{lu,prst}$	$(\bar{L}_p \gamma_\mu L_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{ld,prst}$	$(\bar{L}_p \gamma_\mu L_r)(\bar{d}_s \gamma^\mu d_t)$
$\mathcal{O}_{eu,prst}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{ed,prst}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$
$\mathcal{O}_{qe,prst}$	$(\bar{Q}_p \gamma^\mu Q_r)(\bar{e}_s \gamma_\mu e_t)$	$\mathcal{O}_{ledq,prst}$	$(\bar{L}_p e_r)(\bar{d}_s Q_t)$
$\mathcal{O}_{lequ,prst}^{(1)}$	$(\bar{L}_p^a e_r) \epsilon_{ab} (\bar{Q}_s^b u_t)$	$\mathcal{O}_{lequ,prst}^{(3)}$	$(\bar{L}_p^a \sigma_{\mu\nu} e_r) \epsilon_{ab} (\bar{Q}_s^b \sigma^{\mu\nu} u_t)$
4l operators		Dipole operators	
$\mathcal{O}_{ll,prst}$	$(\bar{L}_p \gamma_\mu L_r)(\bar{L}_s \gamma^\mu L_t)$	$\mathcal{O}_{eW,pr}$	$(\bar{L}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$
$\mathcal{O}_{ee,prst}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$\mathcal{O}_{eB,pr}$	$(\bar{L}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$
$\mathcal{O}_{le,prst}$	$(\bar{L}_p \gamma_\mu L_r)(\bar{e}_s \gamma^\mu e_t)$		
Lepton-Higgs operators			
$\mathcal{O}_{\varphi l,pr}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{L}_p \gamma^\mu L_r)$	$\mathcal{O}_{\varphi l,pr}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{L}_p \gamma^\mu \tau^I L_r)$
$\mathcal{O}_{\varphi e,pr}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$	$\mathcal{O}_{\varphi 3,pr}$	$(\bar{L}_p e_r \varphi)(\varphi^\dagger \varphi)$



Comparative study of the **probing power of quarkonium** (charmonium) **$\mu - e$ cLFV decays** for relevant $\mathcal{C}_{\mu e} = 1$

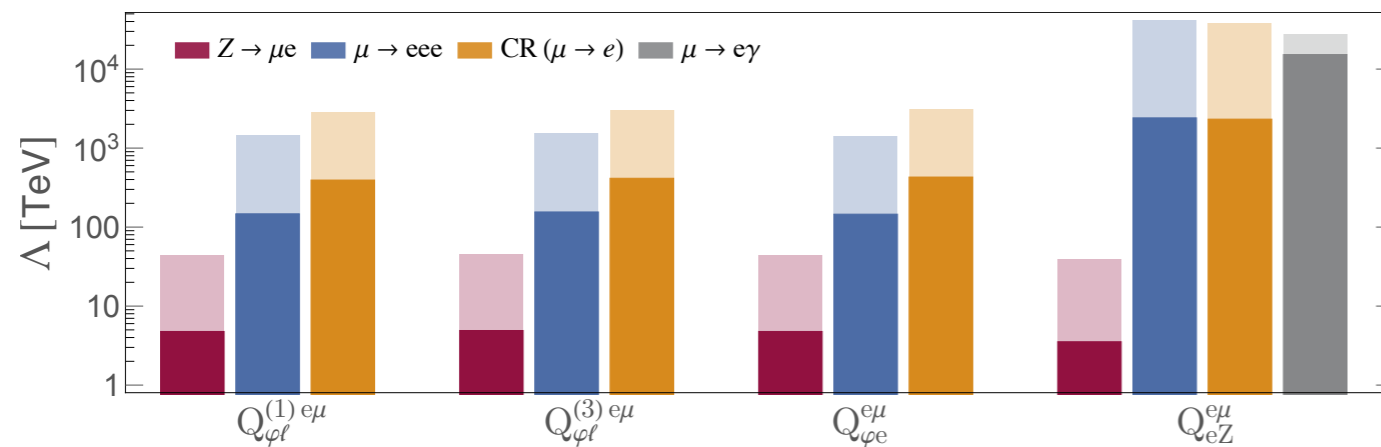
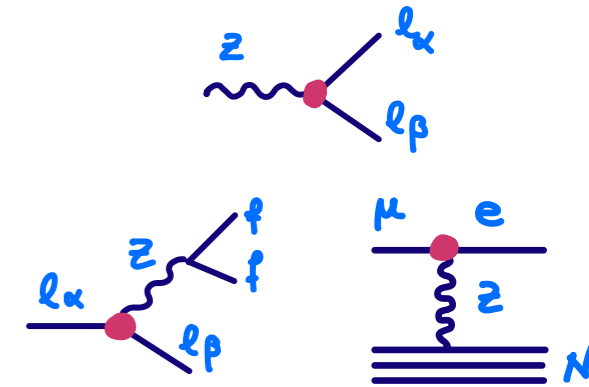
[recent study - Calibbi et al, 2207.10913]

Lepton flavours @ high Tera-Z

High-energy colliders: also **high-intensity frontier** (amazing luminosities!)

LHC \rightsquigarrow abundant sources of flavour in pp collisions
(and also a Higgs-factory...)

TeraZ factory (FCC-ee, CEPC) \rightsquigarrow **EW precision & flavour violation**



TeraZ factory \rightsquigarrow **cLFV Z decays**

[Calibbi et al, 2107.10273]

For $Z \rightarrow \mu e$ better sensitivity of **dedicated (low-energy) cLFV searches**

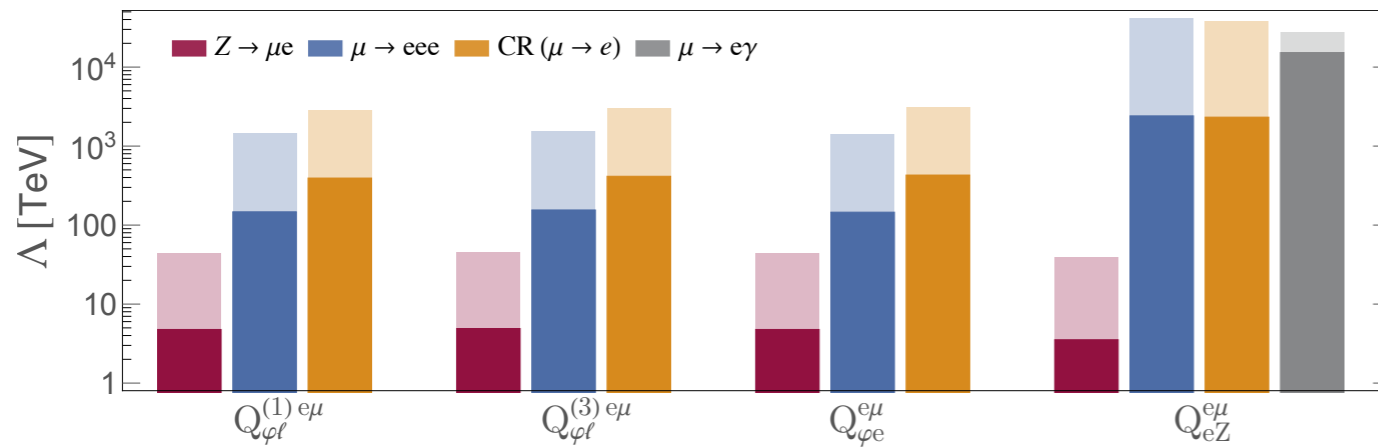
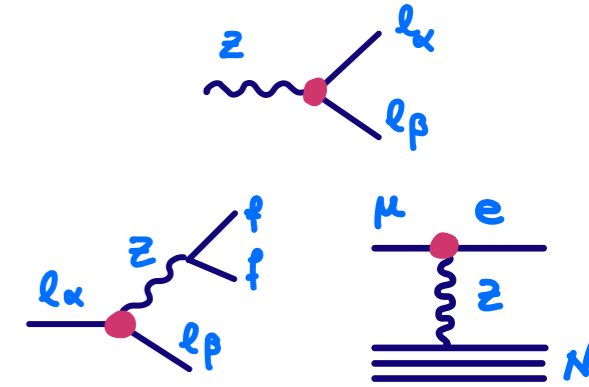
$\mu \rightarrow eee$, $\mu - e$ conversion

Lepton flavours @ high Tera-Z

High-energy colliders: also **high-intensity frontier** (amazing luminosities!)

LHC \rightsquigarrow abundant sources of flavour in pp collisions
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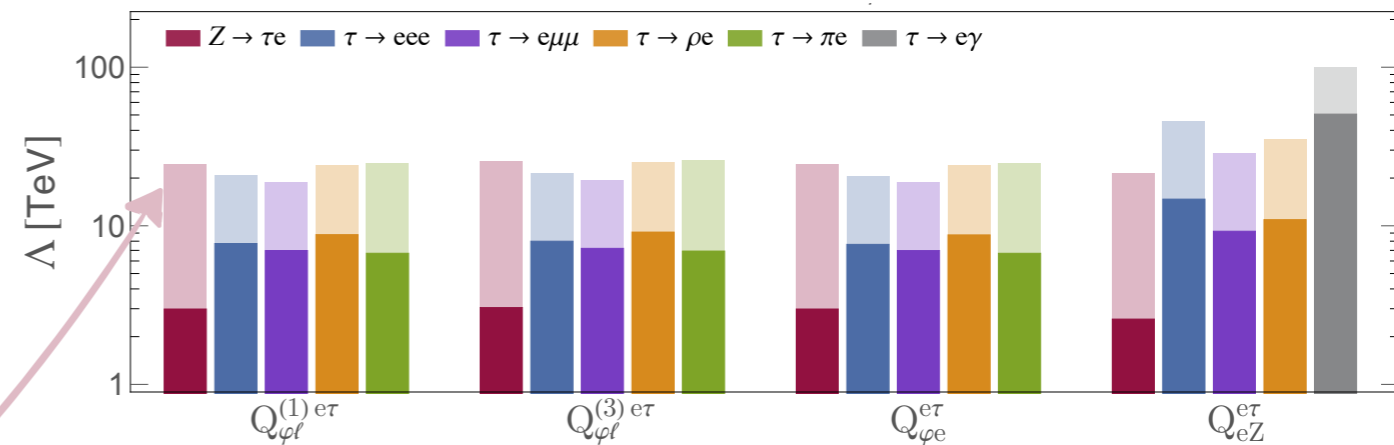
TeraZ factory (FCC-ee, CEPC) \rightsquigarrow EW precision & flavour violation



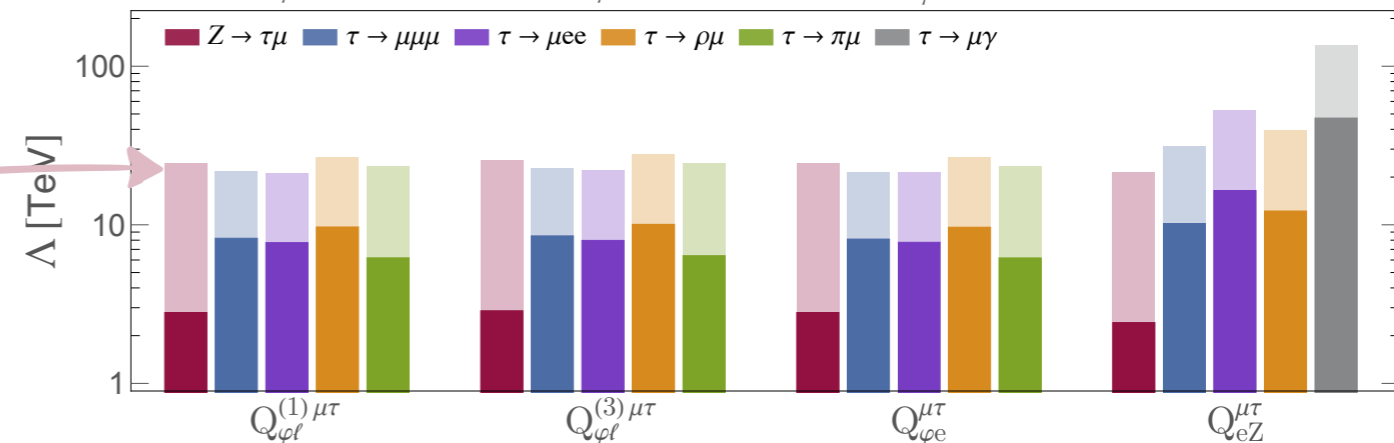
TeraZ factory \rightsquigarrow cLFV Z decays

[Calibbi et al, 2107.10273]

For $Z \rightarrow \mu e$ better sensitivity of dedicated (low-energy) cLFV searches



Promising potential of TeraZ factory to probe NP at the origin of $Z \rightarrow \tau \ell$ decays (competitive with low-energy cLFV)

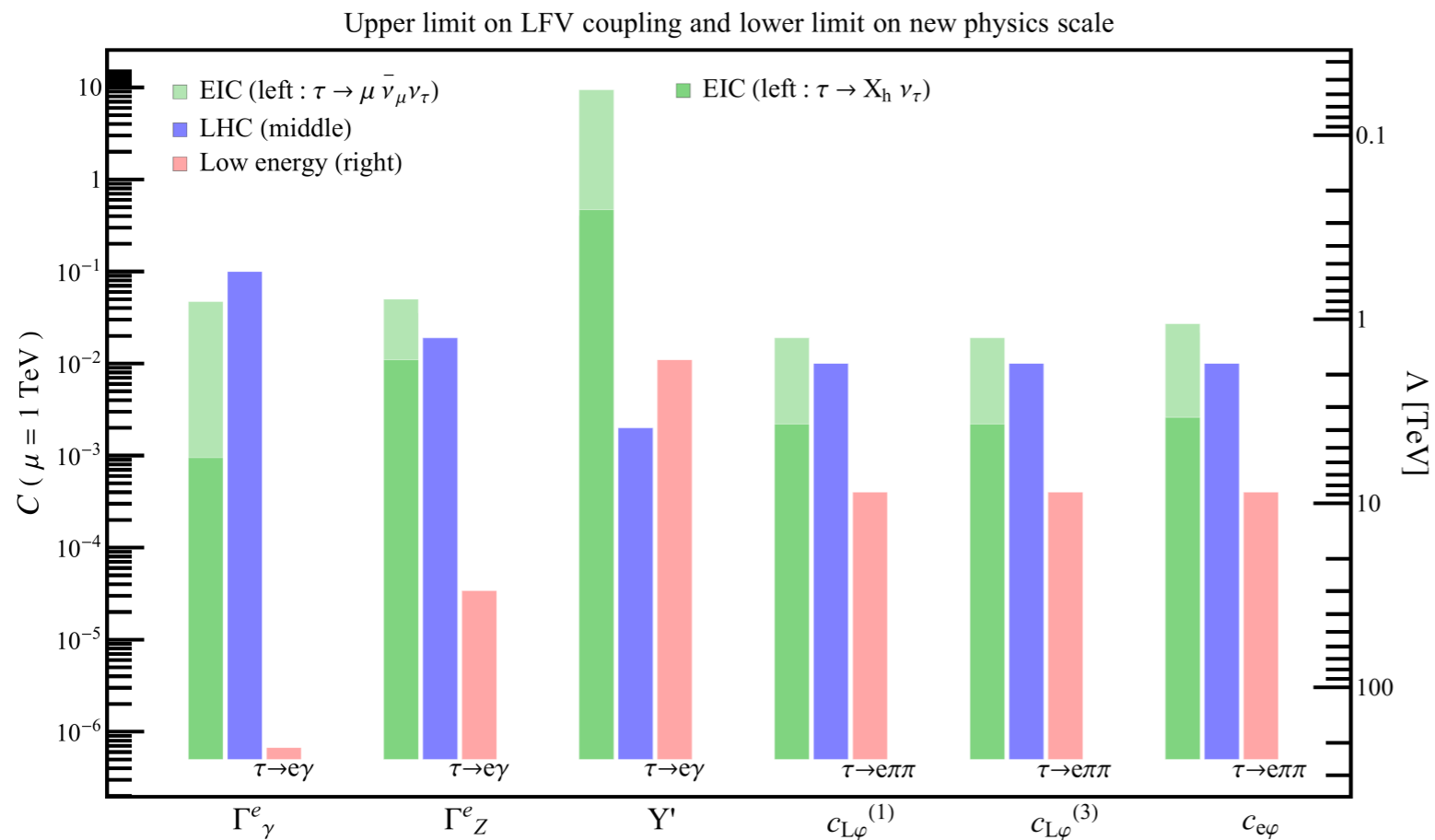


And at electron-ion colliders (EIC)

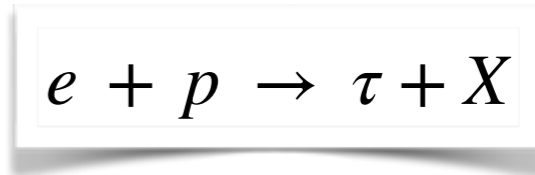
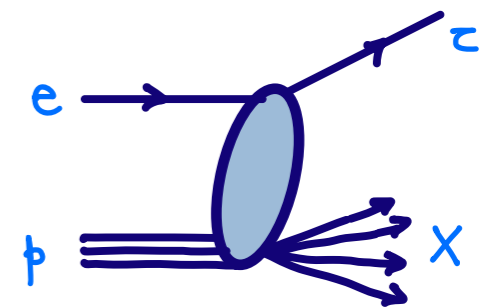
► **Electron-Ion colliders** also offer opportunities to study cLFV!

In general, less ambitious probing power (compared to **LHC**, and especially to dedicated **low-energy** experiments)

For $C \approx 1$, $\Lambda_{\text{NP}} \sim \mathcal{O}(\text{few TeV})$



$$\frac{4 G_F C}{\sqrt{2}} \equiv \frac{1}{\Lambda_{\text{NP}}}$$



[Cirigliano et al, 2102.106176]

After EFT - New Physics models of cLFV!

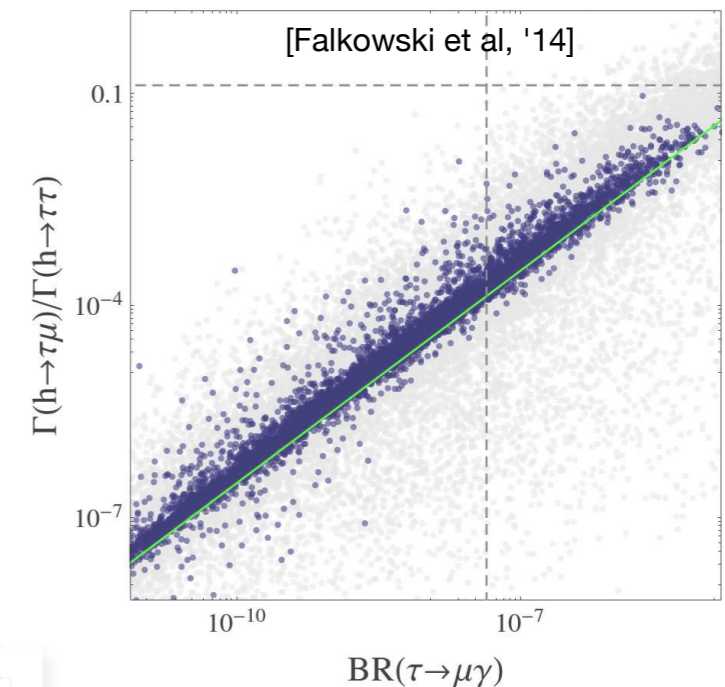
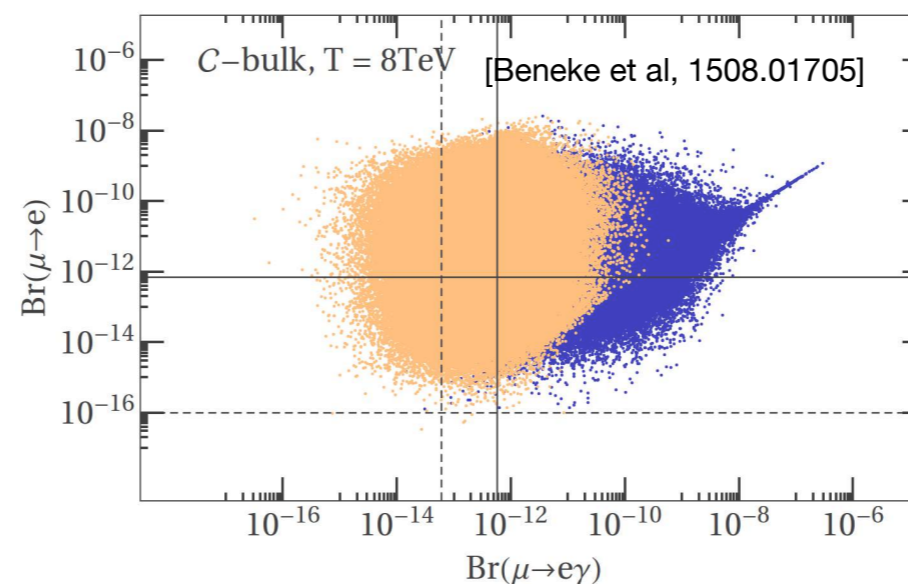
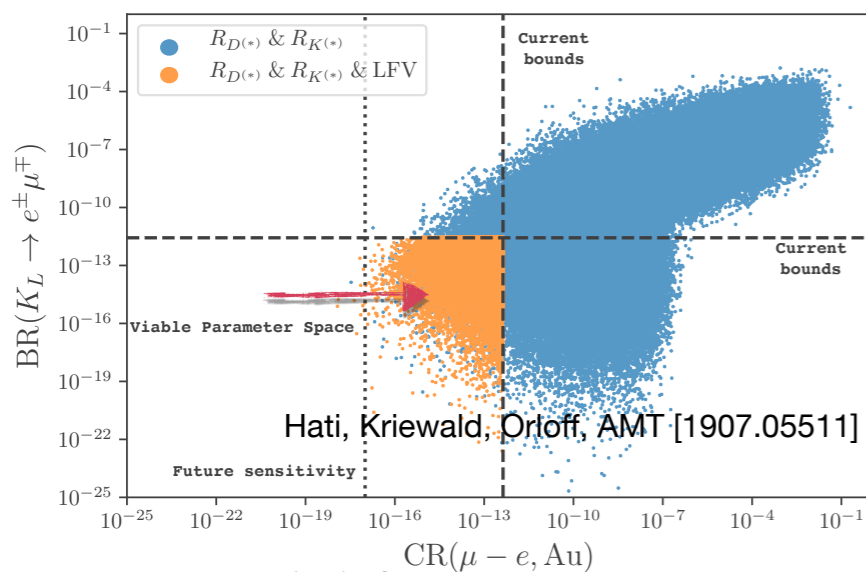
- ▶ **Effective approach:** *first characterisation* (mostly constraints & hints) on scale of NP and nature of new interactions (couplings and currents)
 - ⇒ Ultimately we do *need to unveil the model of NP at work!*

- ▶ Although **oscillation data** (massive neutrinos) do **imply cLFV** (*direct consequence*), cLFV can be **independent** of mechanism of **neutrino mass generation**

Supersymmetry: unconstrained models (beyond cMSSM, pMSSM), new sources of LFV

Rp-violating SUSY

Leptoquark models (extended field content); **Extra-dimensions**; **extended Higgs sectors**, ...



cLFV : powerful means to test/falsify models of NP

After EFT - New Physics models of cLFV!



IN2P3
Les deux infinis

- ▶ **Effective approach:** *first characterisation* (mostly constraints & hints) on scale of NP and nature of new interactions (couplings and currents)
⇒ Ultimately we do *need to unveil the model of NP at work!*
- ▶ Although **oscillation data** (massive neutrinos) do **imply cLFV** (direct consequence), **cLFV contributions** fully **independent** of mechanism of **neutrino mass generation**

- ▶ Here: focus on **SM extensions** aiming at addressing **neutrino mass generation!**

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \frac{\mathcal{C}_5 \mathcal{O}^5}{\Lambda_{\text{LNV}}}(m_\nu) + \frac{\mathcal{C}_6 \mathcal{O}^6}{\Lambda_{\text{cLFV}}^2}(\ell_\alpha \leftrightarrow \ell_\beta) + \dots + \frac{\mathcal{C}_7 \mathcal{O}^7}{\Lambda_{\text{LNV}}^3}(0\nu 2\beta) + \dots$$

Majorana ν masses cLFV and many others LNV (& cLFV), ...

Not always trivial to establish connection between \mathcal{C}_5 , \mathcal{C}_6 (and \mathcal{C}_7)

cLFV : powerful means to test/falsify models of NP (m_ν) (examples ahead!)

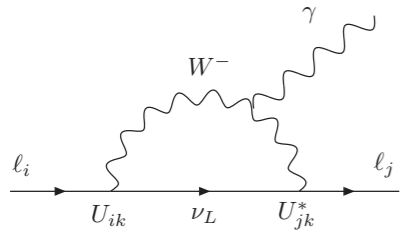
New Physics paths to cLFV: seesaw models of neutrino mass generation



Neutrino masses (brief "how to"...)

Most minimal possibility: SM extended by **Dirac RH neutrinos** (impose **L conservation**)

$$\Rightarrow \mathcal{L}_{m_\nu} \sim - Y^\nu L \tilde{H} \nu_R \text{ but tiny Yukawa couplings, } \mathcal{O}(10^{-13})$$



Successfully **accounting for oscillation data...**

No impact for cLFV; GIM-like suppression due to smallness of m_{ν_i}

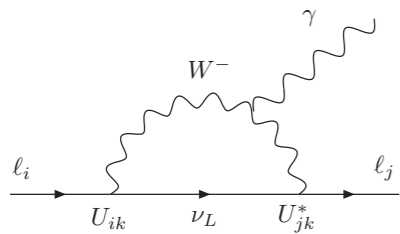
$$\text{BR}(\mu \rightarrow e\gamma) \sim 10^{-54} \text{ and similarly for other observables...}$$

New Physics and (Majorana) m_ν : cLFV

Neutrino masses (brief "how to"...)

Most minimal possibility: SM extended by **Dirac RH neutrinos** (impose **L conservation**)

$$\Rightarrow \mathcal{L}_{m_\nu} \sim -Y^\nu L \tilde{H} \nu_R \text{ but tiny Yukawa couplings, } \mathcal{O}(10^{-13})$$

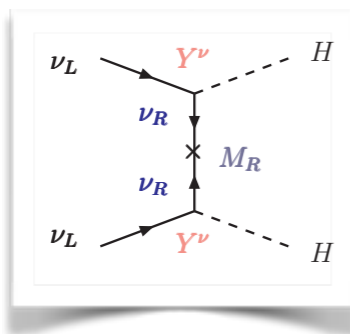


No impact for **cLFV**; **GIM-like suppression** due to smallness of m_{ν_i}

Allow for **L violation**: realisations of Weinberg operator!

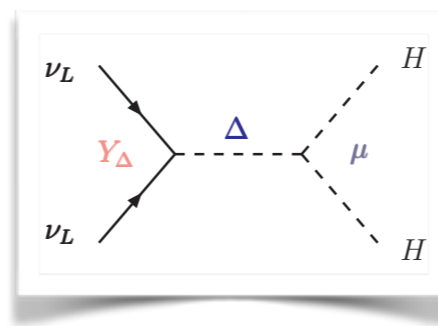
$$\mathcal{L}_{m_\nu}^5 \sim \frac{\mathcal{O}^5}{\Lambda_{\text{NP}}} (\bar{L}^c H H L)$$

Tree-level seesaw realisations



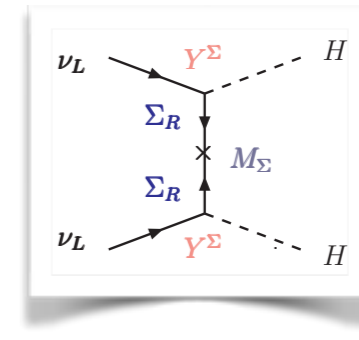
Type I (fermion singlet)

$$m_\nu \sim (Y^\nu \nu)^T \frac{1}{M_R} (Y^\nu \nu)$$



Type II (scalar triplet)

$$m_\nu \sim \frac{Y_\Delta \mu}{2} \frac{v^2}{M_\Delta^2}$$



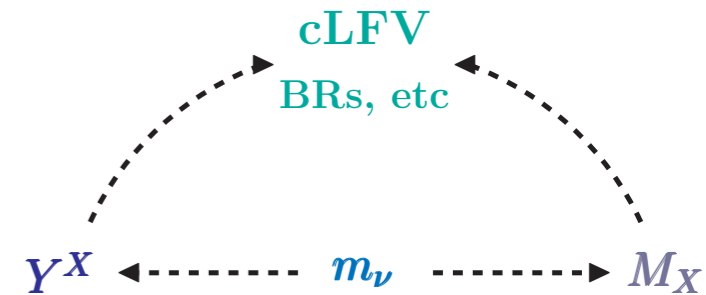
Type III (fermion triplet)

$$m_\nu \sim (Y_\Sigma \nu)^T \frac{1}{M_\Sigma} (Y_\Sigma \nu)$$

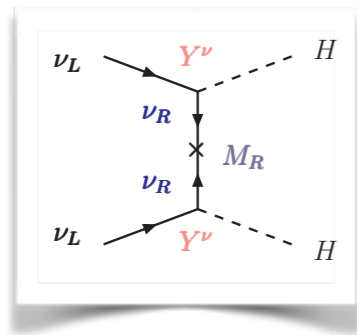
All successfully **accounting for oscillation data...** so far, *no hint from experimental searches!*

Type I seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate "balance" between sources of flavour violation and masses of new propagators
 ⇒ account for oscillation data (observation!)



- ▶ **Type I Seesaw:** extend the SM via (Majorana) sterile fermions



- ~> an enlarged spectrum
- ~> extended mixings

$$m_\nu \sim (Y^\nu \nu)^T \frac{1}{M_R} (Y^\nu \nu)$$

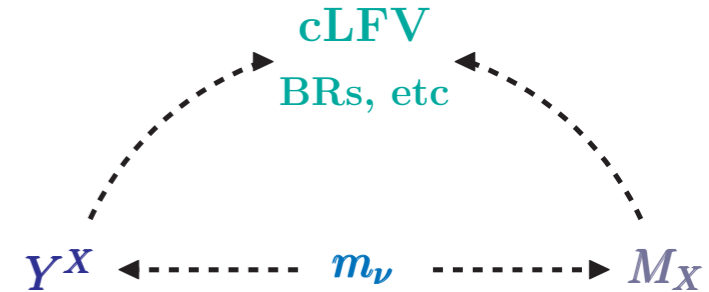
$$U^T \mathcal{M}_\nu^{6 \times 6} U = \text{diag}(m_{\nu_i}) \quad U = \begin{pmatrix} U_{\nu\nu} & U_{\nu N} \\ U_{N\nu} & U_{NN} \end{pmatrix}$$

leptonic mixing $\approx U_{\text{PMNS}}$

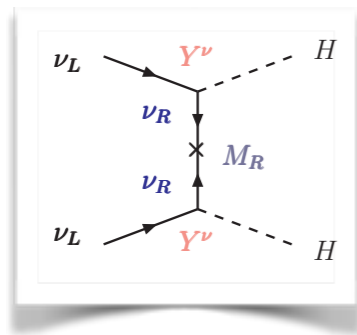
active-sterile mixings ($\theta \propto m_D^\dagger M_R^{-1}$)

Type I seesaw and cLFV

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 ⇒ account for oscillation data (observation!)



- ▶ Type I Seesaw: extend the SM via (Majorana) sterile fermions



$$\mathcal{L}_{\text{Type I}} \supset \nu Y_\nu \bar{\nu}_L \nu_R + \frac{1}{2} M_R \bar{\nu}_R^c \nu_R$$

active-sterile mixings: $\nu_L - \nu_R$

$$\Rightarrow \theta \approx \mathcal{O}(m_D^\dagger M_R^{-1})$$

$$m_\nu \sim (Y^\nu \nu)^T \frac{1}{M_R} (Y^\nu \nu)$$

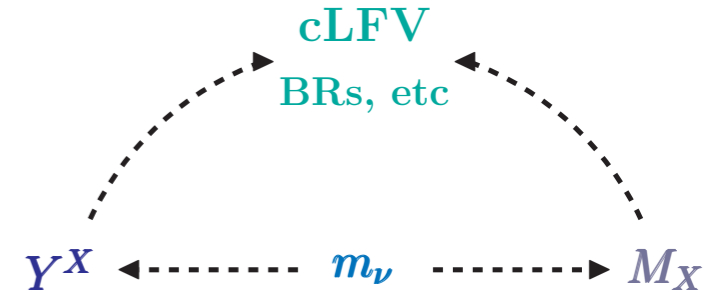
$\propto \sum_{j=1}^{3+n_s} U_{\alpha j}$
 $\propto (1 - \eta) U_{\text{PMNS}}$

$\propto \sum_{i,j=1}^{3+n_s} \left(\sum_{\rho} U_{i\rho}^\dagger U_{\rho j} \right)$
 $\propto (1 - 2\eta) U_{\text{PMNS}}$

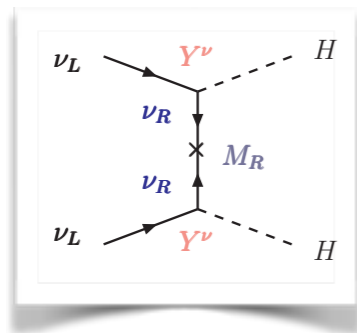
$$\tilde{U}_{\text{PMNS}} = (1 - \eta) U_{\text{PMNS}} ; \eta = \frac{1}{2} \theta \theta^\dagger$$

Type I seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate "balance" between sources of flavour violation and masses of new propagators
 ⇒ account for oscillation data (observation!)



- ▶ **Type I Seesaw:** extend the SM via (Majorana) sterile fermions

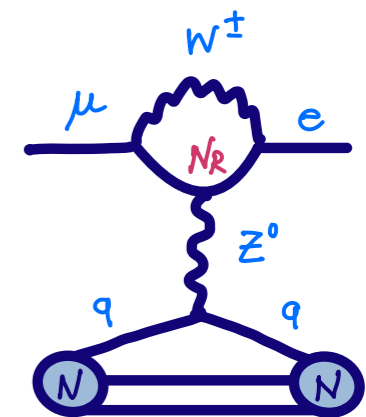
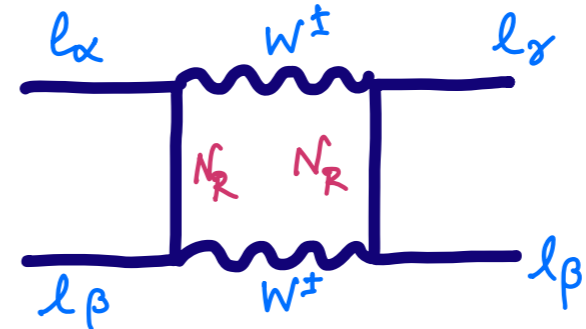
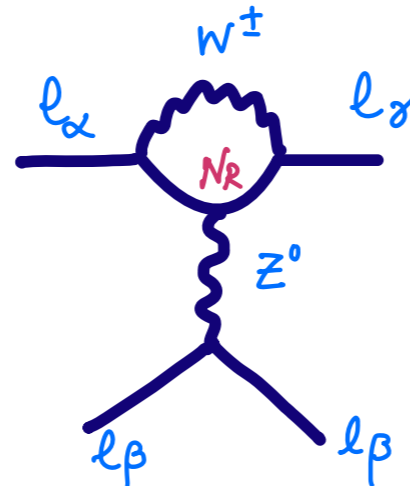
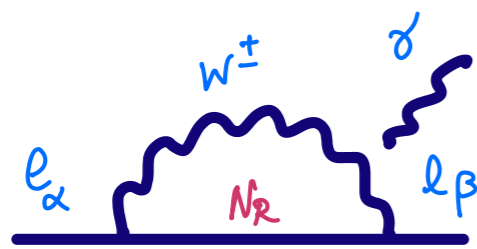


$$\mathcal{L}_{\text{Type I}} \supset \nu Y_\nu \bar{\nu}_L \nu_R + \frac{1}{2} M_R \bar{\nu}_R^c \nu_R$$

active-sterile mixings: $\nu_L - \nu_R$

$$\Rightarrow \theta \approx \mathcal{O}(m_D^\dagger M_R^{-1})$$

$$m_\nu \sim (Y^\nu v)^T \frac{1}{M_R} (Y^\nu v)$$

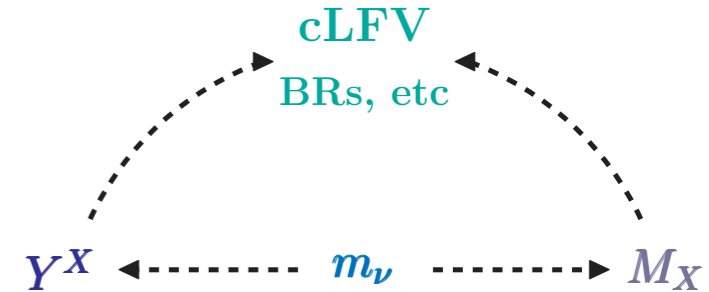


All @ loop-level...

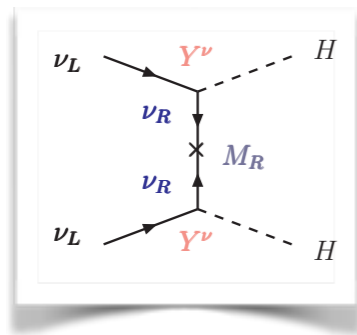
cLFV

Type I seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate **"balance"** between sources of flavour violation and masses of new propagators
 ⇒ account for **oscillation data** (observation!)



- ▶ **Type I Seesaw:** extend the SM via **(Majorana) sterile fermions**



- ~> an enlarged spectrum
- ~> extended mixings

If light neutrino masses generated by **"natural" new physics** ⇒ very **high energy NP scale**

$$Y^\nu \sim \mathcal{O}(1)$$

$$M_R \sim 10^{14-16} \text{ GeV}$$

$$m_\nu \sim (Y^\nu \nu)^T \frac{1}{M_R} (Y^\nu \nu)$$

$$U^T \mathcal{M}_\nu^{6 \times 6} U = \text{diag}(m_{\nu_i})$$

$$U = \begin{pmatrix} U_{\nu\nu} & U_{\nu N} \\ U_{N\nu} & U_{NN} \end{pmatrix}$$

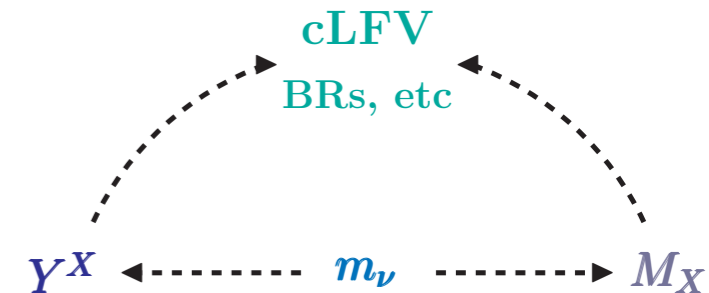
leptonic mixing $\approx U_{\text{PMNS}}$
(unitary to very good approximation)

negligible active-sterile mixings ($\theta \propto m_D^\dagger M_R^{-1}$)

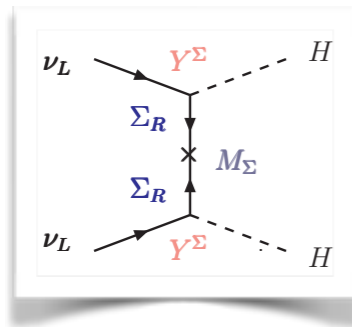
⇒ **Decoupled new physics!** No contributions for cLFV observables, no resonance within collider reach...

Type III seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate "balance" between sources of flavour violation and masses of new propagators
 ⇒ account for oscillation data (observation!)



- ▶ **Type III Seesaw:** extend the SM via SU(2) triplet fermions



- ~> an enlarged spectrum
- ~> extended mixings

$$m_\nu \sim (Y_\Sigma \nu)^T \frac{1}{M_\Sigma} (Y_\Sigma \nu)$$

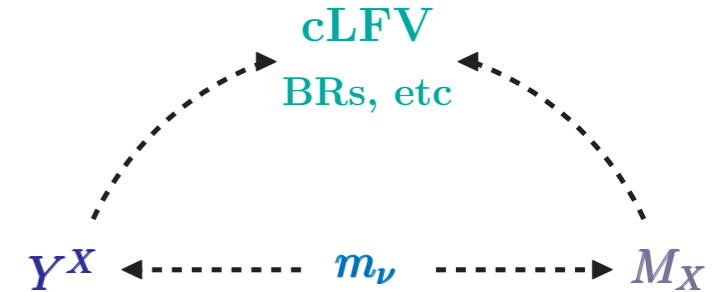
$$\Sigma = \Sigma^+, \Sigma^0, \Sigma^-$$

$$\mathcal{L}_{\text{Type III}} \supset \nu Y_\Sigma \Sigma^+ \ell^- + \nu Y_\Sigma \Sigma^0 \nu + M_\Sigma \bar{\Sigma} \Sigma$$

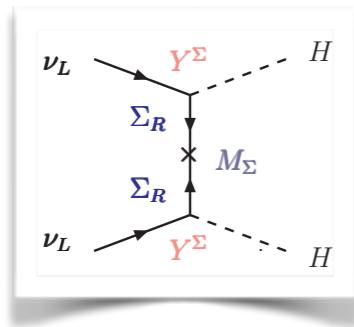
Fermion-triplet mixings: $\Sigma^0 - \nu$ and $\Sigma^{+c} - \ell^-$
 ⇒ $\theta \approx \mathcal{O}(\nu Y_\Sigma M_\Sigma^{-1})$

Type III seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate "balance" between sources of flavour violation and masses of new propagators
 ⇒ account for oscillation data (observation!)



- ▶ **Type III Seesaw:** extend the SM via SU(2) triplet fermions



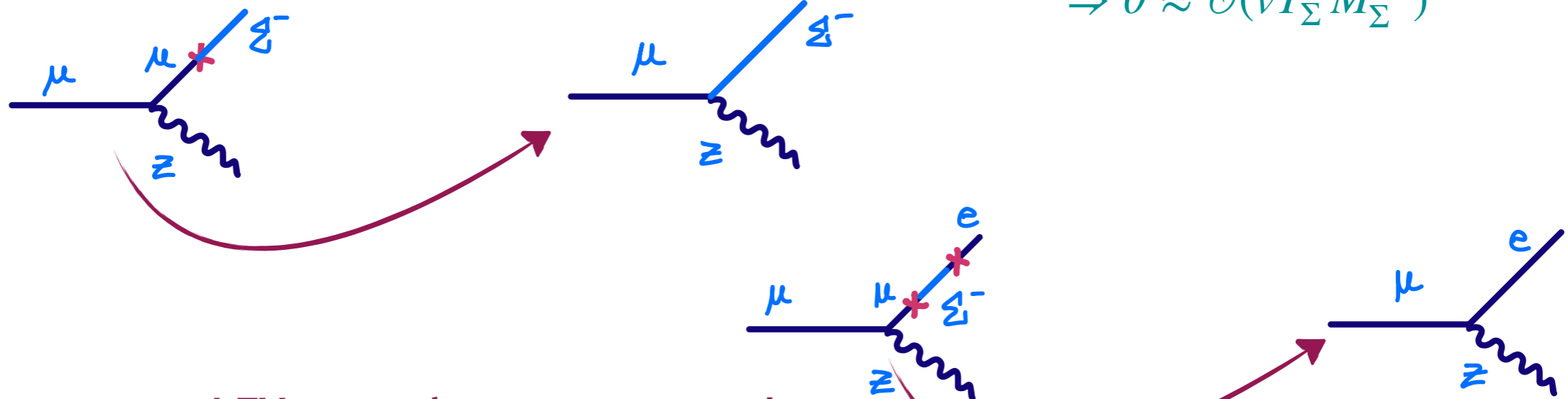
- ~ an enlarged spectrum
- ~ extended mixings

$$\Sigma = \Sigma^+, \Sigma^0, \Sigma^-$$

$$\mathcal{L}_{\text{Type III}} \supset \nu Y_\Sigma \Sigma^+ \ell^- + \nu Y_\Sigma \Sigma^0 \nu + M_\Sigma \bar{\Sigma} \Sigma$$

Fermion-triplet mixings: $\Sigma^0 - \nu$ and $\Sigma^{+c} - \ell^-$
 $\Rightarrow \theta \approx \mathcal{O}(\nu Y_\Sigma M_\Sigma^{-1})$

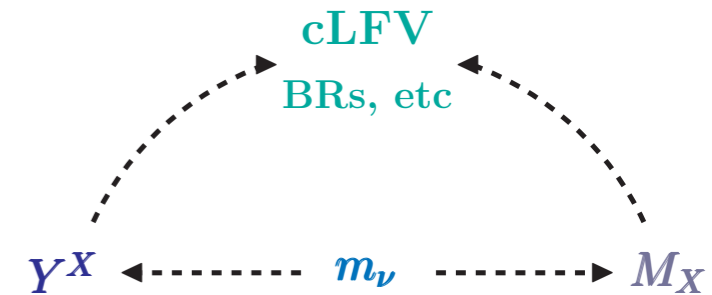
$$m_\nu \sim (Y_\Sigma \nu)^T \frac{1}{M_\Sigma} (Y_\Sigma \nu)$$



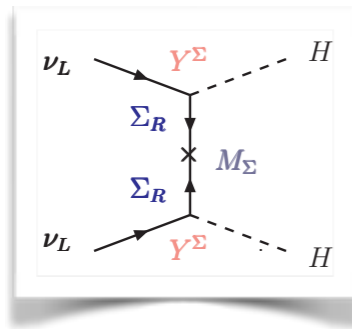
cLFV neutral currents @ tree!

Type III seesaw and cLFV

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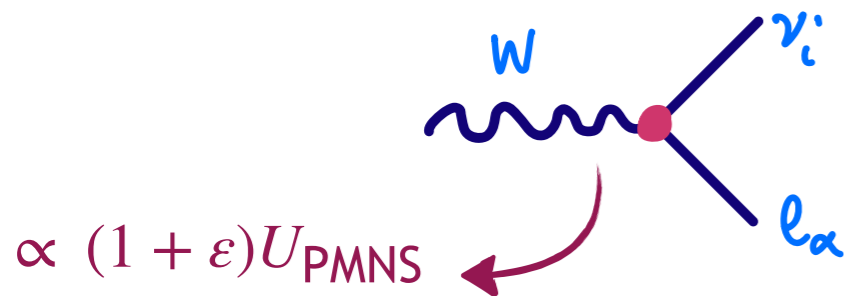
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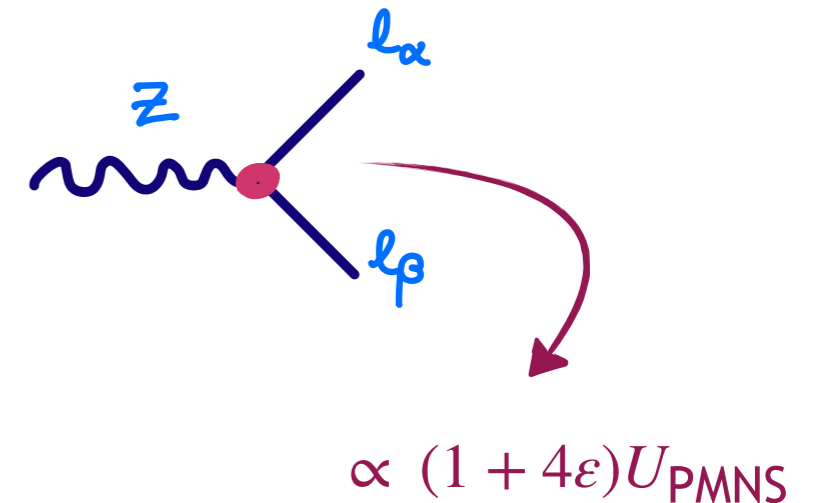
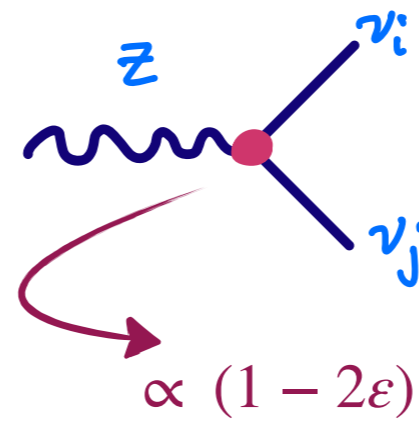
$$\Sigma = \Sigma^+, \Sigma^0, \Sigma^-$$

$$\mathcal{L}_{\text{Type III}} \supset \nu Y_\Sigma \Sigma^+ \ell^- + \nu Y_\Sigma \Sigma^0 \nu + M_\Sigma \bar{\Sigma} \Sigma$$

Fermion-triplet mixings: $\Sigma^0 - \nu$ and $\Sigma^{+c} - \ell^-$
 $\Rightarrow \theta \approx \mathcal{O}(v Y_\Sigma M_\Sigma^{-1})$

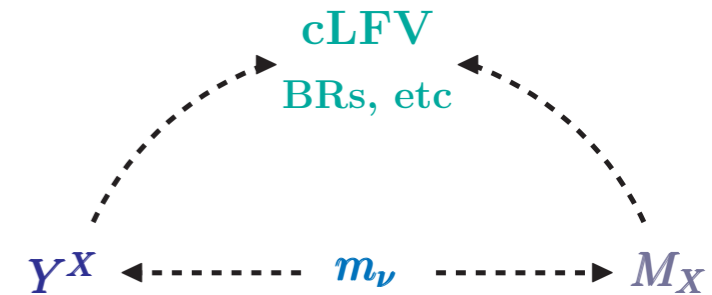


Deviations from unitarity: $\epsilon = \frac{1}{2} m_\Sigma^\dagger M_\Sigma^{-2} m_\Sigma$

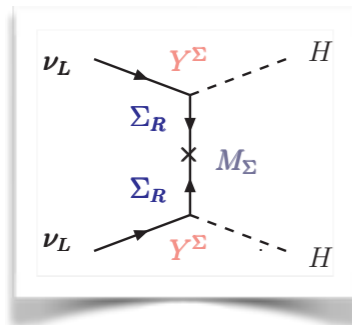


Type III seesaw and cLFV

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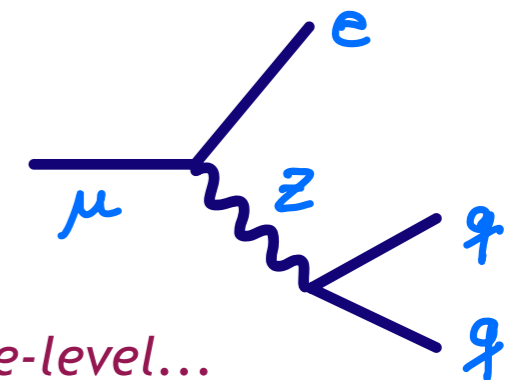
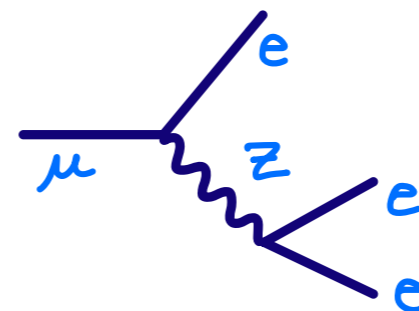
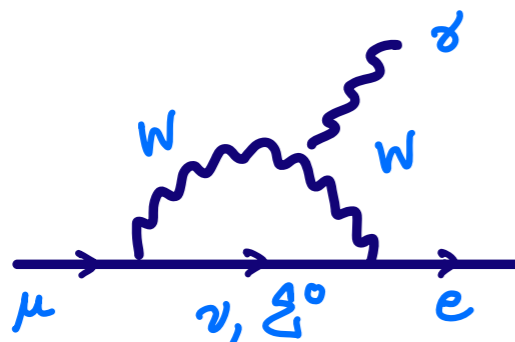
- ~> an enlarged spectrum
- ~> extended mixings

$$\Sigma = \Sigma^+, \Sigma^0, \Sigma^-$$

$$\mathcal{L}_{\text{Type III}} \supset \nu Y_\Sigma \Sigma^+ \ell^- + \nu Y_\Sigma \Sigma^0 \nu + M_\Sigma \bar{\Sigma} \Sigma$$

Fermion-triplet mixings: $\Sigma^0 - \nu$ and $\Sigma^{+c} - \ell^-$
 ⇒ $\theta \approx \mathcal{O}(v Y_\Sigma M_\Sigma^{-1})$

$$m_\nu \sim (Y_\Sigma v)^T \frac{1}{M_\Sigma} (Y_\Sigma v)$$

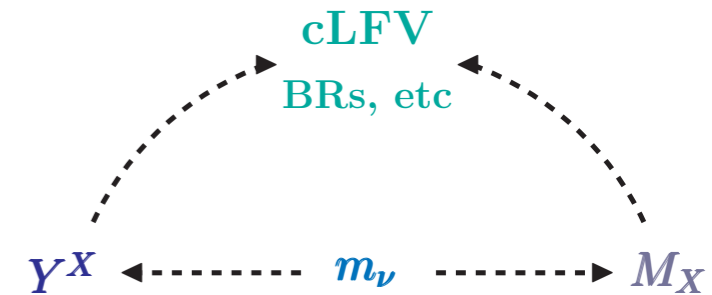


Tree-body decays and conversion in nuclei @ tree-level...
 due to the modified $Z \ell_\alpha \ell_\beta$ vertex!

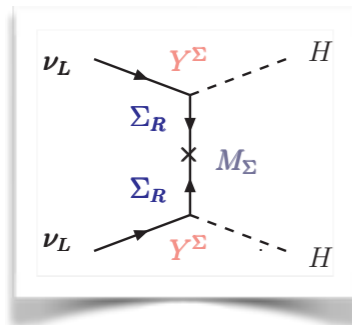
cLFV

Type III seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate **"balance"** between sources of flavour violation and masses of new propagators
 ⇒ account for **oscillation data** (observation!)



- ▶ **Type III Seesaw:** extend the SM via SU(2) **triplet fermions**



- ~> an enlarged spectrum
- ~> extended mixings

$$m_\nu \sim (Y_\Sigma \nu)^T \frac{1}{M_\Sigma} (Y_\Sigma \nu)$$

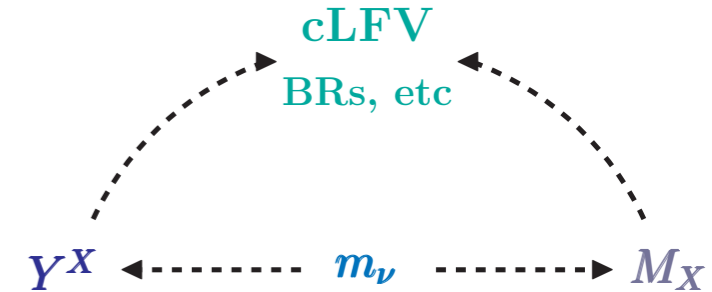
If light neutrino masses generated by **"natural"** new physics ⇒ very **high energy NP scale**
 $Y_\Sigma \sim \mathcal{O}(1)$ $M_\Sigma \sim 10^{14-16}$ GeV

negligible **mixings** between active neutrinos and NP states ($\theta \propto m_\Sigma^\dagger M_\Sigma^{-1}$)

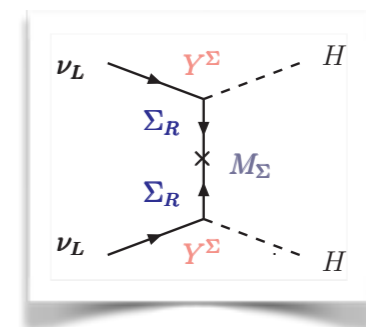
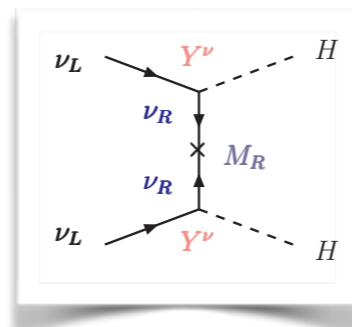
⇒ **Decoupled new physics!** Little contributions for cLFV observables, no resonance within collider reach...

Fermionic seesaws and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate **"balance"** between sources of flavour violation and masses of new propagators
 \Rightarrow account for **oscillation data** (observation!)



- ▶ **Type I & III Seesaw**: a quick **EFT detour** - integrate out the heavy mediators (N_R, Σ)



Dimension 5
(Weinberg operator)

$$Y_\nu^T M_R^{-1} Y_\nu (\bar{L}_L^c \tilde{\phi}^*) (\tilde{\phi}^\dagger L_L)$$

$$m_\nu \sim (Y^\nu \nu)^T \frac{1}{M_R} (Y^\nu \nu)$$

[see Broncano et al, 0210271]

Dimension 6

$$Y_\nu^\dagger M_R^{-2} Y_\nu (\bar{L}_L \tilde{\phi}^*) \not{\partial} (\tilde{\phi}^\dagger L_L)$$

$$Y_\Sigma^T M_\Sigma^{-1} Y_\Sigma (\bar{L}_L^c \tilde{\phi}^*) (\tilde{\phi}^\dagger L_L)$$

$$m_\nu \sim (Y_\Sigma \nu)^T \frac{1}{M_\Sigma} (Y_\Sigma \nu)$$

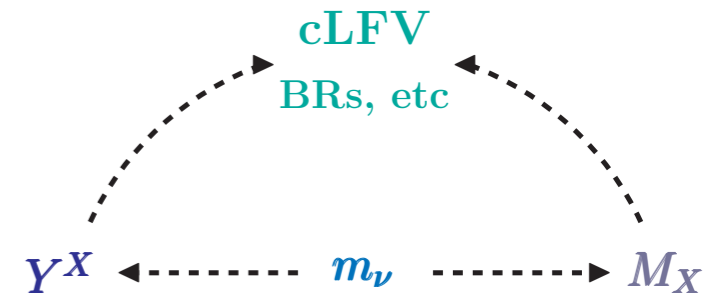
[see Abada et al, 0707.4058]

$$Y_\Sigma^\dagger M_\Sigma^{-2} Y_\Sigma (\bar{L}_L \vec{\tau} \tilde{\phi}) \not{D} (\tilde{\phi}^\dagger \vec{\tau} L_L)$$

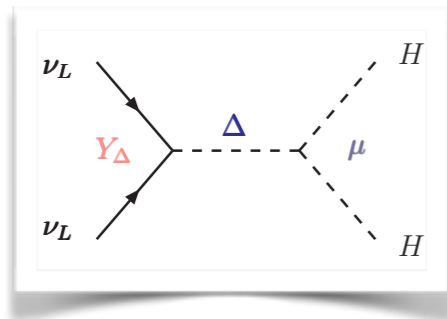
\Rightarrow suppression of **"light neutrino masses"** entails strong suppression of **NP effects!**

Type II seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate "balance" between sources of flavour violation and masses of new propagators
 ⇒ account for oscillation data (observation!)



- ▶ Type II Seesaw: extend the SM via SU(2) triplet scalars



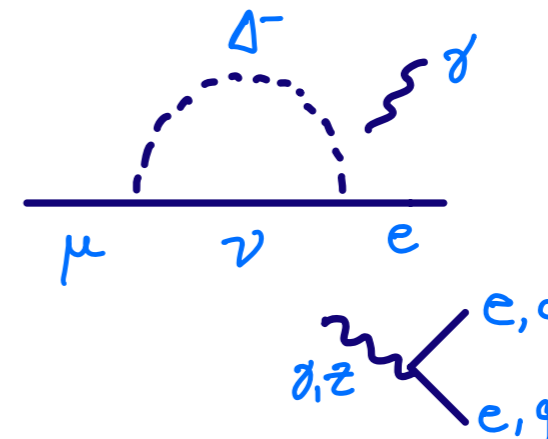
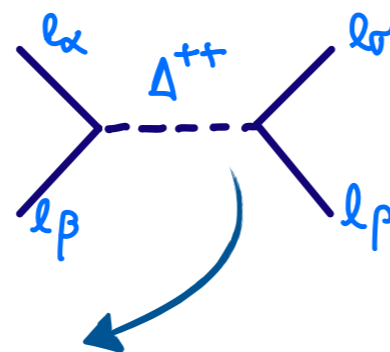
- ~> an enlarged spectrum
- ~> extended mixings

$$\Delta = \Delta^0, \Delta^+, \Delta^{++}$$

$$\mathcal{L}_{\text{Type II}} \supset Y_{\Delta} \bar{L}_L \vec{\tau} L_L^c \vec{\Delta} + \mu_{\Delta} \phi^{\dagger} \vec{\tau} \tilde{\phi} \vec{\Delta}$$

$$m_{\nu} \sim \frac{Y_{\Delta} \mu}{2} \frac{v^2}{M_{\Delta}^2}$$

Interactions with gauge bosons; direct cLFV couplings

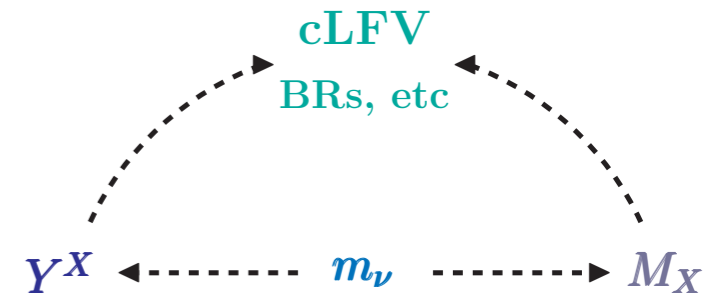


Tree-body decays @ tree-level...

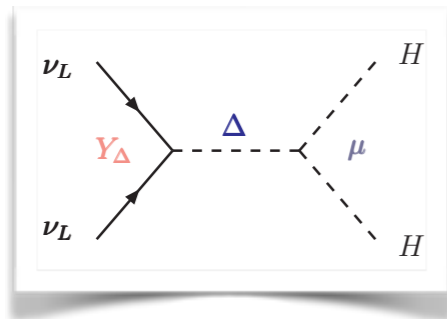
cLFV

Scalar seesaw and cLFV

- ▶ Mechanisms for neutrino mass generation: delicate **"balance"** between sources of flavour violation and masses of new propagators
 ⇒ account for **oscillation data** (observation!)



- ▶ **Type II Seesaw:** extend the SM via SU(2) **triplet scalars**



- ~> an enlarged spectrum
- ~> extended mixings

$$m_\nu \sim \frac{Y_\Delta \mu}{2} \frac{v^2}{M_\Delta^2}$$

A **different scenario:** additional ingredient!

"natural" new physics \nRightarrow very **high energy NP scale**

Smallness of m_ν **also** from (tiny) μ coupling for "natural" Y_Δ and not "too heavy" M_Δ

[see Abada et al, 0707.4058]

Dimension 5 $4 Y_\Delta \mu M_\Delta^{-2} (\bar{L}_L^c \tilde{\phi}^*) (\tilde{\phi}^\dagger L_L)$

Dimension 6 $Y_\Delta Y_\Delta^\dagger M_\Delta^{-2} (\bar{L}_L \gamma_\mu L_L) (\bar{L}_L \gamma^\mu L_L)$

⇒ suppression of **"light neutrino masses"** decorrelated from contribution to **NP effects!**

cLFV and the seesaw: peculiar patterns

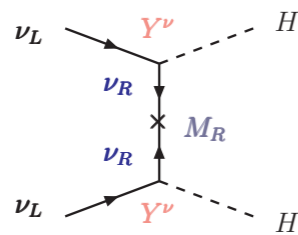


IN2P3
Les deux infinis

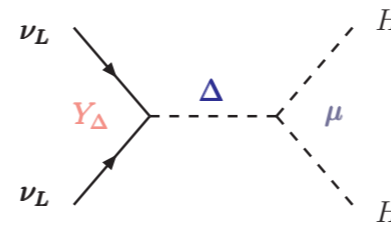
- ▶ **Seesaw realisations:** distinctive expectations for numerous **cLFV observables**
If **observable/measurable cLFV** - what can we learn?

cLFV and the seesaw: peculiar patterns

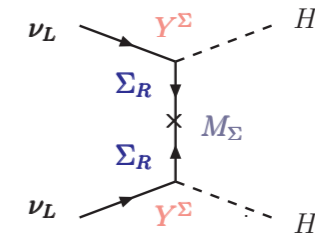
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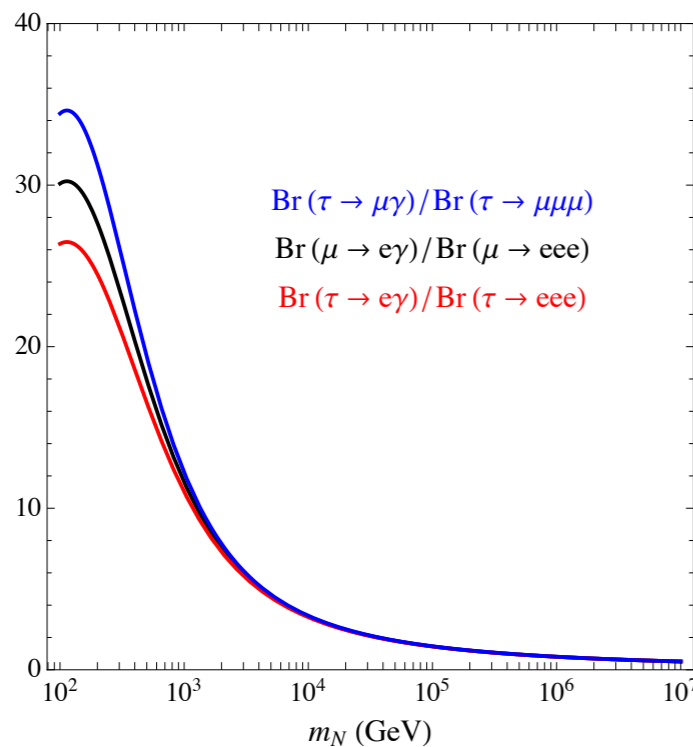
Type I (fermion singlet)



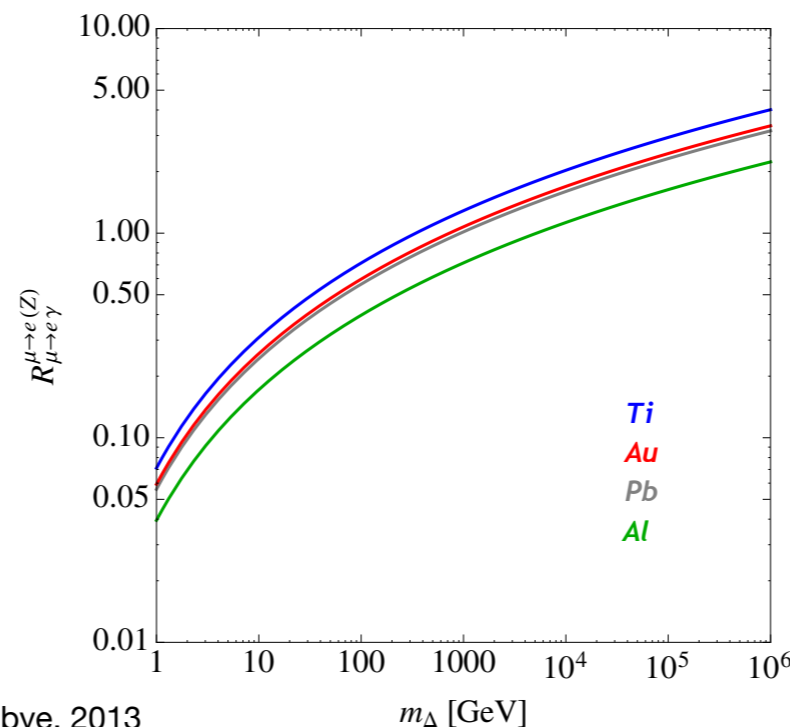
Type II (scalar triplet)



Type III (fermion triplet)



Hambye, 2013



$$\frac{\text{BR}(\mu \rightarrow e \gamma)}{\text{BR}(\mu \rightarrow 3e)} = 1.3 \times 10^{-3}$$

$$\frac{\text{BR}(\tau \rightarrow \mu \gamma)}{\text{BR}(\tau \rightarrow 3\mu)} = 1.3 \times 10^{-3}$$

$$\frac{\text{BR}(\mu \rightarrow e \gamma)}{\text{CR}(e-\mu, \text{Ti})} = 3.1 \times 10^{-4}$$

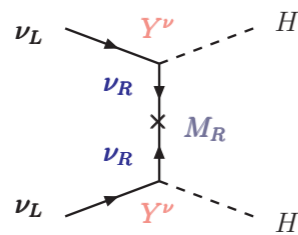
[adapted from Calibbi et al, 1709.00294]

cLFV patterns reflect the **topology** of contributions associated with the new mediators
 (dipole or Z-dominated, tree vs. loop, ...)

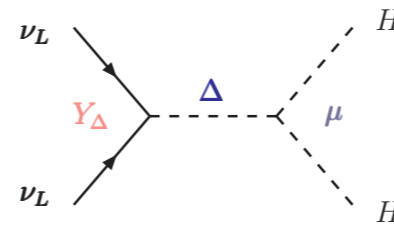
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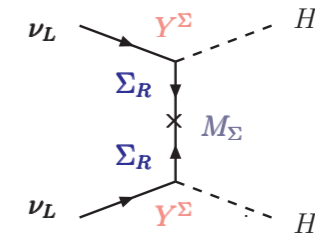
If **observable/measurable cLFV** - what can we learn?



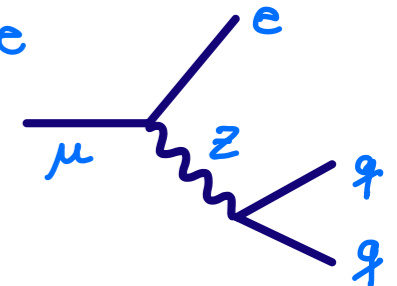
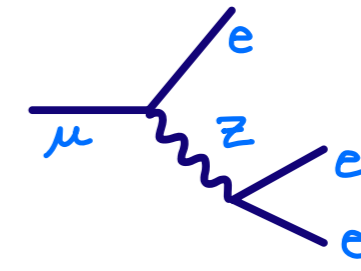
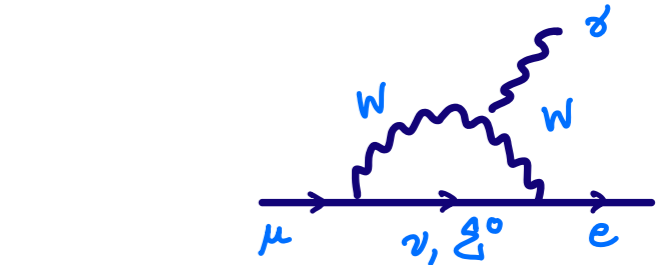
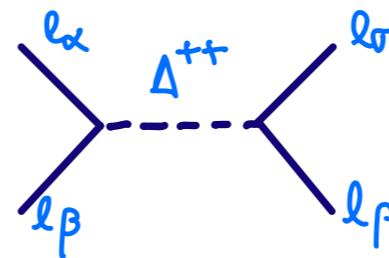
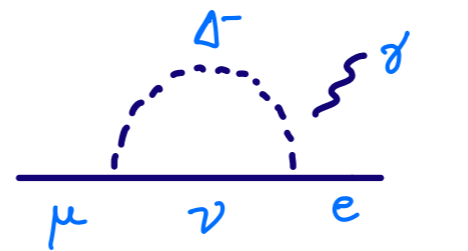
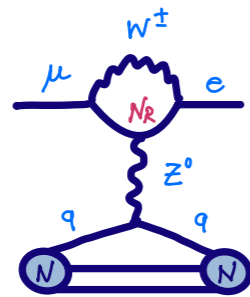
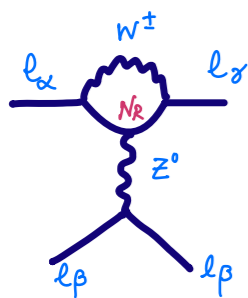
Type I (fermion singlet)



Type II (scalar triplet)



Type III (fermion triplet)



cLFV patterns reflect the **topology** of contributions associated with the new mediators
(dipole or Z-dominated, tree vs. loop, ...)

cLFV and the seesaw: peculiar patterns

► **Seesaw realisations:** distinctive expectations for numerous **cLFV observables**

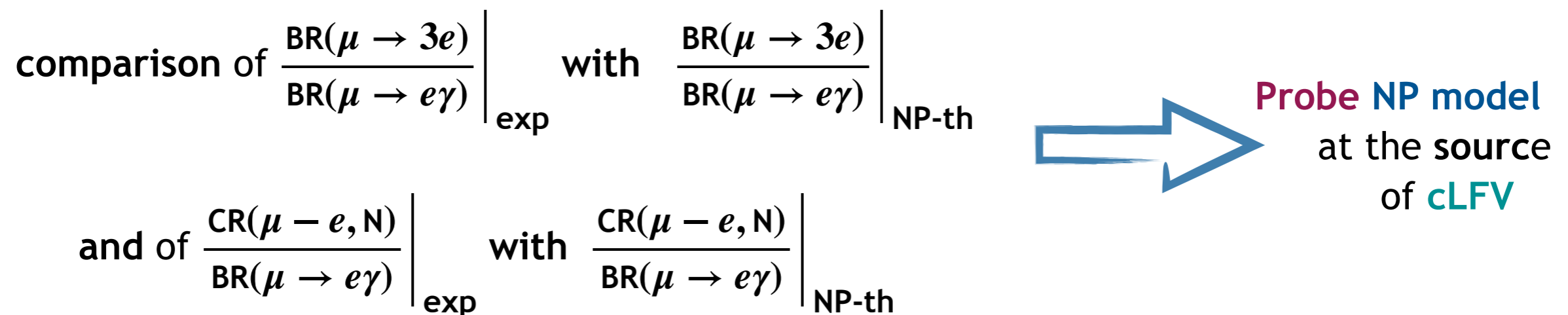
⇒ ratios of **observables** to **identify seesaw mediators** & constrain masses...

cLFV patterns reflect the **topology** of contributions associated with the new mediators
(**dipole** or **Z-dominated**, **tree** vs. **loop**, ...)

Model	$\mu \rightarrow eee$	$\mu N \rightarrow eN$	$\frac{\text{BR}(\mu \rightarrow eee)}{\text{BR}(\mu \rightarrow e\gamma)}$	$\frac{\text{CR}(\mu N \rightarrow eN)}{\text{BR}(\mu \rightarrow e\gamma)}$
MSSM	Loop	Loop	$\approx 6 \times 10^{-3}$	$10^{-3} - 10^{-2}$
Type-I seesaw	Loop*	Loop*	$3 \times 10^{-3} - 0.3$	0.1–10
Type-II seesaw	Tree	Loop	$(0.1 - 3) \times 10^3$	$\mathcal{O}(10^{-2})$
Type-III seesaw	Tree	Tree	$\approx 10^3$	$\mathcal{O}(10^3)$
LFV Higgs	Loop [†]	Loop* [†]	$\approx 10^{-2}$	$\mathcal{O}(0.1)$
Composite Higgs	Loop*	Loop*	0.05 – 0.5	2 – 20

[adapted from Calibbi et al, 1709.00294]

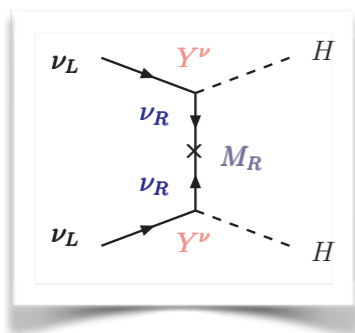
Upon **experimental determination** of rates for **cLFV transitions**:



New Physics paths to cLFV: low-scale seesaws



- Addition of 3 "heavy" Majorana right-handed neutrinos ν_R to the SM
but explore considerably lighter range for M_R $\text{MeV} \leq M_R \leq 10^{\text{few}} \text{TeV}$



Type I (fermion singlet)

$$m_\nu \sim (Y^\nu v)^T \frac{1}{M_R} (Y^\nu v)$$

After EW symmetry breaking, 6 states in the neutral lepton spectrum

$$\mathcal{M}_\nu^{6 \times 6} = \begin{pmatrix} 0 & Y^\nu v \\ (Y^\nu)^T v & M_R \end{pmatrix}$$

3 light neutrinos $m_\nu \approx -v^2 Y_\nu^T M_R^{-1} Y_\nu$

3 heavy states $m_N \approx M_R$

Enlarged 6×6 mixing matrix $U^T \mathcal{M}_\nu^{6 \times 6} U = \text{diag}(m_{\nu_i})$

$$U = \begin{pmatrix} U_{\nu\nu} & U_{\nu N} \\ U_{N\nu} & U_{NN} \end{pmatrix}$$

Non-negligible active-sterile mixings! ($\theta \propto m_D^\dagger M_R^{-1}$)

Non-unitary leptonic mixing \tilde{U}_{PMNS}

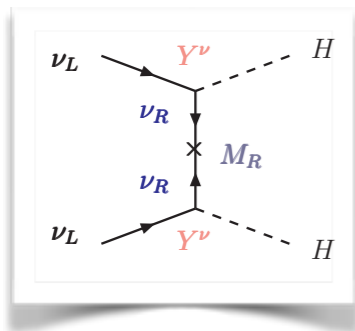
Low-scale realisations of the **Type I seesaw** open door to a **very rich phenomenology** from cLFV signals, to collider searches

Similar implications for low-scale Type III

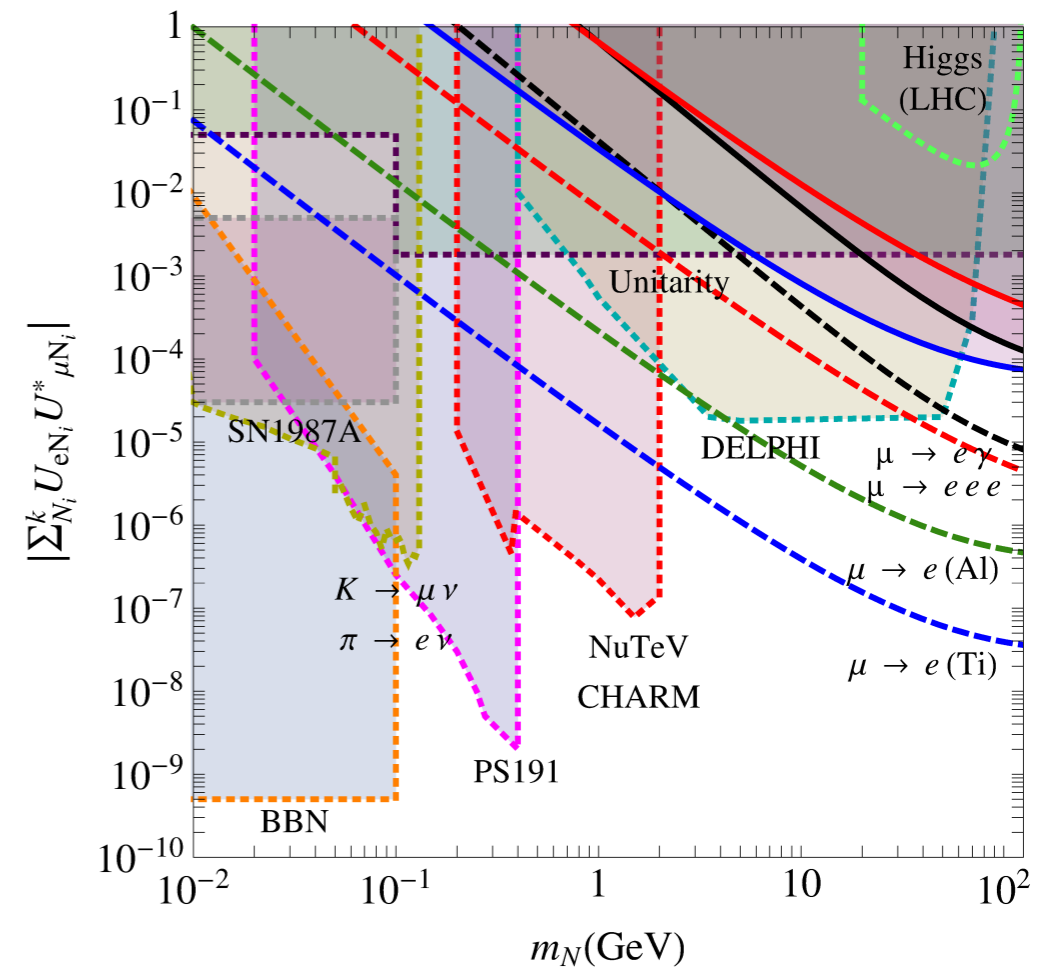
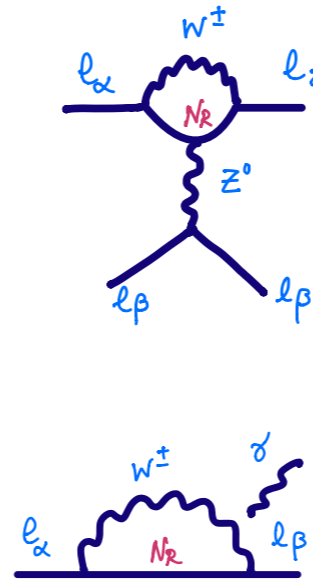
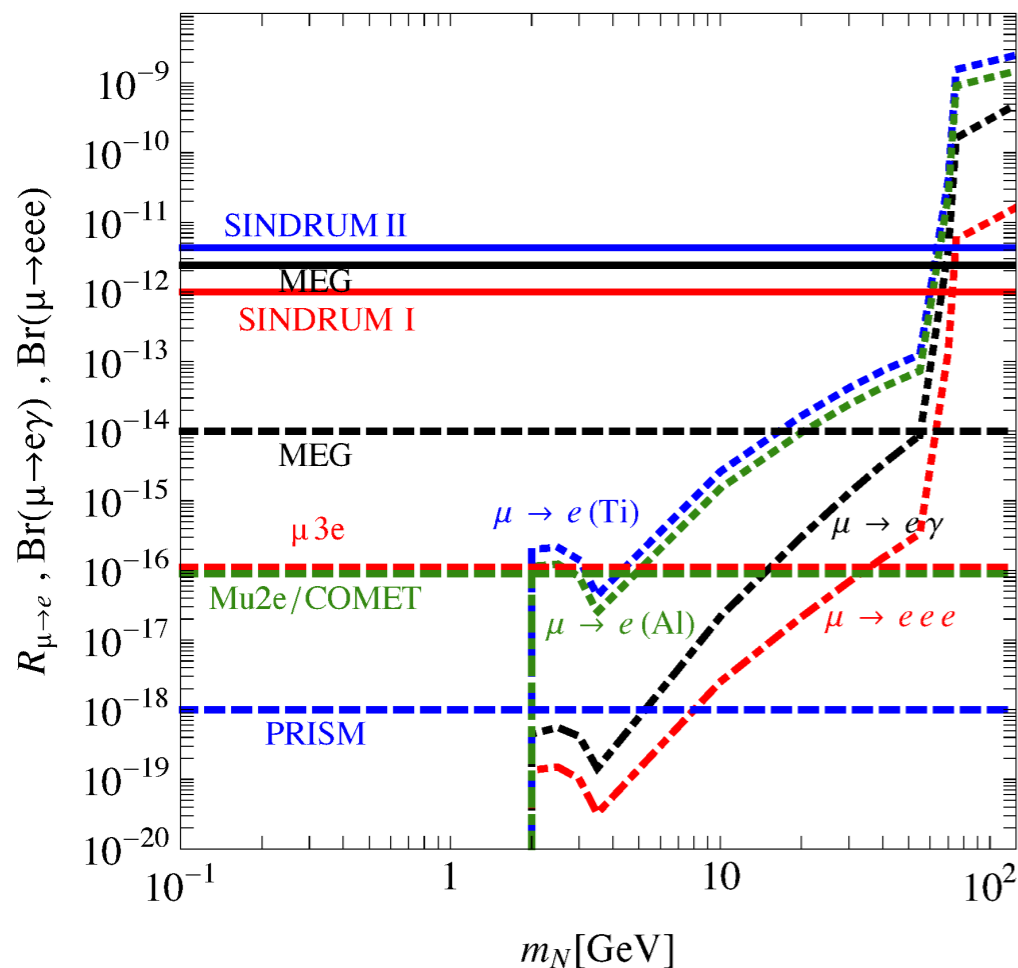
(but important direct/indirect constraints due to the *non-singlet nature of new states...*)

Low-scale models of m_ν generation: type I seesaw

- Addition of 3 "heavy" Majorana right-handed neutrinos ν_R to the SM
but explore considerably lighter range for M_R $\text{MeV} \leq M_R \leq 10^{\text{few}} \text{TeV}$



Low-scale realisations of the **Type I seesaw**: very rich phenomenology
 \Rightarrow cLFV signals (more promising than collider searches)



Alonso, Dehn, Gavela, Hambye [1209.2679]

Low-scale models for m_ν : Inverse Seesaw

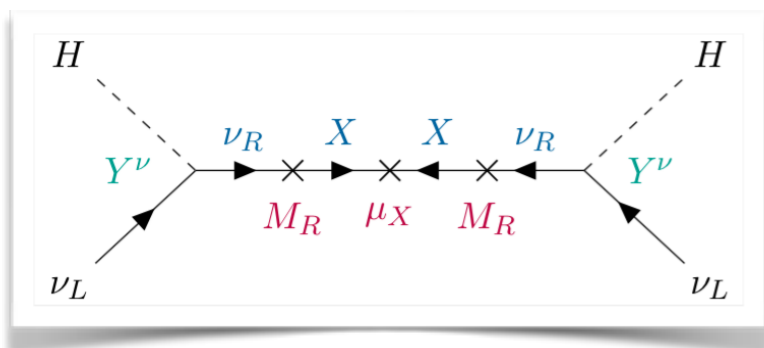
- ▶ Variants of **type I seesaw** aiming at a **natural** realisation of a low-scale m_ν mechanism
- ▶ Addition of **two new species** of **fermionic gauge singlets**

n_R right-handed neutrinos ν_R ($L_{\nu_R} = 1$) and n_X extra sterile states X ($L_X = -1$)

$$\mathcal{L}_{ISS}^{(3,3)} = - Y^\nu \bar{L} \tilde{H} \nu_R - M_R \bar{\nu}_R^c X - \frac{1}{2} \mu_X \bar{X}^c X$$

[Mohapatra and Valle, '86]

lepton number violating!



$$\mathcal{M}_{ISS}^{9 \times 9} = \begin{pmatrix} 0 & Y_\nu v & 0 \\ Y_\nu^T v & 0 & M_R \\ 0 & M_R & \mu_X \end{pmatrix} \Rightarrow \begin{cases} 3 \text{ light } \nu : m_\nu \approx \frac{(Y_\nu v)^2}{(Y_\nu v)^2 + M_R^2} \mu_X \\ 3 \text{ pseudo-Dirac pairs} : m_{N\pm} \approx M_R \pm \mu_X \end{cases} \quad ISS(3,3)$$

$$m_\nu \sim (Y^\nu v)^T \frac{\mu_X}{M_R^2} (Y^\nu v)$$

Interplay of **two scales** driving smallness of m_ν : M_R and μ_X

For natural values of $Y^\nu \sim \mathcal{O}(1)$

comparatively "light" heavy spectrum ($\Lambda_{EW} \leftrightarrow \text{TeV}$) for small values of μ_X (around eV - keV)

Natural ('t Hooft criterium) since **B-L conservation restored** when $\mu_X \rightarrow 0$!

Symmetry protected "smallness" of m_ν - approximate LNC

Low-scale models for m_ν : Inverse Seesaw

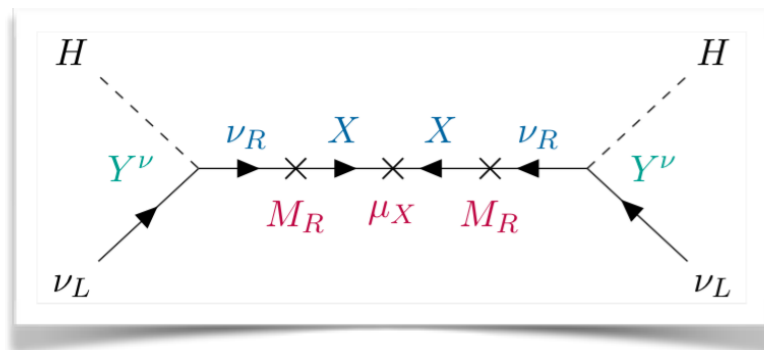
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lepton number violating!



$$\mathcal{M}_{ISS}^{9 \times 9} = \begin{pmatrix} 0 & Y_\nu v & 0 \\ Y_\nu^T v & 0 & M_R \\ 0 & M_R & \mu_X \end{pmatrix} \Rightarrow \begin{cases} 3 \text{ light } \nu : m_\nu \approx \frac{(Y_\nu v)^2}{(Y_\nu v)^2 + M_R^2} \mu_X \\ 3 \text{ pseudo-Dirac pairs : } m_{N\pm} \approx M_R \pm \mu_X \end{cases} \quad ISS(3,3)$$

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Interplay of **two scales** driving smallness of m_ν : M_R and μ_X

For natural values of $Y^\nu \sim \mathcal{O}(1)$

comparatively "light" heavy spectrum ($\Lambda_{EW} \leftrightarrow \text{TeV}$) for small values of μ_X (around eV - keV)

\Rightarrow Despite small $m_\nu \sim \mu_X \frac{m_D^2}{M_R^2}$, a "low" NP scale $\sim M_R$, and **sizeable mixings** ($\theta \propto m_D^\dagger M_R^{-1}$) !

cLFV and EW precision in the ISS

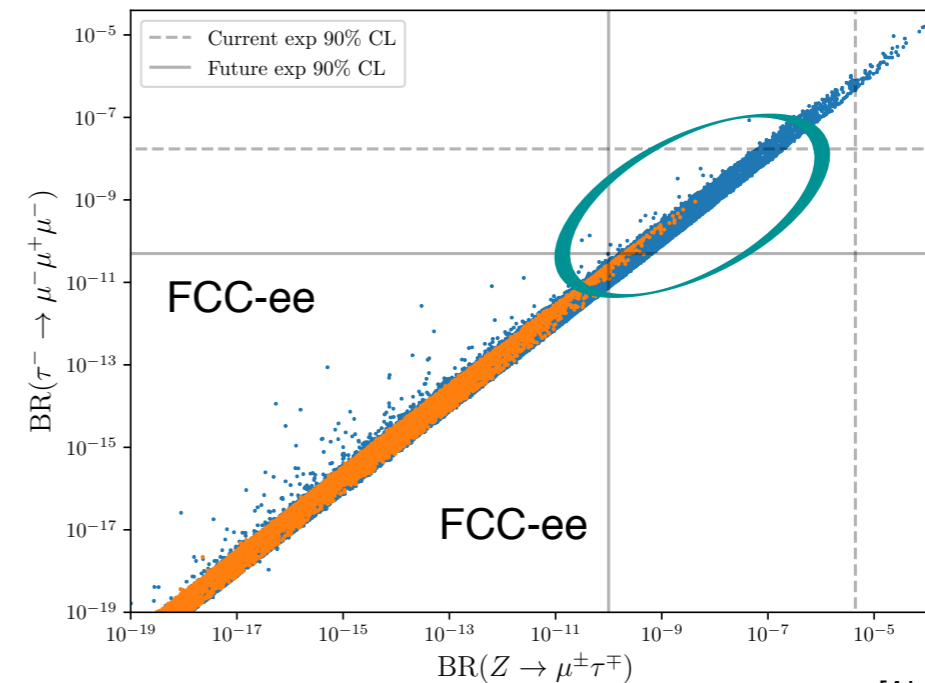
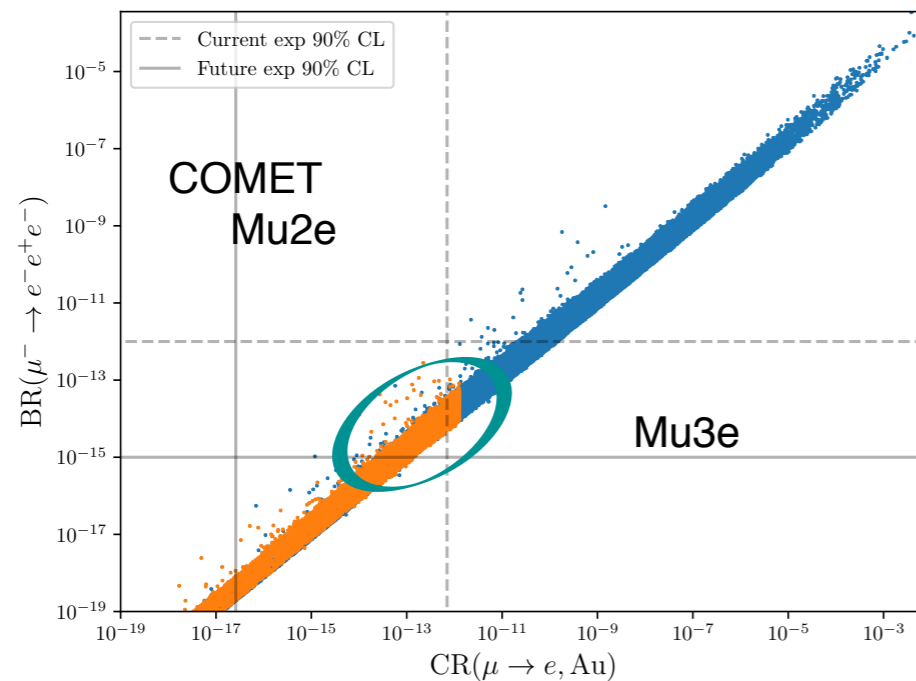


Low-scale models for m_ν : Inverse Seesaw

Inverse seesaw: well-motivated low-scale mechanism of neutrino mass generation

$$\text{ISS}(3,3) \Rightarrow \text{SM} + 3 \nu_R + 3 X$$

(rich phenomenology \Rightarrow *testability!*)



[Abada, Kriewald, Pinsard, Rosauero, AMT, '23]

\Rightarrow Abundant "flavour" signals: cLFV transitions (at low and high energies)

Regimes *already disfavoured* from current bounds!

cLFV actively **constrains** parameter space of ISS

\Rightarrow Opportunities to **observe cLFV** in (near-)future facilities:

$\mu - e$ sector @ Mu3e, COMET & Mu2e

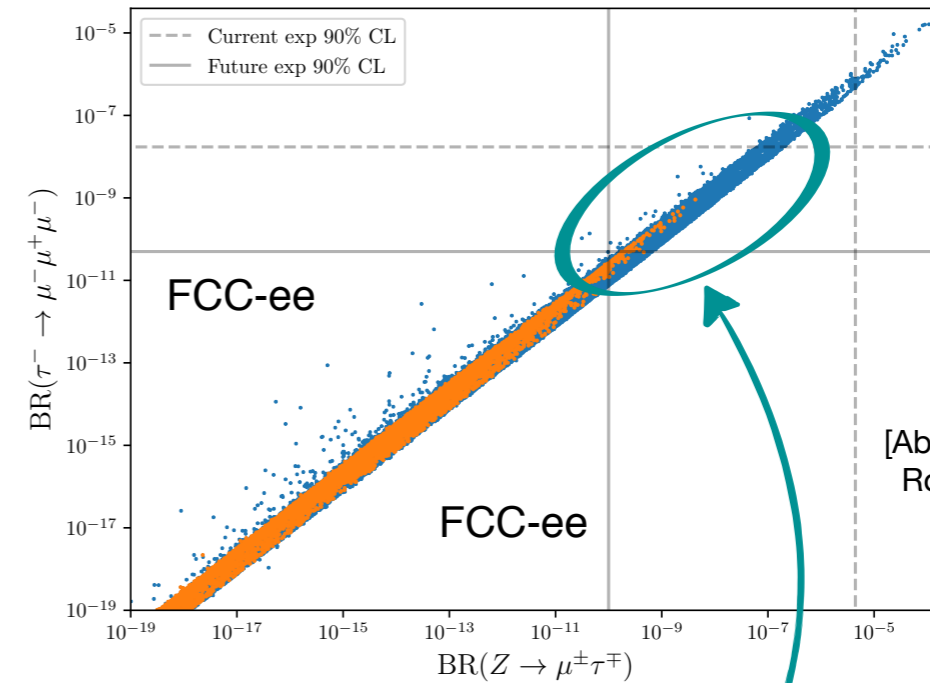
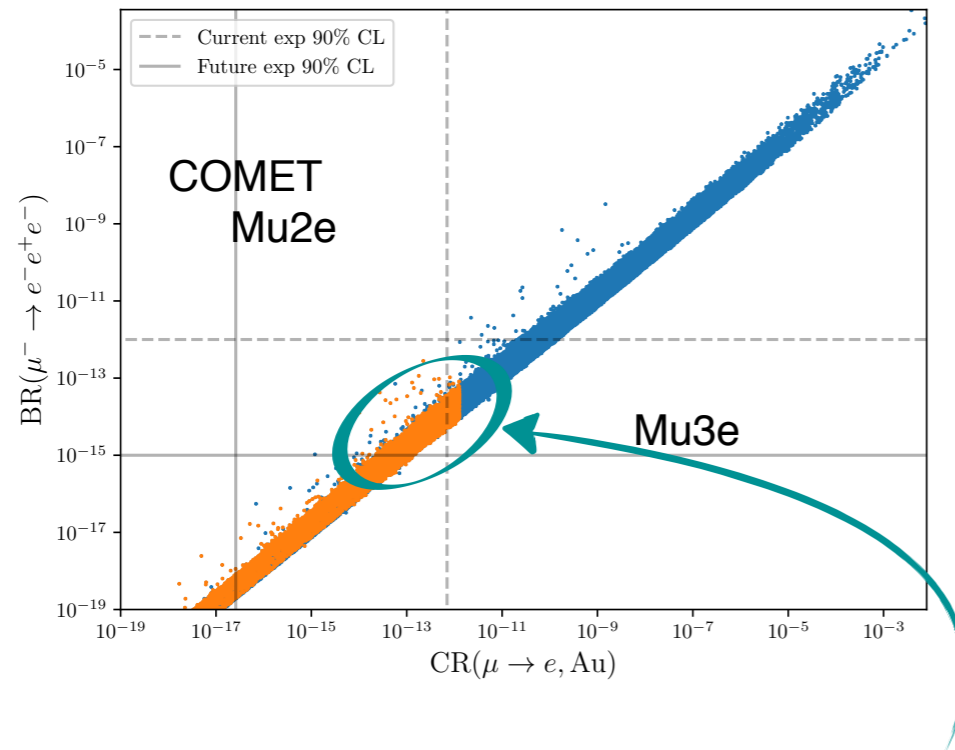
$\tau - \mu$ sector @ Belle II, FCC-ee, ...

Low-scale models for m_ν : Inverse Seesaw

Inverse seesaw: well-motivated low-scale mechanism of neutrino mass generation

$$\text{ISS}(3,3) \Rightarrow \text{SM} + 3 \nu_R + 3 X$$

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\Rightarrow Correlated observables! $\mu \rightarrow 3e$ vs. $\mu - e$ conversion

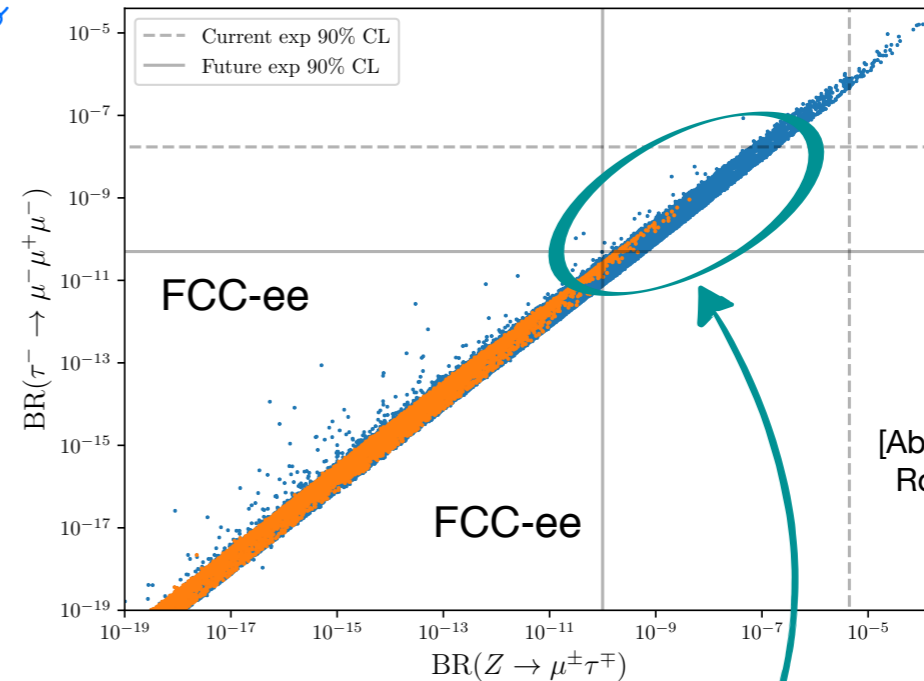
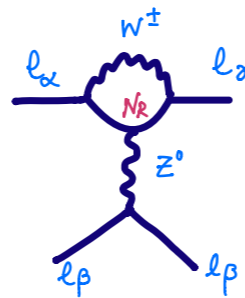
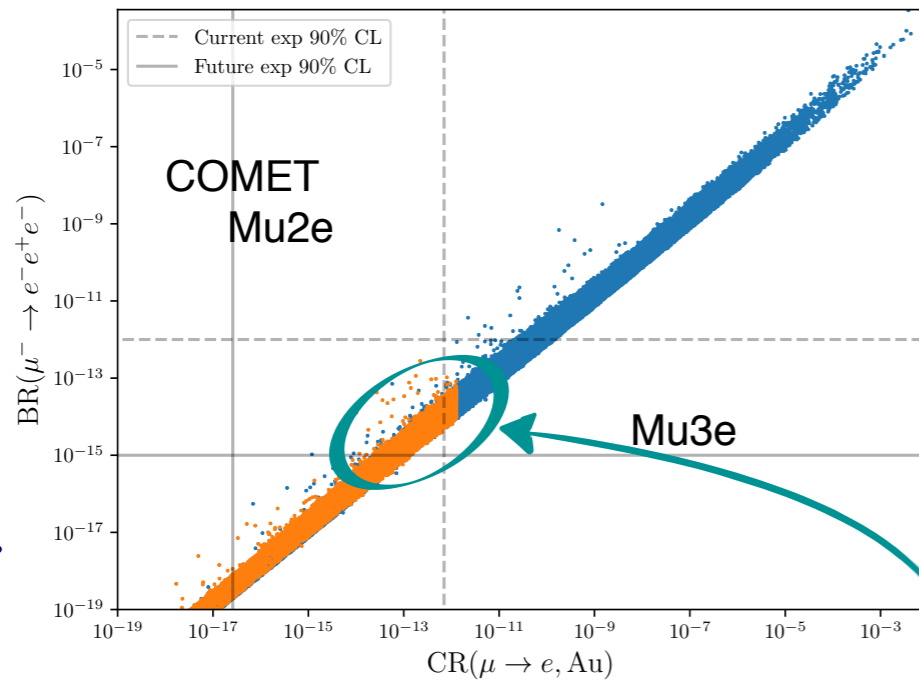
and $Z \rightarrow \mu\tau$ vs. $\tau \rightarrow 3\mu$

Low-scale models for m_ν : Inverse Seesaw

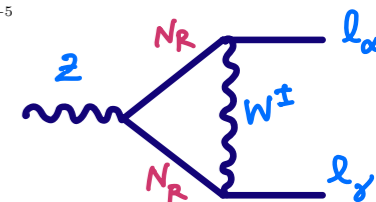
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\Rightarrow Correlated observables! $\mu \rightarrow 3e$ vs. $\mu - e$ conversion

and $Z \rightarrow \mu\tau$ vs. $\tau \rightarrow 3\mu$

A consequence of the **dominant contribution of Z-penguins** in the 3-body decays and in neutrinoless conversion in nuclei (for the most "observable" regimes...)

Observation of $\mu \rightarrow 3e \Rightarrow$ observation of $\mu - e$ conversion
 $\tau \rightarrow 3\mu \Rightarrow$ observation of $Z \rightarrow \mu\tau$

testability!?

Low-scale models for m_ν : Inverse Seesaw



IN2P3
Les deux infinis

Inverse seesaw: well-motivated **low-scale mechanism of neutrino mass generation**

$$\text{ISS}(3,3) \Rightarrow \text{SM} + 3 \nu_R + 3 X \quad (\text{rich phenomenology} \Rightarrow \text{testability!})$$

⇒ Abundant **"flavour" signals**: **cLFV** transitions (at low and high energies)

⇒ **Precision tests**: **electroweak observables** and **lepton flavour universality**

$$\text{LFUV-sensitive: } R_Z^{\alpha\beta} = \frac{\Gamma(Z \rightarrow \ell_\alpha^+ \ell_\alpha^-)}{\Gamma(Z \rightarrow \ell_\beta^+ \ell_\beta^-)}, \quad R_Z^{\alpha\beta} \Big|_{\text{SM}} \simeq 1 \quad \& \quad R_W^{\alpha\beta} = \frac{\Gamma(W \rightarrow \ell_\alpha \nu)}{\Gamma(W \rightarrow \ell_\beta \nu)}, \quad R_W^{\alpha\beta} \Big|_{\text{SM}} \simeq 1$$

EWPO: $\Gamma(Z \rightarrow \text{inv.})$, oblique parameters, ...

New parametrisation (beyond Casas-Ibarra) to access key-regimes of **large LFUV**
(escaping **cLFV** constraints!)

$$\mu_X = M_R^T m_D^{(-1)} U_{\text{PMNS}}^* m_\nu^{\text{diag}} U_{\text{PMNS}}^\dagger (m_D^T)^{(-1)} M_R, \quad Y_D = y_i^d \mathcal{V} M_R^\dagger$$

Massive (semi-)analytical computation of form-factors, renormalisation, ...

[Abada, Kriewald, Pinsard, Rosauero, AMT, 2307.02558]

► **Non-negligible NLO corrections** to **ISS** vertices! **Zℓℓ**, **Zνν** and **Wℓν**

Low-scale models for m_ν : Inverse Seesaw

Inverse seesaw: well-motivated **low-scale mechanism of neutrino mass generation**

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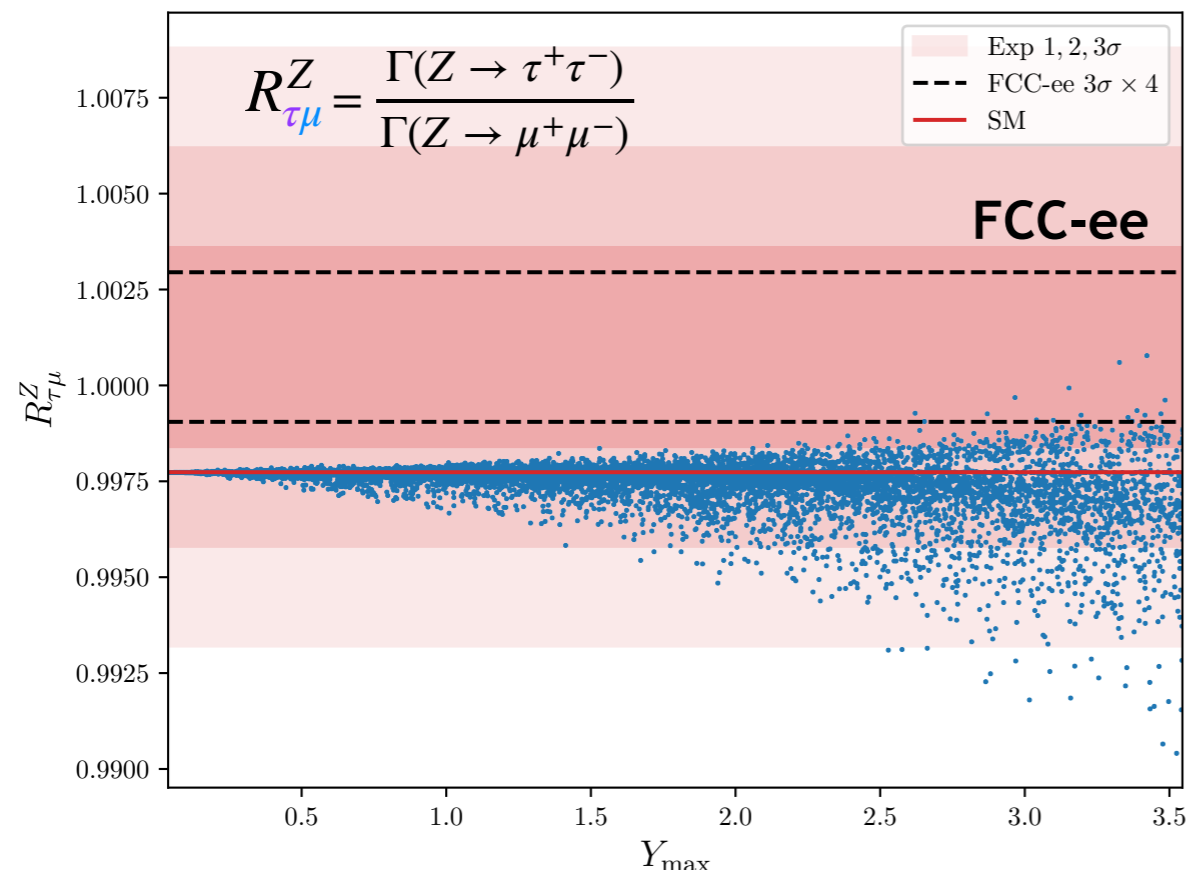
\Rightarrow **Abundant "flavour" signals:** **cLFV** transitions (at low and high energies)

\Rightarrow **Precision tests:** **LFUV** in **Z decays**

large deviations from **SM** possible (sizeable **NLO** corrections)

$$R_{\tau\mu}^Z = \frac{\Gamma(Z \rightarrow \tau^+\tau^-)}{\Gamma(Z \rightarrow \mu^+\mu^-)}$$

[Abada, Kriewald, Pinsard, Rosauero, AMT, 2307.02558]



Contributions around **SM prediction**
(already in tension with measurement)

Large deviations for sizeable Yukawas
(corresponding to masses ~ 5 TeV)

Significant NLO corrections to LFU in Z decays!

\Rightarrow **FCC-ee** expected to probe these regimes
(increase in experimental precision)

Low-scale models for m_ν : Inverse Seesaw

Inverse seesaw: well-motivated **low-scale mechanism of neutrino mass generation**

$$\text{ISS}(3,3) \Rightarrow \text{SM} + 3 \nu_R + 3 X$$

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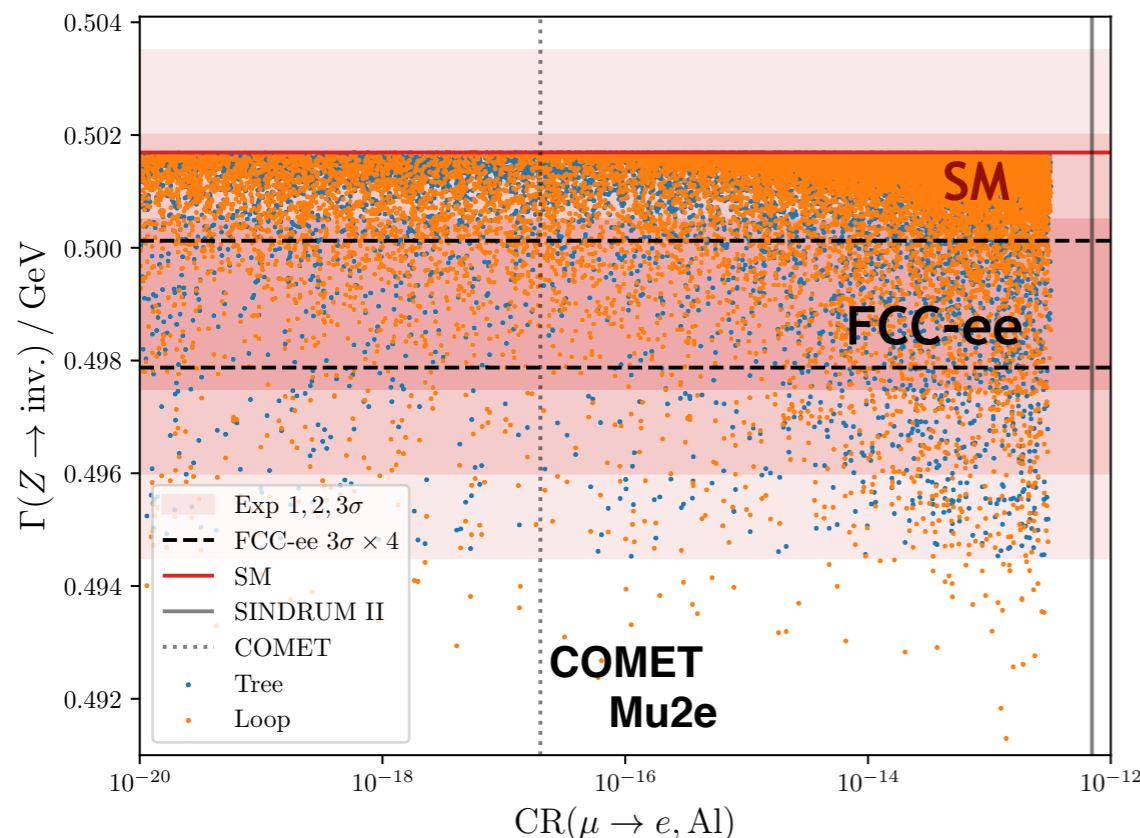
\Rightarrow Abundant **"flavour" signals: cLFV** transitions (at low and high energies)

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large deviations from SM possible (sizeable **NLO corrections**)

\Rightarrow (EW) precision tests: **Invisible Z decays**

[Abada, Kriewald, Pinsard, Rosauero, AMT, EPJC 84 (2024) 2]



Large deviations from SM prediction

Significant NLO corrections to invisible Z decays!

$$\Gamma_{\text{tree}}(Z \rightarrow \text{inv.}) - \Gamma_{\text{loop}}(Z \rightarrow \text{inv.}) \text{ up to } 5 \text{ MeV!}$$

(current exp. uncertainty 1.5 MeV...)

NLO corrections: master theory uncertainties, on par with experiment!

\Rightarrow **FCC-ee** expected to probe important regimes (increase in experimental precision)

Low-scale models for m_ν : Inverse Seesaw

Inverse seesaw: well-motivated **low-scale mechanism of neutrino mass generation**

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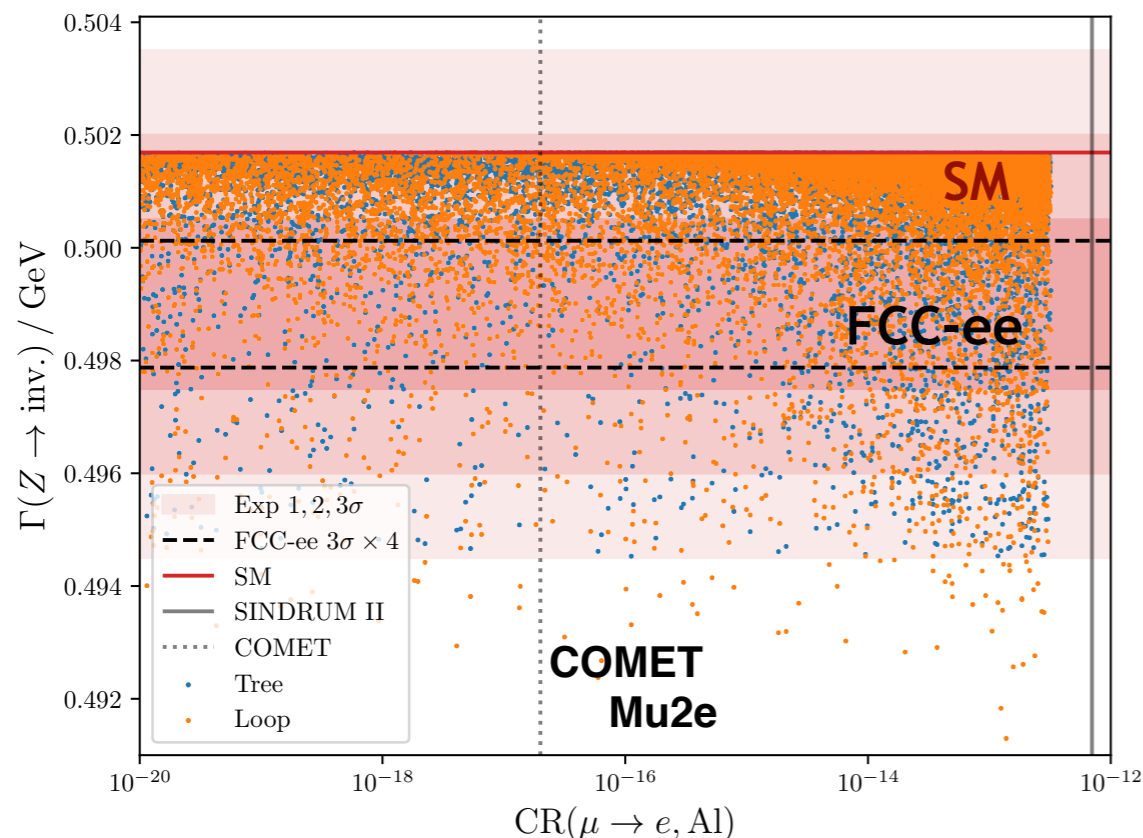
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large deviations from SM possible (sizeable **NLO corrections**)

\Rightarrow (EW) precision tests: **Invisible Z decays** - large **deviations from SM!**

[Abada, Kriewald, Pinsard, Rosauero, AMT, EPJC 84 (2024) 2]



EWPO vs. **cLFV** - complementary probes

cLFV in $\mu - e$: usually most stringent constraints
(e.g. $\mu \rightarrow e\gamma$, $\mu - e$ conversion...)

Invisible Z decays @ FCC-ee

\Rightarrow explore regimes beyond **cLFV future reach**

\Rightarrow probe **ISS(3,3)** regimes with sizeable
or negligible **cLFV!**

New Physics paths to cLFV: neutrino masses beyond "standard seesaw"



Beyond "standard" seesaw realisations

Seesaw (and its variants) - one of the most appealing mechanisms for m_ν generation

Further ways to account for **tiny neutrino masses**: **higher order, higher dimension!**

Tree \rightarrow n loops \rightarrow

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda_{\text{NP}}} \left[\mathcal{L}_{d=5}^{(0)} + \delta \mathcal{L}_{d=5}^{(1)} + \delta \mathcal{L}_{d=5}^{(2)} + \dots \right]$$

Type I-II-III seesaw (dim5 @tree level) \rightarrow

$$+ \frac{1}{\Lambda_{\text{NP}}^3} \left[\mathcal{L}_{d=7}^{(0)} + \delta \mathcal{L}_{d=7}^{(1)} + \delta \mathcal{L}_{d=7}^{(2)} + \dots \right]$$

$$+ \frac{1}{\Lambda_{\text{NP}}^5} \left[\mathcal{L}_{d=9}^{(0)} + \delta \mathcal{L}_{d=9}^{(1)} + \delta \mathcal{L}_{d=9}^{(2)} + \dots \right]$$

$$+ \frac{1}{\Lambda_{\text{NP}}^7} \left[\mathcal{L}_{d=11}^{(0)} + \delta \mathcal{L}_{d=11}^{(1)} + \delta \mathcal{L}_{d=11}^{(2)} + \dots \right]$$

$$+ \dots$$

Higher order operators:
dim=7, 9, 11...

Higher order contributions:
loop-suppressed ($1/16\pi^2$)

LNV operators dim > 5 \downarrow

Several **other** interesting and theoretically well-motivated possibilities exist:

Tree-level realisations via **higher-dimension** operators, dynamical "seesaws", ...

Higher order realisations (Dirac/Majorana): from first Zee model, to $R_p V$ SUSY, ...

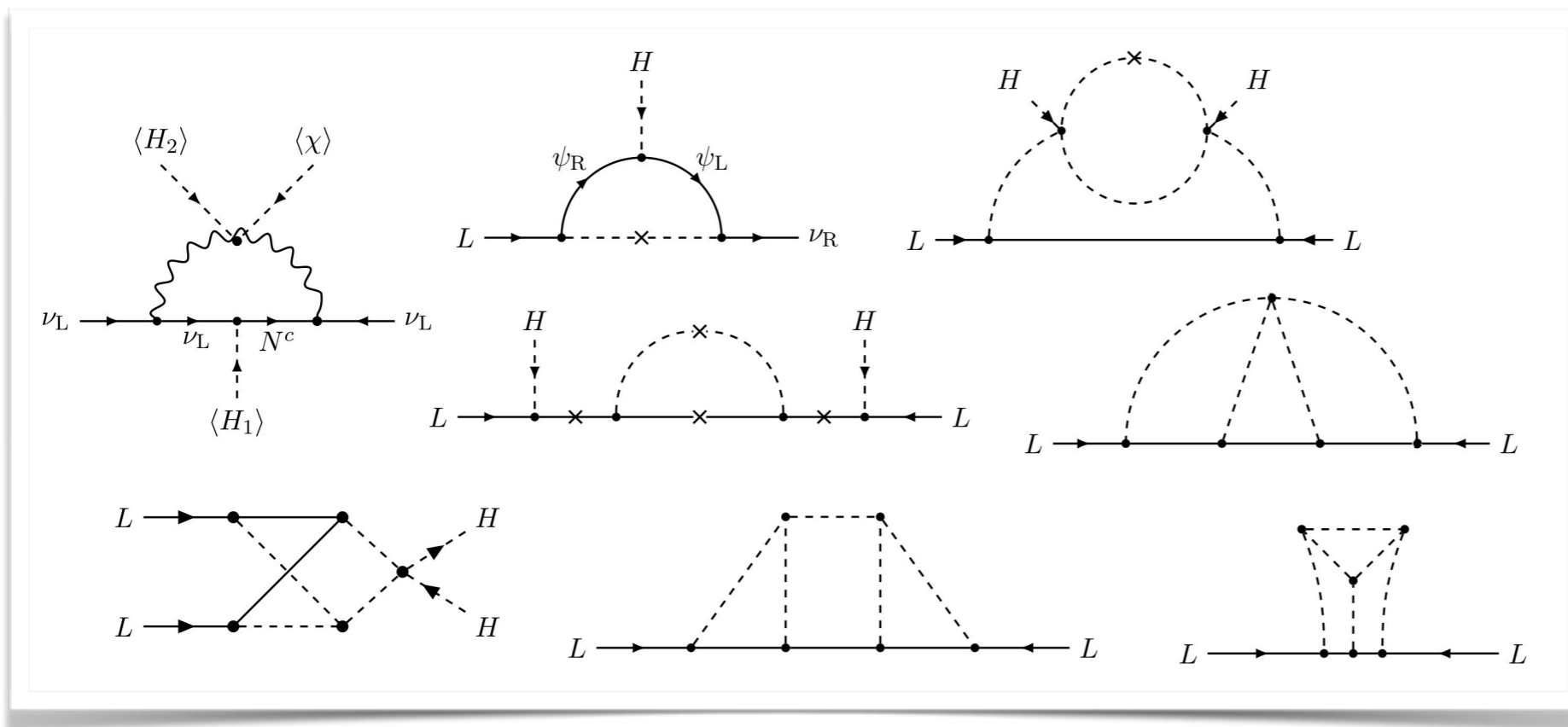
scotogenic models, to **3-loops** and more!

Beyond "standard" seesaw realisations

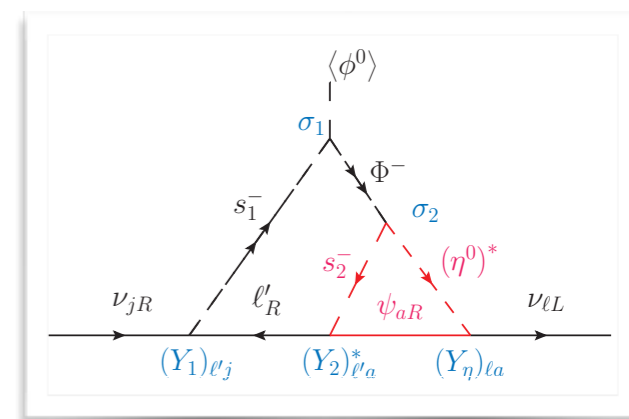
Seesaw (and its variants) - one of the most appealing mechanisms for m_ν generation

Several **other** interesting and theoretically well-motivated possibilities exist:

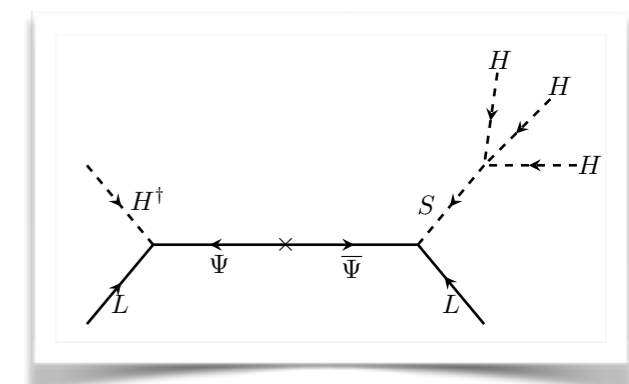
often aiming at addressing m_ν and other **SM observational** issues...



[Cai et al, 1706.08524]



Enomoto et al [1904.07039]



[Arbelaez et al, 2007.11007]

$\Rightarrow m_\nu$ from large **NP scale**, small **couplings**, approximate **symmetries**, **loop** suppression

cLFV and EW precision in scotogenic models



Low-scale models for m_ν : DM connection (?)

Scotogenic models: a link between **neutrino mass** generation and **dark matter!**

[Ma, 2006]

Additional Z_2 **symmetry:** stabilises dark matter candidate ... but

\Rightarrow **neutrino masses @ 1-loop**

[Review on phenomenology of generalised scotogenic models: Hagedorn et al, 1804.04117]

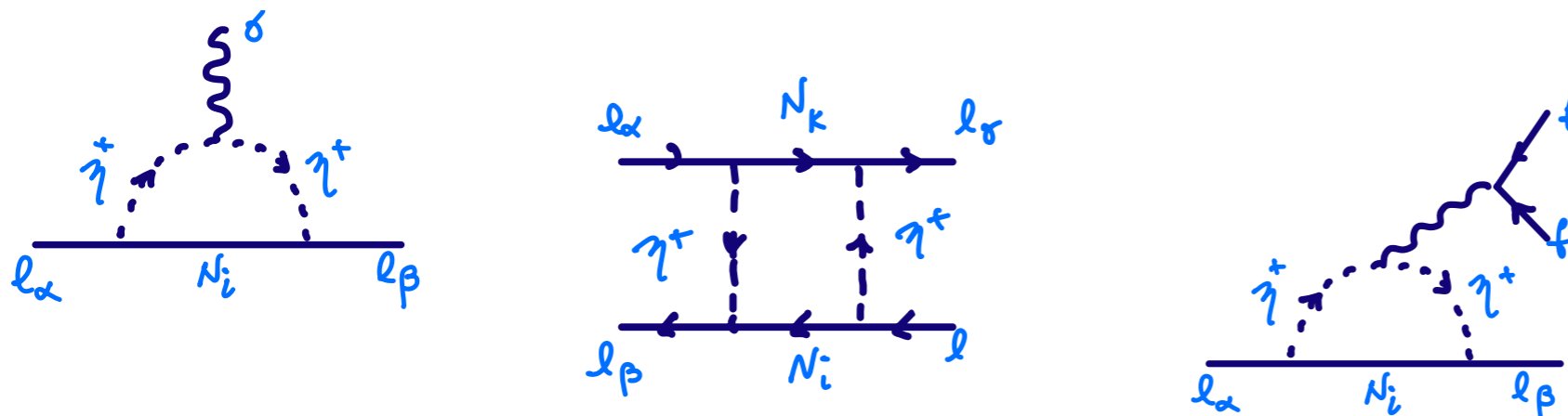
A minimal realisation: extend SM by inert scalar doublet η and RH neutrinos N_R

$$\mathcal{M}_{ij}^\nu \simeq \frac{\lambda_5}{16\pi^2} \frac{2 Y_{ik}^\eta Y_{jk}^\eta v^2}{M_{N_i}^2} \left[\frac{M_{N_k}^2}{m_0^2 - M_{N_k}^2} + \frac{M_{N_k}^4}{(m_0^2 - M_{N_k}^2)^2} \log \left(\frac{M_{N_k}^2}{m_0^2} \right) \right]$$

(for $\lambda_5 \ll 1$, with $m_0 = m_{\eta_R} \sim m_{\eta_1}$)

Suppression of neutrino masses:
smallness of λ_5 and loop factors!

cLFV observables: numerous contributions from η and/or N_R



[Toma and Vicente, 1312.2840]

Low-scale models for m_ν : DM connection (?)

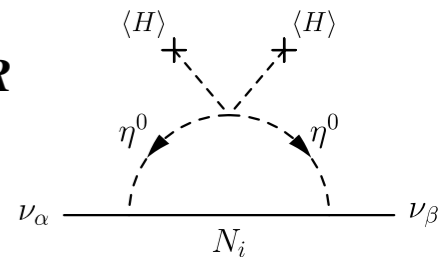
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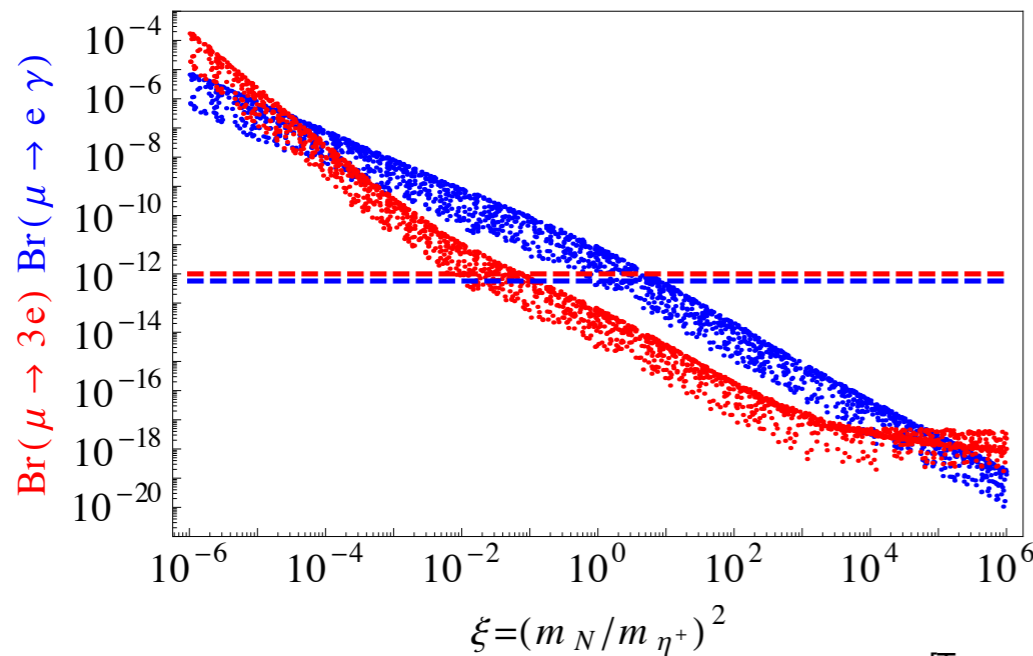
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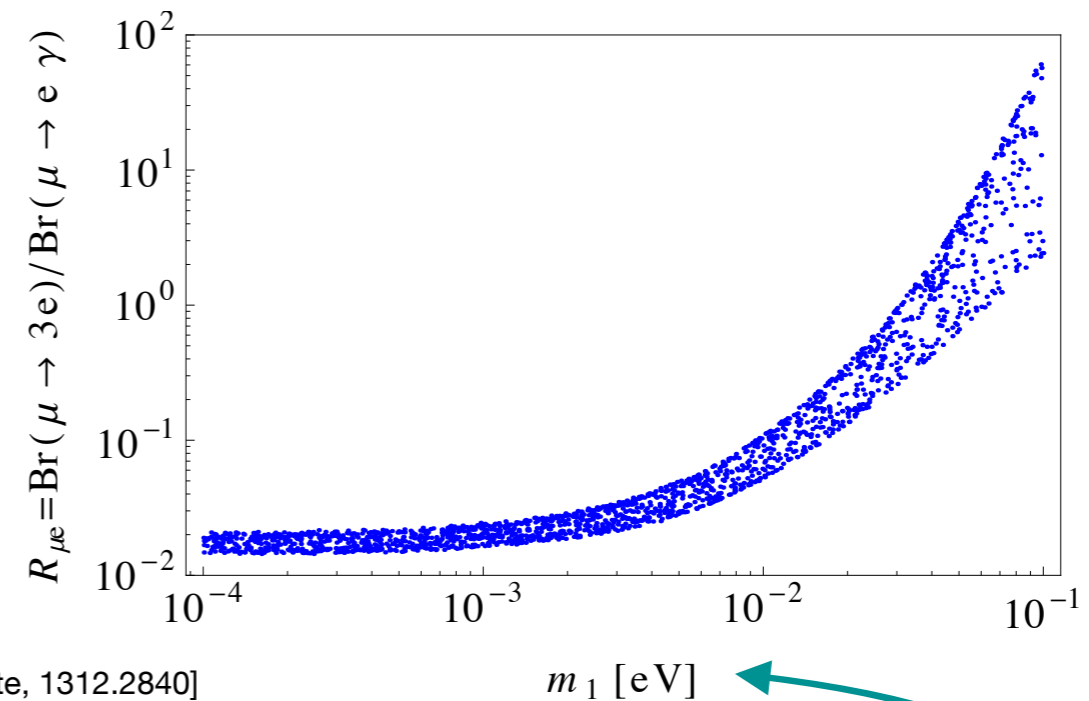
A minimal realisation: extend SM by inert scalar doublet η and RH neutrinos N_R



cLFV observables: hints on the **nature of the DM candidate** (η or N_R) and ν **mass scale**



[Toma and Vicente, 1312.2840]



Determination of $R_{\mu e} = \text{BR}(\mu \rightarrow 3e) / \text{BR}(\mu \rightarrow e\gamma) \Rightarrow$ hints on **lightest neutrino mass** m_{ν_1}

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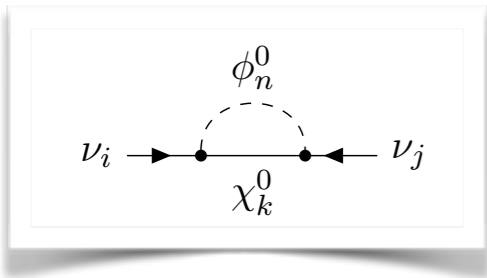
Additional Z_2 **symmetry:** stabilises dark matter candidate ... but

\Rightarrow **neutrino masses @ 1-loop**

"T1-2-A" variant: SM extended by $SU(2)_L$ Weyl fermions, Majorana singlets & extra scalars

\Rightarrow ν **mass generation**, **DM candidates**, $(g-2)_\mu$ and **BAU** via leptogenesis

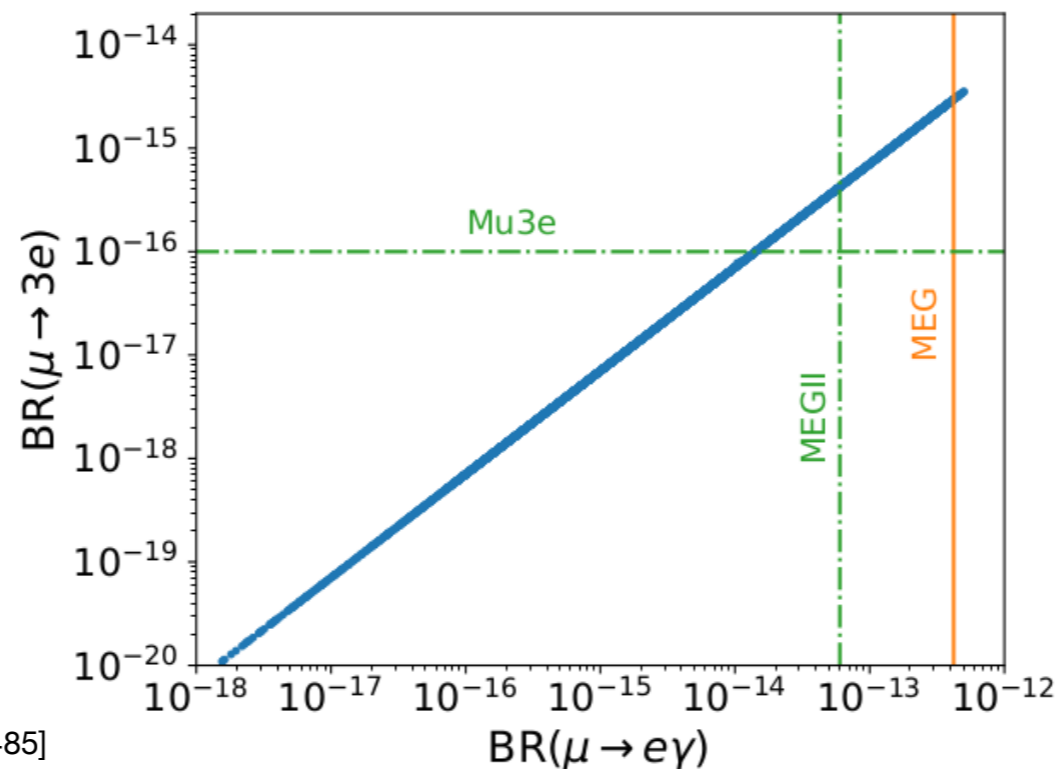
[Alvarez et al, 2301.08485]



cLFV observables - strict correlation between

$BR(\mu \rightarrow e\gamma)$, $BR(\mu \rightarrow 3e)$

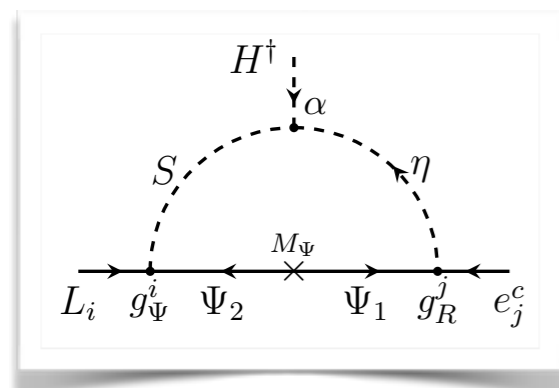
[dipole dominated]



muon cLFV decays

\Rightarrow **falsify model**

@ **MEG II** and **Mu3e** !



[Alvarez et al, 2301.08485]

Low-scale models for m_ν : DM connection (?)



IN2P3
Les deux infinis

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"T1-2-A" variant: SM extended by $SU(2)_L$ Weyl fermions, Majorana singlets & extra scalars

[Alvarez et al, 2301.08485]

Relax certain (driving) assumptions \Rightarrow **ν masses, DM candidates** (scalar and fermionic)

\Rightarrow generic Δa_μ - from *SM-like* to *NP* (at $\sim 5\sigma$) ; no BAU

Avoid **theoretically disfavoured** regimes (large hierarchical "Yukawas" & scalar couplings)?

Thorough exploration of flavoured and electroweak precision observables!

\Rightarrow **cLFV decays:** leptonic, $Z \rightarrow \ell_\alpha \ell_\beta$, $H \rightarrow \ell_\alpha \ell_\beta$

\Rightarrow **EW observables:** sensitive probes of new interactions (scalar, vector, fermion...)

$Z \rightarrow \text{inv}$, LFUV in $Z \rightarrow \ell_\alpha \ell_\alpha$ (*in progress*)

New contributions to $H \rightarrow \ell_\alpha \ell_\alpha$

[Darricau, Lee, Orloff, AMT to appear soon]

Full computation of **NLO** contributions to Z and Higgs interactions!

[see e.g. Grimus et al, 0802.4353]

Low-scale models for m_ν : DM connection (?)

Scotogenic models: a link between neutrinos

generation and dark matter!

Stabilises dark matter candidate ... but

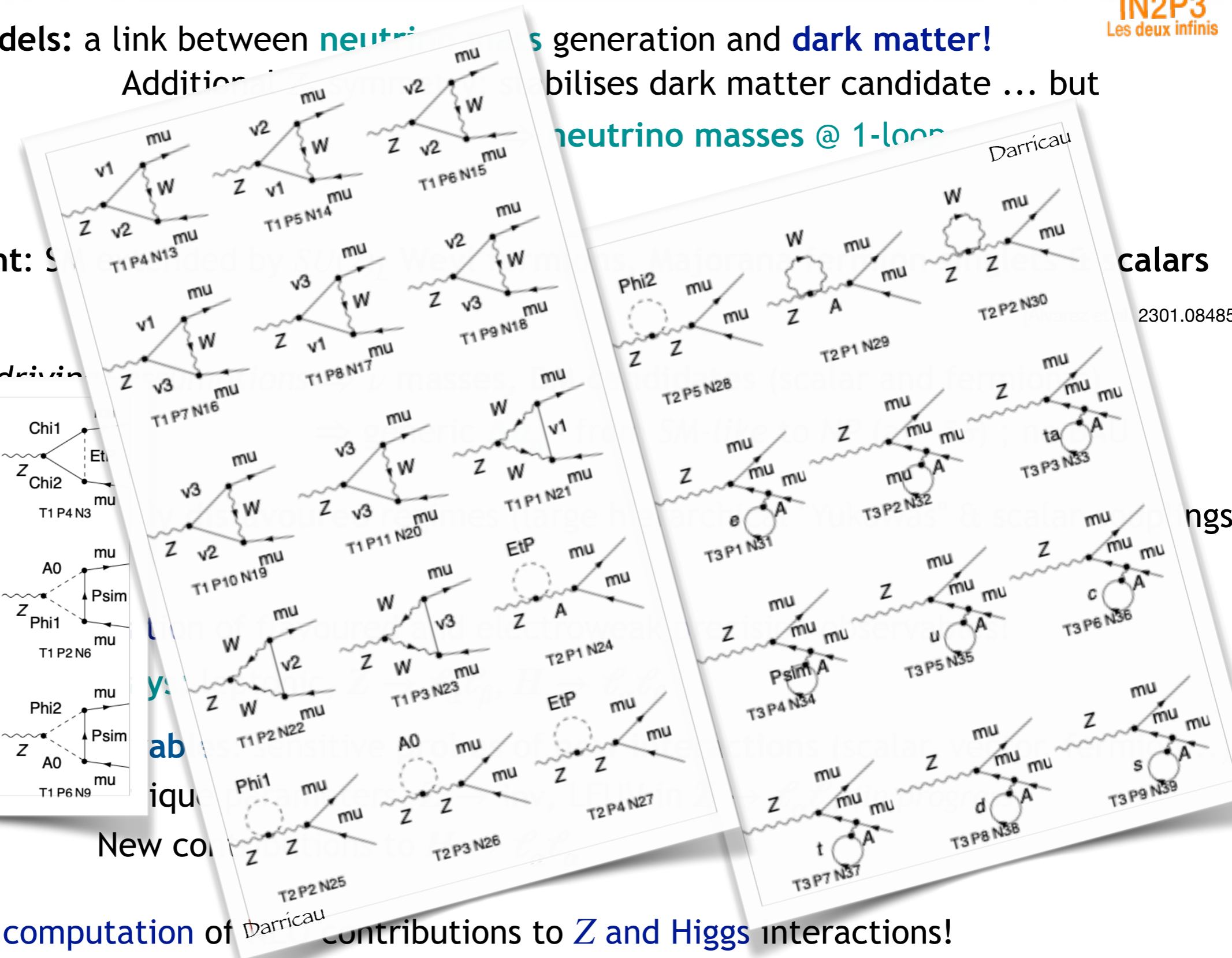
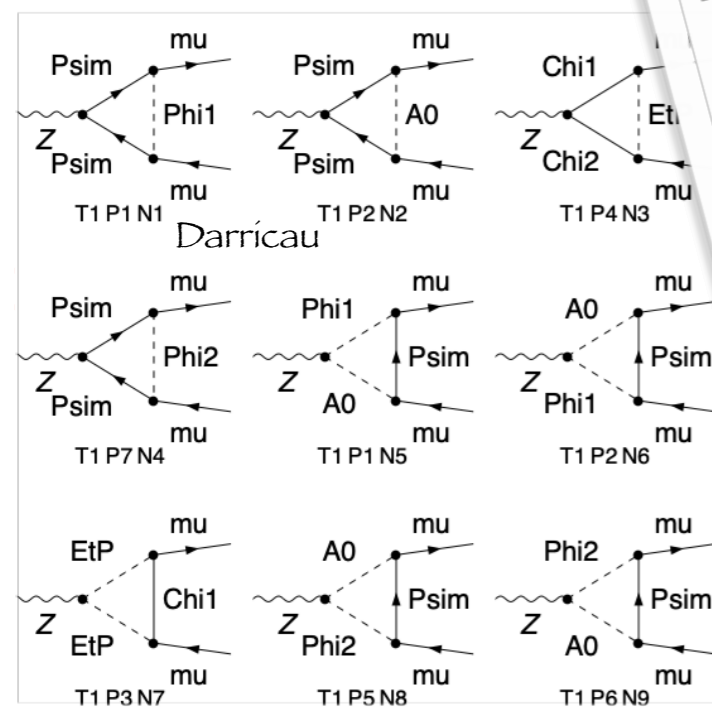
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[Ma, 2006]

Additional

"T1-2-A" variant: S

Delay certain (drivin



calars

2301.08485]

ngs)?

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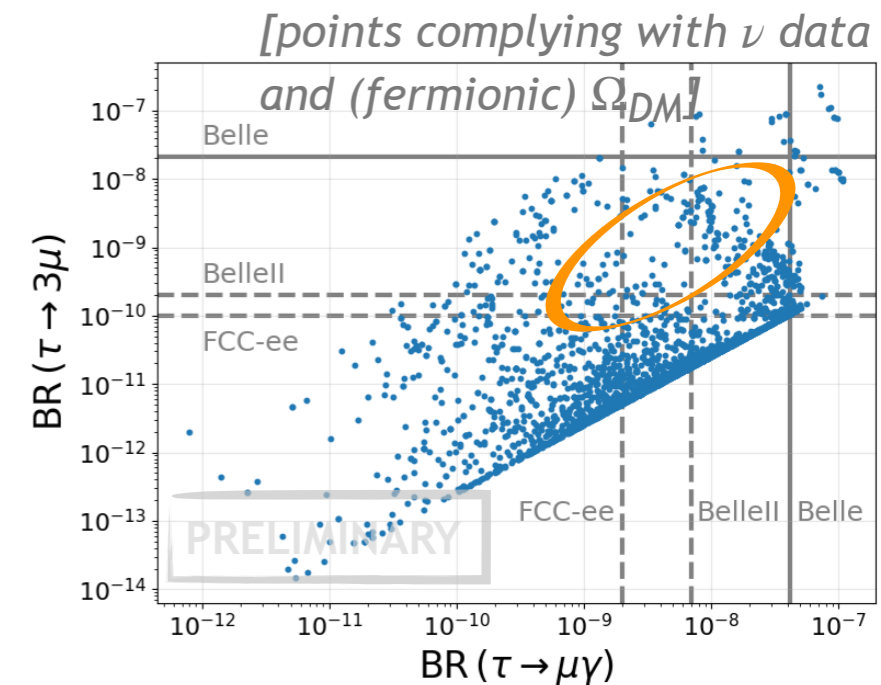
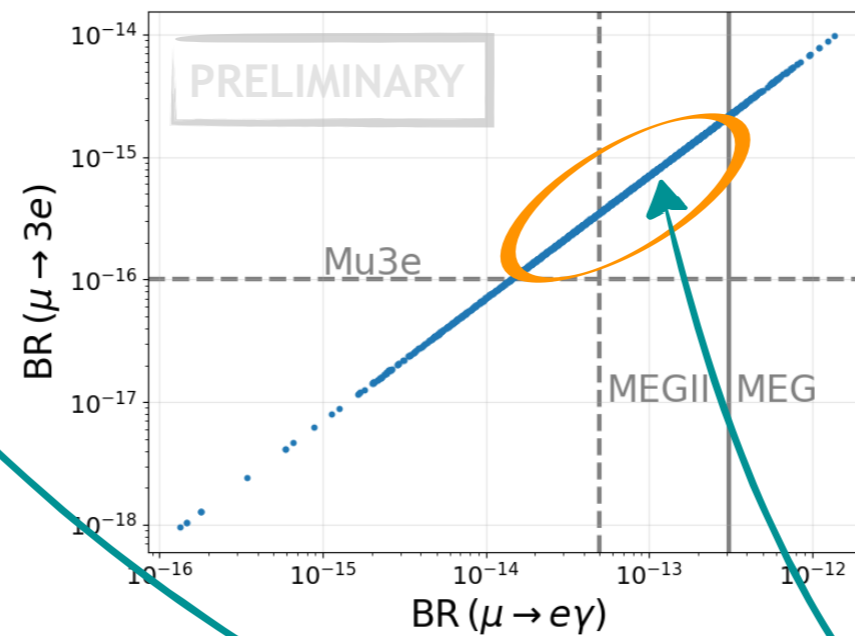
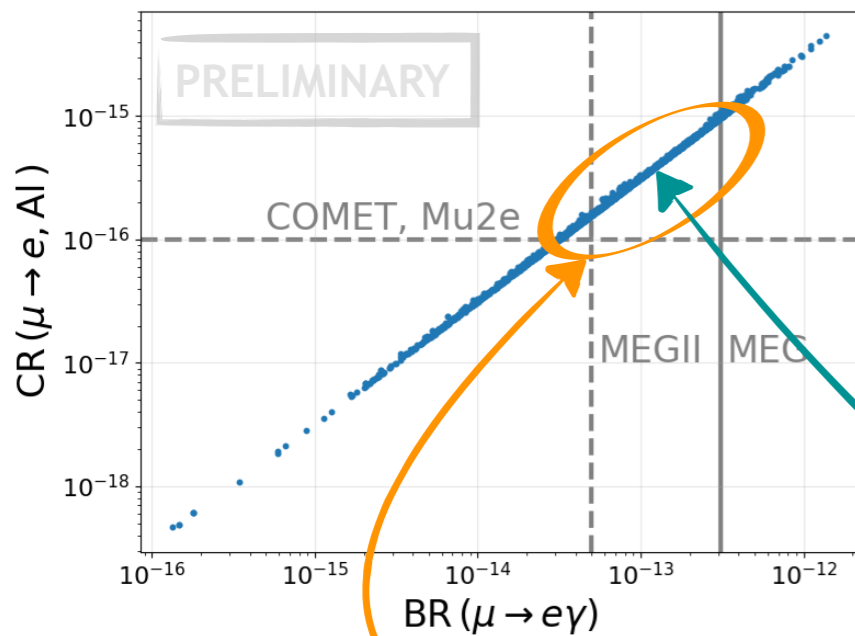
Low-scale models for m_ν : DM connection (?)

► Scotogenic models: a link between **neutrino mass** generation and **dark matter**!

► cLFV leptonic decays:

⇒ correlated **muon-electron** decays (dipole dominance)

⇒ sizeable **box-contributions** for **3-body tau** decays



⇒ well within **future reach**!

⇒ **muon cLFV decays** \rightsquigarrow **test & falsify model @ MEG II, Mu3e, Mu2e & COMET**

[Darricau, Lee, Orloff, AMT, to appear soon]

Low-scale models for m_ν : DM connection (?)

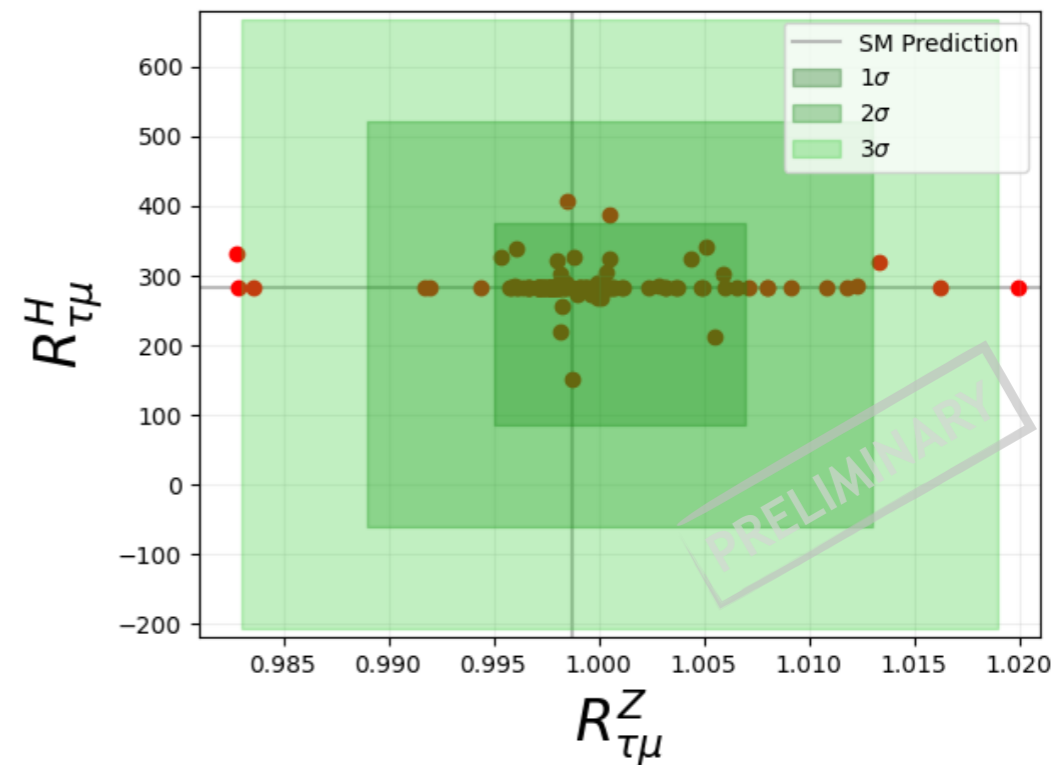
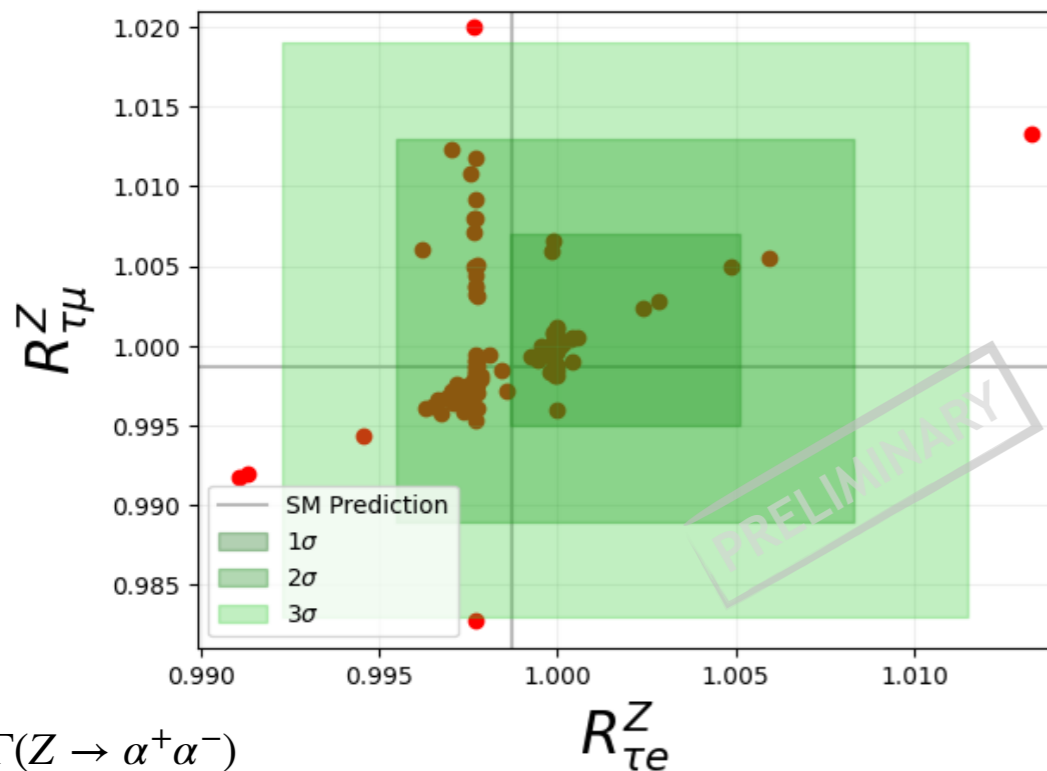
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muon cLFV decays \leadsto test model @ MEG II, Mu3e, Mu2e & COMET

► LFUV in Z decays (and Higgs)

\Rightarrow Typically within 2σ (although certain regimes in stronger tension)



[Darricau, Lee, Orloff, AMT, to appear soon]

$$R_{\alpha\beta}^Z = \frac{\Gamma(Z \rightarrow \alpha^+\alpha^-)}{\Gamma(Z \rightarrow \beta^+\beta^-)}$$

Low-scale models for m_ν : DM connection (?)

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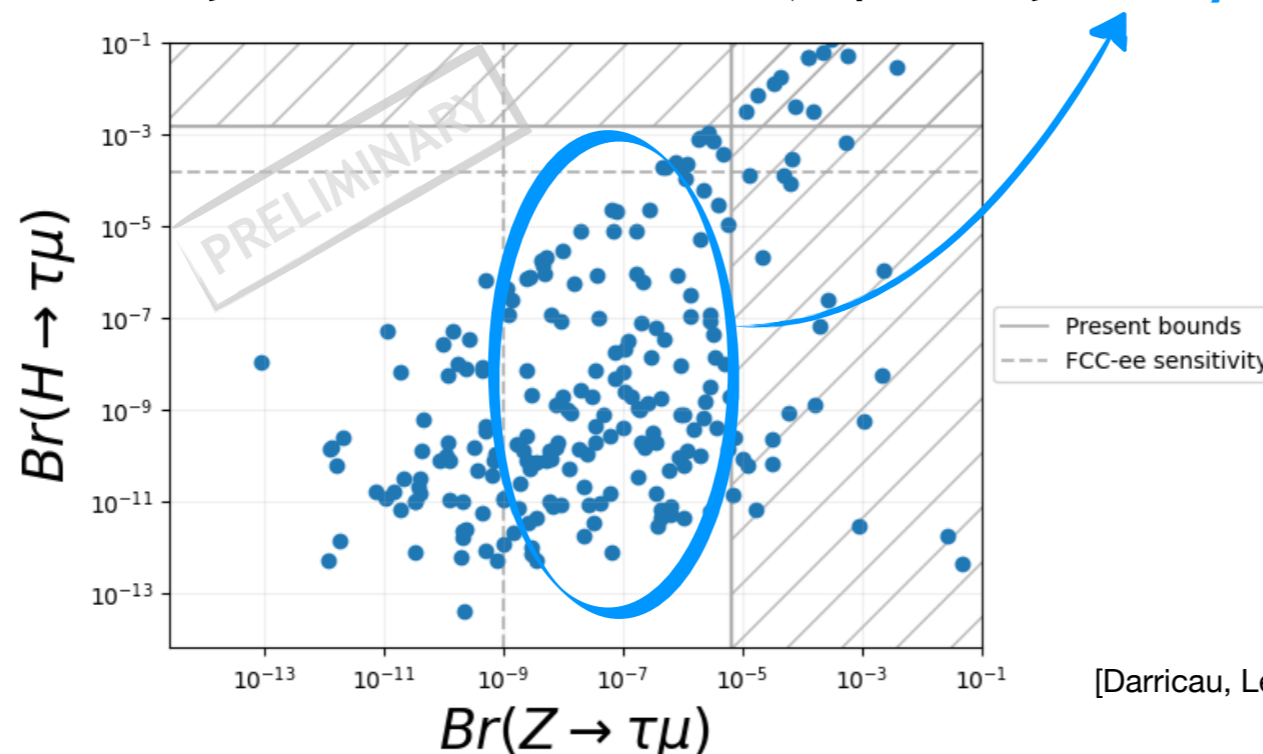
► LFUV in Z decays (and Higgs)

\Rightarrow Typically within 2σ (although certain regimes in stronger tension)

► cLFV in Z and Higgs decays (and Higgs)

\Rightarrow Constraints on certain regimes!

\Rightarrow Potentially testable at FCC-ee (especially $Z \rightarrow \mu\tau$)



[Darricau, Lee, Orloff, AMT, to appear soon]

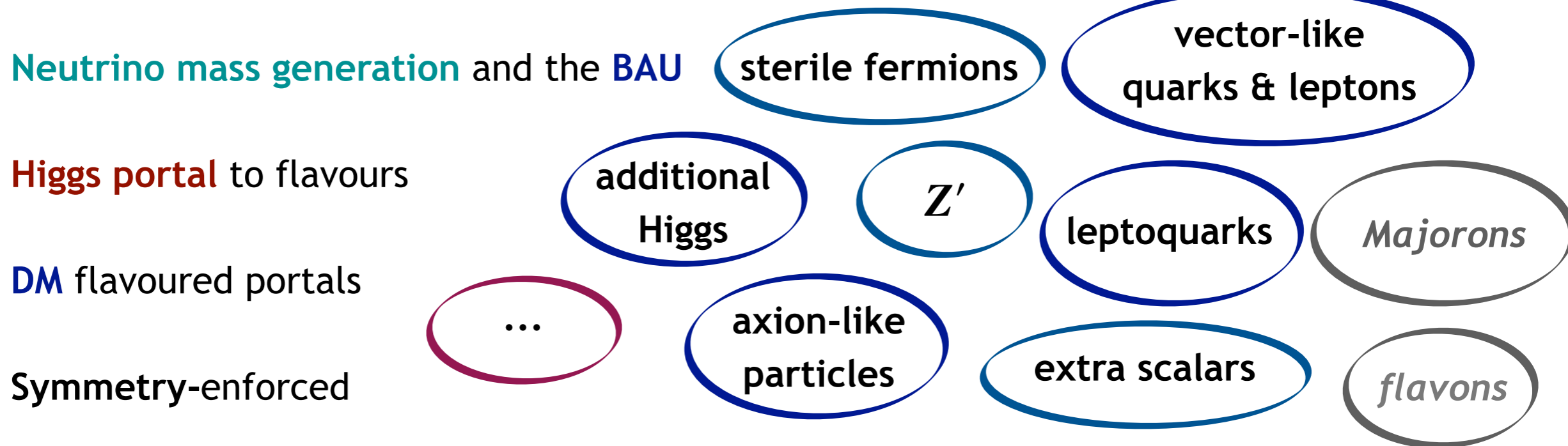
cLFV beyond (minimal) neutrino masses...



NP models of flavour: so many possibilities!

Extensive contributions in recent years - driven by **NP hints** (m_ν , AMMs, B-anomalies...)

⇒ exploring *flavoured signatures of BSM* realisations



UV-complete models: GUTs, Supersymmetry, extra dimensions, ...

Ultimately addressing all (several) SM problems, and testable (via flavours?!) ??

Can lepton flavours help us disentangle the NP model at work?

Or falsify candidates?

NP models of flavour: so many possibilities!

Extensive contributions in recent years - driven by NP hints (m_ν , AMMs, B-anomalies...)

⇒ exploring *flavoured signatures of BSM* realisations

	AC	RVV2	AKM	δ LL	FBMSSM	LHT	RS
$D^0 - \bar{D}^0$	★★★★	★	★	★	★	★★★★	?
ϵ_K	★	★★★★	★★★★	★	★	★★	★★★★
$S_{\psi\phi}$	★★★★	★★★★	★★★★	★	★	★★★★	★★★★
$S_{\phi K_S}$	★★★★	★★	★	★★★★	★★★★	★	?
$A_{CP}(B \rightarrow X_s \gamma)$	★	★	★	★★★★	★★★★	★	?
$A_{7,8}(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★★★★	★★★★	★★	?
$A_9(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★	★	★	?
$B \rightarrow K^{(*)} \nu \bar{\nu}$	★	★	★	★	★	★	★
$B_s \rightarrow \mu^+ \mu^-$	★★★★	★★★★	★★★★	★★★★	★★★★	★	★
$K^+ \rightarrow \pi^+ \nu \bar{\nu}$	★	★	★	★	★	★★★★	★★★★
$K_L \rightarrow \pi^0 \nu \bar{\nu}$	★	★	★	★	★	★★★★	★★★★
$\mu \rightarrow e \gamma$	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★
$\tau \rightarrow \mu \gamma$	★★★★	★★★★	★	★★★★	★★★★	★★★★	★★★★
$\mu + N \rightarrow e + N$	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★
d_n	★★★★	★★★★	★★★★	★★	★★★★	★	★★★★
d_e	★★★★	★★★★	★★	★	★★★★	★	★★★★
$(g-2)_\mu$	★★★★	★★★★	★★	★★★★	★★★★	★	?

- AC: RH currents & U(1) flavour sym
- RVV2: SU(3)-flavoured MSSM
- AKM: RH currents & SU(3) family sym
- δ LL: CKM-like currents
- FBMSSM: flavour-blind MSSM
- LHT: Little Higgs (T-parity)
- RS: Warped extra dimensions

Expected impact for observables:

- ★★★★ large effects
- ★★ small, but visible
- ★ unobservable

[Altmannshofer et al, '10]

Densely populated sector!

cLFV transitions amongst the most sensitive observables to numerous NP models!

► Embed **4dim space-time** into **5dim AdS space** (extra dim compactified on orbifold)

Two branes (UV, IR) and bulk between them

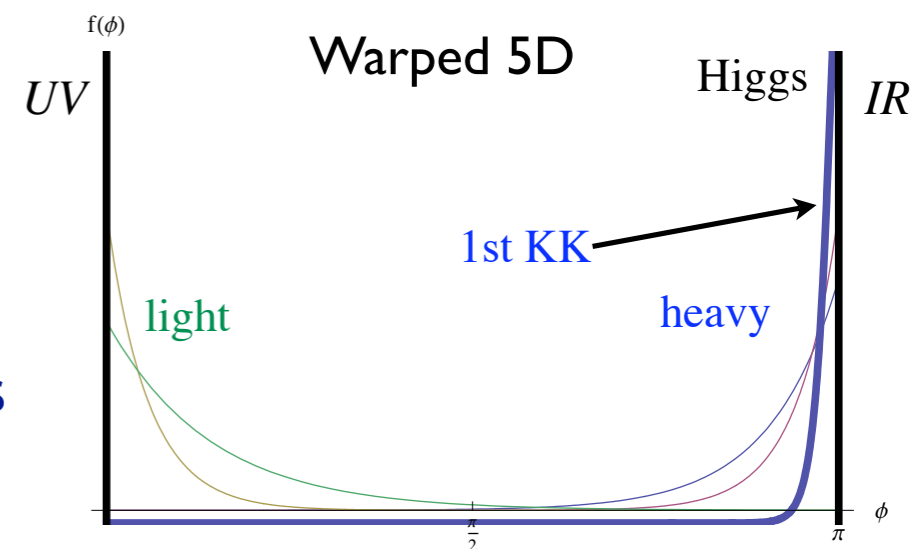
Localise fields: **Higgs** close to *IR brane*

SM fermions & **gauge bosons** on *bulk*

KK excitations close to *IR brane*

Interactions \leftrightarrow **overlap of wave functions**

(L)FV from **couplings** of light fermions to **KK excitations**



$$M_{\text{TeV}} \simeq M_{\text{Planck}} e^{-\pi k L_5}$$

Geometrical distribution of **fermions in bulk**:

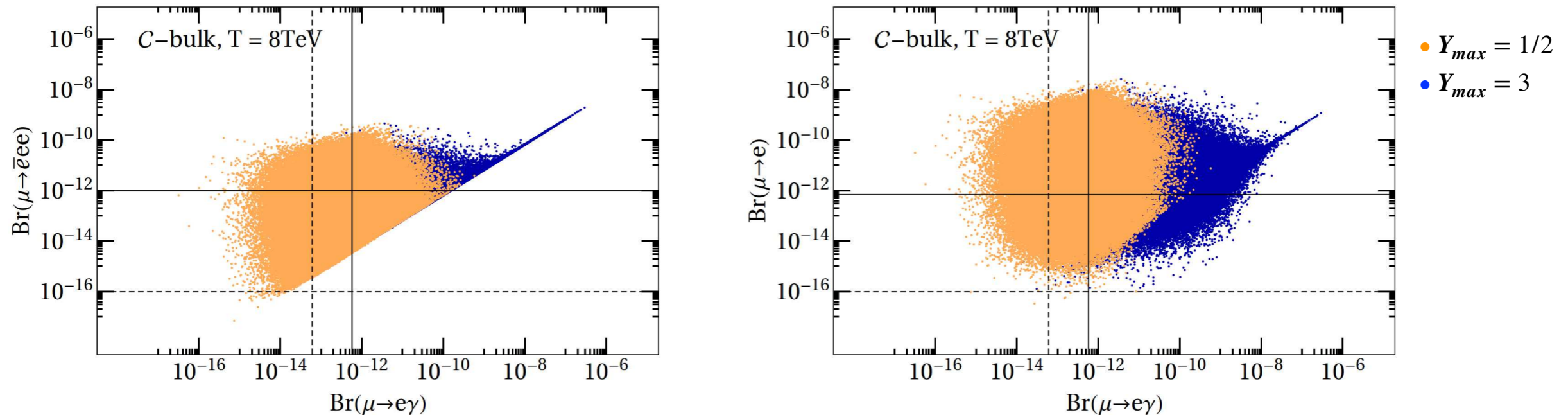
reproduce hierarchy in **4dim Yukawas** from "**anarchic**" $\mathcal{O}(1)$ **dim5 couplings!**

Non-negligible phenomenological issues:

enlarge bulk symmetry to prevent violation of custodial SU(2) symmetry
 additional "rescue" strategies to avoid excessive FCNCs,
 to protect EW precision observables, ..., among other issues

[Burdman, '02; Agashe et al, '04; Csaki et al, '08; Blanke et al & Buras et al, '08-'09;
 Bauer et al, '10; Vempati et al, '12; Beneke et al, '12-'15; and many others...]

- ▶ **Example: custodially protected** model, full inclusion of **all dim-6 cLFV** operators
 generical anarchic Yukawa couplings
 new gauge fields & KK-excitations of lepton fields \Rightarrow **cLFV transitions**



[Beneke et al, 1508.01705]

Most stringent constraints from $\mu \rightarrow e\gamma$ and $\mu - e$ conversion

τ decays comparatively less restrictive

Current $\mu - e$ cLFV bounds constrain NP scale to be very heavy, **beyond LHC reach**

$$T_{KK} \gtrsim 4 \text{ TeV} \quad (\text{corresponding to } m_{KK}^1 \gtrsim 10 \text{ TeV})$$

Future $\mu - e$ sensitivities: exclude anarchic RS models (without additional symmetries)

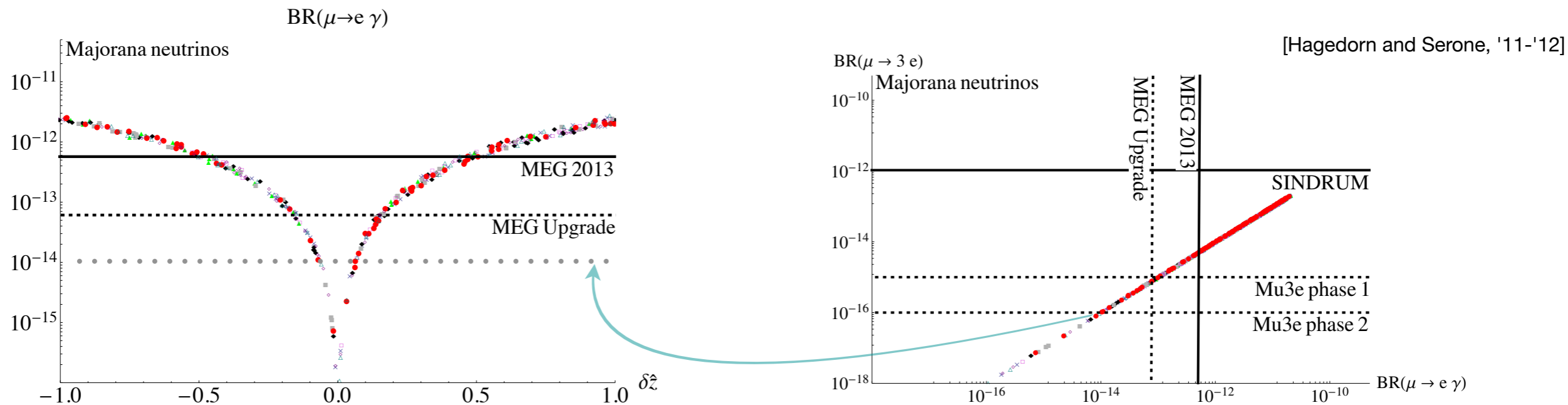
up to **8 TeV** (corresponding to KK gluon masses around 20 TeV)

► **Holographic composite Higgs** model based on enlarged symmetry, $\mathcal{G}_{SM} \times G_f$

$$G_f = X \times Z_N, \text{ with } X = S_4, A_4, \Delta(96,384)$$

(Discrete) symmetries - predict the **lepton mixing pattern** (masses unconstrained)

Applied to **5dim** model in warped space; both cases of **Dirac and Majorana** neutrinos



cLFV observables (as well as **EDMs**) typically below experimental bounds ($m_{KK}^1 \sim 3 - 4$ TeV)

MEG (I & II) bounds on $\mu \rightarrow e \gamma$ \rightsquigarrow constrain the **size of boundary kinetic terms!**

Important role played in the future by **Mu3e data**

\Rightarrow **cLFV** allows to infer relevant **information on fundamental parameters**

Concluding remarks



Confirmed observations and several "tensions" suggest the need to go **beyond the SM**

In the **lepton sector**, ν -masses provided the 1st laboratory **evidence of NP**

Many experimental "tensions" nested in **lepton-related observables**

Lepton physics might offer valuable hints in **constructing and probing NP models**

New Physics can be manifest via **cLFV, LNV, ...** even before any **direct discovery!**

(Synergy of) lepton observables can provide information on the underlying NP model

New Physics is there! **Lepton physics** might be a perfect portal to address SM problems

⇒ First hints on preferred paths to NP from **EFT approach**

⇒ Attempt at identifying the **underlying model** capable of accounting

for **all SM problems** (m_ν , DM and BAU) and further "tensions" with observation!

cLFV emerges as extremely **powerful probe** to test and falsify **NP in the lepton sector**

Explore **different paths**, and profit from amazing **experimental prospects** in the near future!

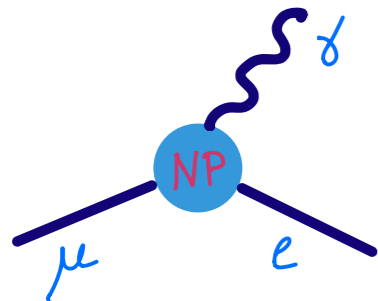
Additional material



cLFV observables



cLFV muon channels: radiative decays



► cLFV decay: $\mu^+ \rightarrow e^+ \gamma$

► Event signature: $E_e = E_\gamma = m_\mu/2$ (~ 52.8 MeV)

Back-to-back $e^+ - \gamma$ ($\theta \sim 180^\circ$); Time coincidence

► Backgrounds \Rightarrow prompt physics & accidental

Prompt: radiative μ decays ($\mu \rightarrow e \bar{\nu}_e \nu_\mu \gamma$, very low E_ν)

[$\propto R_\mu$]

Accidental: coincidence of γ with positron from Michel decays $\mu \rightarrow e \bar{\nu}_e \nu_\mu$:

photon from $\mu \rightarrow e \bar{\nu}_e \nu_\mu \gamma$; γ from in-flight e^+e^- annihilation

[$\propto R_\mu^2$]

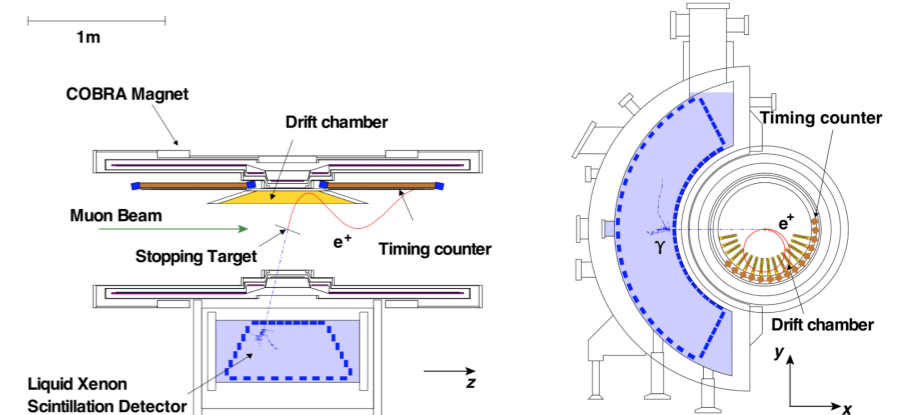
► Experimental status:

First searches (!) in 1940's

[MEG Coll., 1605.05081]

Advent of intense muon beams in 2000's MEG @ PSI

$BR(\mu^+ \rightarrow e^+ \gamma) \leq 4.2 \times 10^{-13}$ (90% CL)



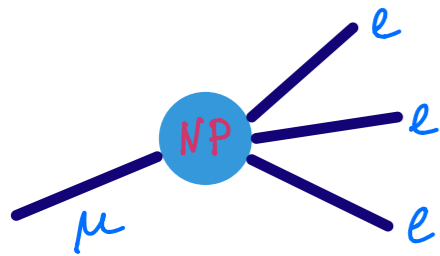
► Future prospects:

[MEG II Coll., 2201.008200]

MEG II (@ PSI): $BR(\mu^+ \rightarrow e^+ \gamma) \leq 6 \times 10^{-14}$

very hard to go beyond 10^{-15} without conceptually different approach

cLFV muon channels: 3-body decays



► cLFV decay: $\mu^+ \rightarrow e^+ e^- e^+$

► Event signature: $\Sigma E_e = m_\mu; \Sigma \vec{P}_e = \vec{0}$

common vertex; Time coincidence

► Backgrounds \Rightarrow physics & accidental

Physics: multi-body μ decays ($\mu \rightarrow e \bar{\nu}_e \nu_\mu e^+ e^-$, very low E_ν)

Accidental: Bhabha scattering of Michel e^+ from $\mu \rightarrow e \bar{\nu}_e \nu_\mu$ decays with atomic $e^+ e^-$
Michel positrons with $e^+ e^-$ from γ conversion

► Experimental status:

SINDRUM @ PSI

[SINDRUM Coll., '88]

$$BR(\mu^+ \rightarrow e^+ e^- e^+) \leq 1.0 \times 10^{-12} \quad (90\% \text{ CL})$$

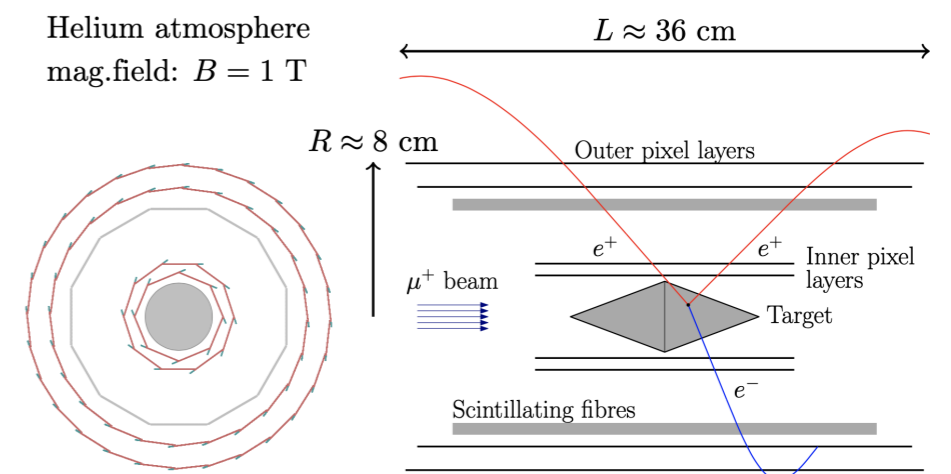
► Future prospects:

[Mu3e Coll., 2009.11690]

Mu3e (@ PSI): expected sensitivity $\mathcal{O}(10^{-15})$ for Phase I

with HIMB, $\mathcal{O}(10^{-16})$ for Phase II

[Aiba et al, 2111.05788]



cLFV in muonic atoms: $\mu - e$ conversion

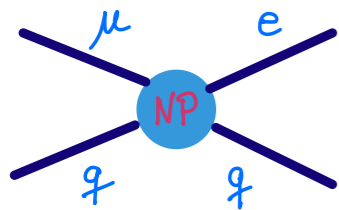
► **Muonic atoms:** 1s bound state formed when μ^- stopped in target

SM allowed processes: decay in orbit (DIO) $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$

nuclear capture $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$

► In the presence of New Physics - **cLFV neutrinoless $\mu^- - e^-$ conversion**

$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$



► **Event signature:** single mono-energetic electron

$$E_{\mu e} = m_\mu - E_B(A, Z) - E_R(A, Z)$$

For Aluminium, Lead, Titanium $\sim E_{\mu e} \approx \mathcal{O}(100 \text{ MeV})$

Which target? ** For coherent conversion, maximal rates for $30 \leq Z \leq 60$

► **Backgrounds** \Rightarrow **Only physics!** μ decay in orbit, beam purity, cosmic rays, ...

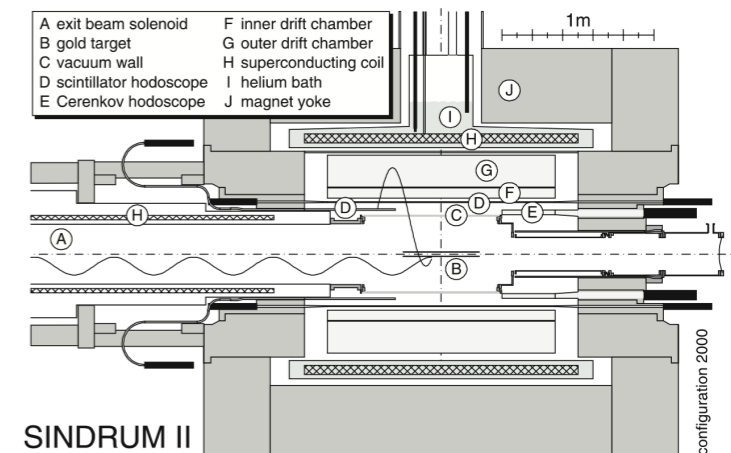
► **Experimental status:** [SINDRUM II Coll., '06]

$$\text{SINDRUM @ PSI: } CR(\mu^- - e^-, \text{Au}) \leq 7.1 \times 10^{-13} \text{ (90\% CL)}$$

► **Future prospects:**

$$\text{Mu2e (@ FNAL) - } \mathcal{O}(10^{-17}), \text{ [Bartoszek et al, 1501.05241]}$$

$$\text{[Abramishvili et al, '20] COMET (@ JPARC) - } \mathcal{O}(10^{-15} - 10^{-17}), \dots$$



cLFV in muonic atoms: $\mu - e$ conversion

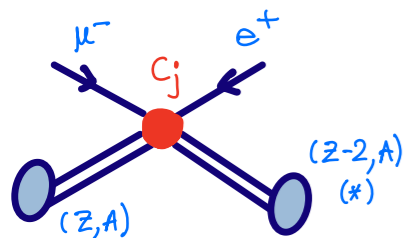
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nuclear capture $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$

► In the presence of New Physics - **cLFV & LNV** ($\Delta L = 2$) **neutrinoless $\mu^- - e^+$ conversion**

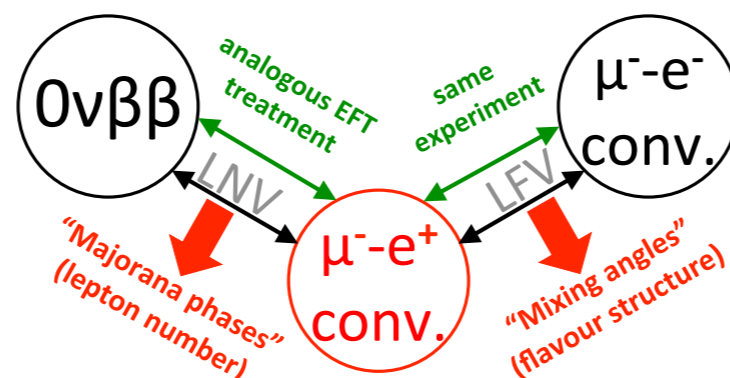
$$\mu^- + (A, Z) \rightarrow e^+ + (A, Z - 2)^*$$



$\mu^- - e^-$ conversion: coherent process, single nucleon, nuclear ground states

$\mu^- - e^+$ conversion: 2 nucleons ($\Delta Q = 2$), possibly excited final state

A unique connection between **LNV** (in association with **Majorana** nature and possibly, neutrino mass generation) and **cLFV**



LNV-Alternatives:
 $\mu^- - \mu^+$ conversion
 $K^+ \rightarrow \pi^+ \mu^- \mu^-$

LFV-Alternatives:
 $\mu \rightarrow e + \gamma$
 $\mu \rightarrow 3e$

[see e.g. Geib et al, 1609.09088]

cLFV in muonic atoms: $\mu - e$ conversion

► **Muonic atoms:** 1s bound state formed when μ^- stopped in target

SM allowed processes: decay in orbit (DIO) $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$

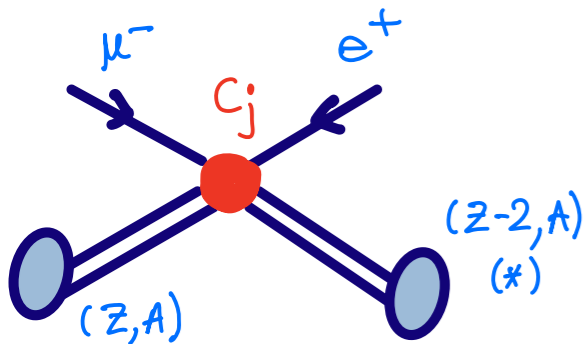
nuclear capture $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$

► In the presence of New Physics - **cLFV & LNV** ($\Delta L = 2$) **neutrinoless $\mu^- - e^+$ conversion**

$$\mu^- + (A, Z) \rightarrow e^+ + (A, Z - 2)^*$$

$\mu^- - e^-$ conversion: coherent process, single nucleon, nuclear ground states

$\mu^- - e^+$ conversion: 2 nucleons ($\Delta Q = 2$), possibly excited final state



► **Event signature:** single positron - but complex energy spectrum

$$E_{\mu e}^{N^*} = m_\mu - E_B(A, Z) - E_R(A, Z) - \Delta_{Z-2}^{(*)}$$

For Aluminium (giant dipole resonance) $\sim E_{\mu-e^+}^{\text{Al, GDR}} \approx \mathcal{O}(83.9 \text{ MeV})$

► **Experimental status:**

Collaboration	year	Process	Bound
PSI/SINDRUM	1998	$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}^*$	3.6×10^{-11}
PSI/SINDRUM	1998	$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}$	1.7×10^{-12}

► **Future prospects:**

Best sensitivity expected for Ca, S and Ti targets (possibly $\sim \mathcal{O}(\text{few} \times 10^{-15})$)

▶ **Muonic atoms:** 1s bound state formed when μ^- stopped in target

▶ In the presence of **New Physics** - **cLFV muonic atom decay** $\mu^- e^- \rightarrow e^- e^-$

Initial μ^-, e^- : 1s states bound in **Coulomb field** of muonic atom's nucleus

Coulomb interaction increases wave function overlap
rate strongly enhanced in **large Z** atoms, $\Gamma \gtrsim (Z - 1)^3$
Larger phase space (compared with $\mu \rightarrow 3e$)

▶ **Event signature:** back-to-back electrons, $E_{e^-} \approx m_\mu/2$

▶ **Backgrounds** \Rightarrow similar to neutrinoless conversion

▶ **Experimental status** - **new** observable!

possibly included in future physics runs (e.g. COMET)

cLFV muonium decays

► Muonium: $\mu^+ e^-$

Hydrogen-like Coulomb bound state, free of hadronic interactions!
Powerful laboratory for EW tests and cLFV

► In the presence of New Physics - Muonium oscillations and Muonium decays

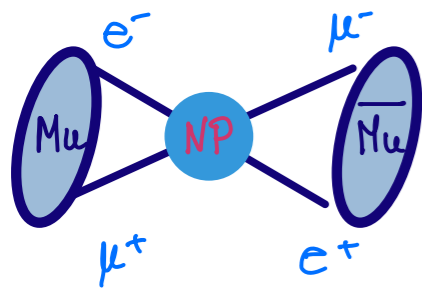
► Mu- $\overline{\text{Mu}}$ oscillation

Spontaneous conversion $\mu^+ e^- \leftrightarrow \mu^- e^+$

Reflects a double (individual) lepton number violation $|\Delta L_e| = |\Delta L_\mu| = 2$

Rate (typically) suppressed by external magnetic fields

Detection: reconstruct Michel electron from μ^- decays and shell positron



Experimental status: MACS - $P(\text{Mu} - \overline{\text{Mu}}) < 8.3 \times 10^{-11}$ [Willmann et al, 1999]

Future prospects: MACE, AMF (@FNAL)

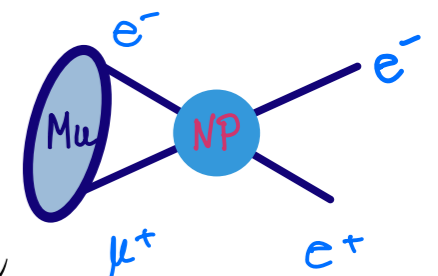
[Bai et al, 2203.11406]

► Mu decays

$$\mu^+ e^- \rightarrow e^+ e^-$$

Clear signal compared to SM-allowed muonium decay, $\text{Mu} \rightarrow e^+ e^- \bar{\nu}_\mu \nu_e$

No available bounds, no clear roadmap...



cLFV tau decays: leptonic and more

Tau leptons - heaviest of all charged leptons! *Cannot have "intense tau beams"* 😊

Copious production at B-factories (BaBar, Belle, LHCb, Belle II, ...)

Production and decay: $e^+ e^- \rightarrow \tau^+ \tau^-$ signal "hemisphere"
tagging "hemisphere" (e.g. $\tau^+ \rightarrow \bar{\nu}_\tau \nu_e e^+$)

► cLFV tau decays: abundant modes! Pure leptonic, semileptonic (2- and 3-body), ...

► Radiative decay: $\tau^\pm \rightarrow \ell^\pm \gamma$

► **Event signature:** $E_{\text{final}} - \sqrt{s}/2 = \Delta E \sim 0$; $M_{\text{final}} = M_{\ell\gamma} \sim m_\tau$

► **Backgrounds** \Rightarrow coincidence of isolated leptons with γ (ISR, FSR); mistagging

► 3-body leptonic decay: $\tau^\pm \rightarrow \ell_i^\pm \ell_j^\mp \ell_k^\pm$

► **Event signature:** $E_{3\ell} - \sqrt{s}/2 \sim 0$; $M_{3\ell} \sim m_\tau$

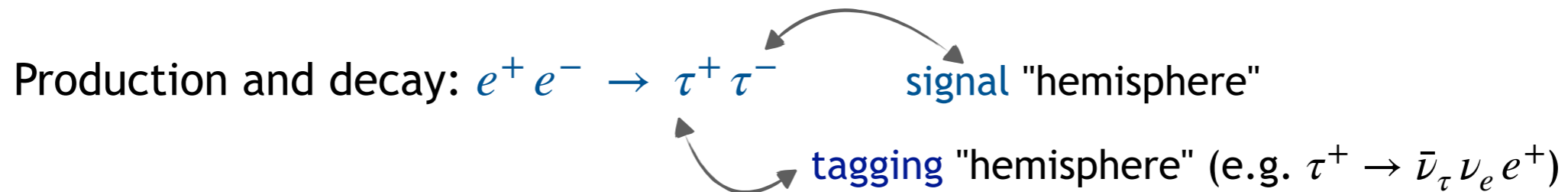
► **Backgrounds** \Rightarrow No irreducible backgrounds!

Small background from $q\bar{q}$ and Bhabha pairs, ...

cLFV tau decays: leptonic and more

Tau leptons - heaviest of all charged leptons! *Cannot have "intense tau beams"* :)

Copious production at B-factories (BaBar, Belle, LHCb, Belle II, ...)



► cLFV tau decays: abundant modes! Pure leptonic, semileptonic (2- and 3-body), ...

► Semi-leptonic cLFV tau decays

2-body final state: $\tau \rightarrow \ell h^0$ (pseudoscalar, scalar or vector neutral meson)

3-body final state: $\tau \rightarrow \ell h_i h_j$ ($h \leftrightarrow \pi^\pm, K^\pm, K_s^0$)

► cLFV exotic modes (also lepton & baryon number violating)

$\tau^- \rightarrow \ell^+ h_i^\pm h_j^\pm$ ($h \leftrightarrow \pi^\pm, K^\pm$) \Rightarrow LNV

$\tau^- \rightarrow \Lambda h^-$ ($h \leftrightarrow \pi^\pm, K^\pm$) \Rightarrow LNV & BLV

$\tau \rightarrow p \ell_i \ell_j$ \Rightarrow LNV & BLV