



The gradient flow extended to the Standard Model

Janosch Borgulat, Jonas Kohnen, Robert Harlander
RWTH Aachen University



Institut für
Theoretische
Teilchenphysik und
Kosmologie



Why gradient flow?

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Energy-momentum tensor $T_{\mu\nu}^{\text{YM}} = G_{\mu\rho}^a G_{\nu\rho}^a - \frac{1}{4} \delta_{\mu\nu} G_{\rho\sigma}^a G_{\rho\sigma}^a$

❖ Energy density: $\epsilon \sim \langle T_{00} \rangle$, Pressure: $P \sim \langle T_{ii} \rangle$

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❖ EMT = Conserved current

$$\begin{aligned} x_\mu &\rightarrow x_\mu + \epsilon_\mu \\ \Rightarrow \partial_\mu T_{\mu\nu} &= 0 + \dots \\ \Rightarrow T_{\mu\nu} &= \text{fin.} \end{aligned}$$

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... for hadrons? \longrightarrow hadronic energies \longrightarrow lattice

❖ EMT = Conserved current

❖ Lattice **breaks translational invariance**

\Rightarrow EMT not finite on lattice

[Suzuki 2013]

~~$x_\mu \rightarrow x_\mu + \epsilon_\mu$~~
 $\Rightarrow \partial_\mu T_{\mu\nu} = 0 + \dots$
 $\Rightarrow T_{\mu\nu} = \text{fin.}$

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❖ **Solution?** *Remove UV divergence of composite operators...*

❖ UV divergence \longrightarrow high fluctuations / momenta

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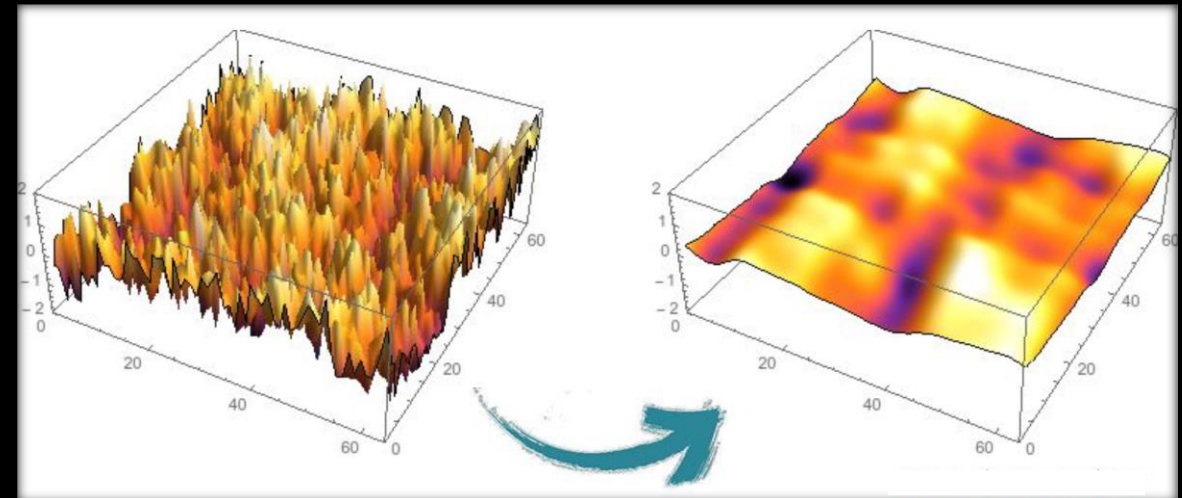
❖ **Solution?** *Remove UV divergence of composite operators...*

❖ UV divergence \longrightarrow high fluctuations / momenta

❖ Need **low-pass filter** for QCD

\longrightarrow cutoff?

\longrightarrow breaks Lorentz invariance

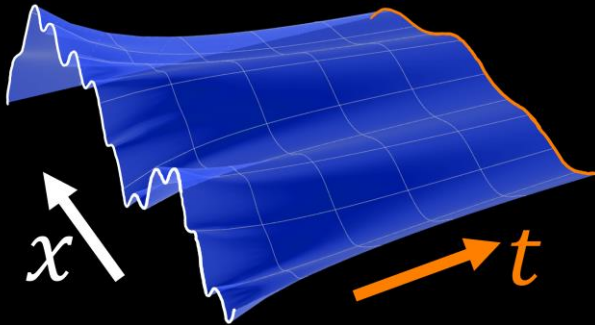


Gradient flow: Definition

❖ Flow equation:
(Heat equation)

$$\partial_t \chi(t, x) = \mathcal{D}^2(t) \chi(t, x), \quad \chi(0, x) = \psi(x)$$

$$\partial_t B_\mu(t, x) = \mathcal{D}_\nu(t) G_{\nu\mu}(t), \quad B_\mu(0, x) = A_\mu(x)$$



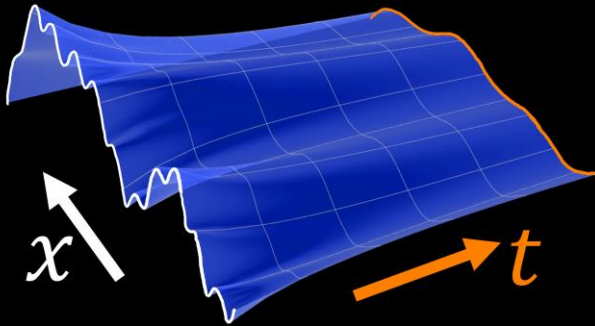
[Narayanan, Neuberger 2006; Lüscher 2010; Lüscher, Weisz 2011; Lüscher 2013]

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[Narayanan, Neuberger 2006; Lüscher 2010; Lüscher, Weisz 2011; Lüscher 2013]

❖ LO solution: $\partial_t \chi(t, x) = \partial^2 \chi(t, x) \Rightarrow$

$$\hat{\chi}(t, p) = \hat{\psi}(p) e^{-tp^2}$$

Gradient flow: Definition

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu} + \bar{\psi}(\not{D} + m)\psi$$

[Lüscher, Weisz 2011]

$$\mathcal{L}_{\text{flowed}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}_{\chi} + \mathcal{L}_{\text{B}}$$

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$$\int_0^{\infty} dt \bar{\lambda} (\partial_t \chi - \mathcal{D}^2 \chi)$$

$$s, j \xrightarrow{p} t, i = \delta_{ij} \frac{-i\not{p} + m}{p^2 + m^2} e^{-(t+s)p^2}$$

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$$\int_0^{\infty} dt L_{\mu} (\partial_t B_{\mu} - \mathcal{D}_{\nu} G_{\nu\mu})$$

$$s, j \xrightarrow{p} t, i = \delta_{ij} \frac{-i\not{p} + m}{p^2 + m^2} e^{-(t+s)p^2}$$

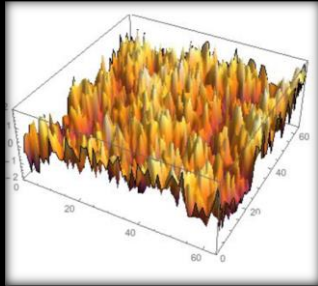
$$s, \nu, b \xrightarrow{p} t, \mu, a = \frac{\delta^{ab}}{p^2} \left(\delta_{\mu\nu} - \frac{p_{\mu} p_{\nu}}{p^2} \right) e^{-(t+s)p^2}$$

+ vertices & mixed propagators

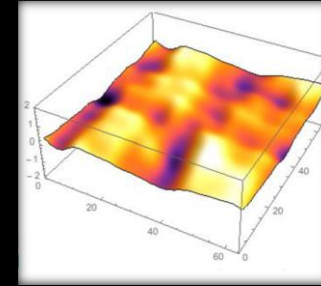
Application

Energy-momentum tensor

$$T_{\mu\nu}^{\text{YM}} = G_{\mu\rho}^a G_{\nu\rho}^a - \frac{1}{4} \delta_{\mu\nu} G_{\rho\sigma}^a G_{\rho\sigma}^a$$



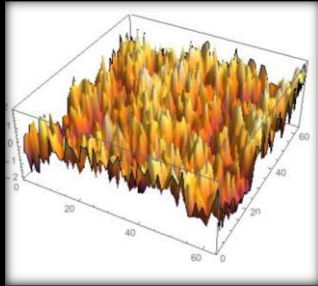
$$G_{\mu\nu}(x) \xrightarrow{A(x) \rightarrow B(t, x)} \mathcal{G}_{\mu\nu}(t, x)$$



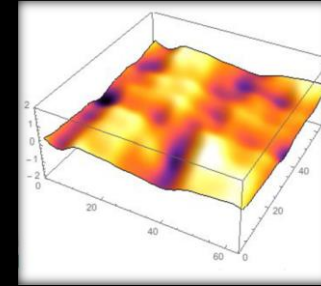
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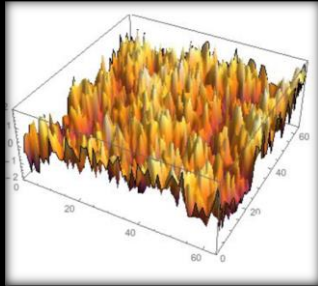
Short flow-time expansion (SFTX):
[Lüscher, Weisz 2011]

$$\tilde{\mathcal{O}}_i(t, x) = \zeta_{ij}(t) \mathcal{O}_j(x) + \dots$$

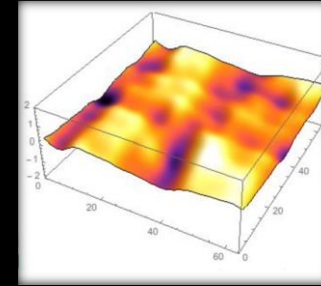
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Short flow-time expansion (SFTX):

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$$\tilde{\mathcal{O}}_i(t, x) = \zeta_{ij}(t) \mathcal{O}_j(x) + \dots$$

$$T_{\mu\nu}^{\text{YM}}(x) = c_1(t) \left[\mathcal{G}_{\mu\rho} \mathcal{G}_{\nu\rho}(t, x) \right] + c_2(t) \left[\delta_{\mu\nu} \mathcal{G}_{\rho\sigma} \mathcal{G}_{\rho\sigma}(t, x) \right]$$

Perturbative

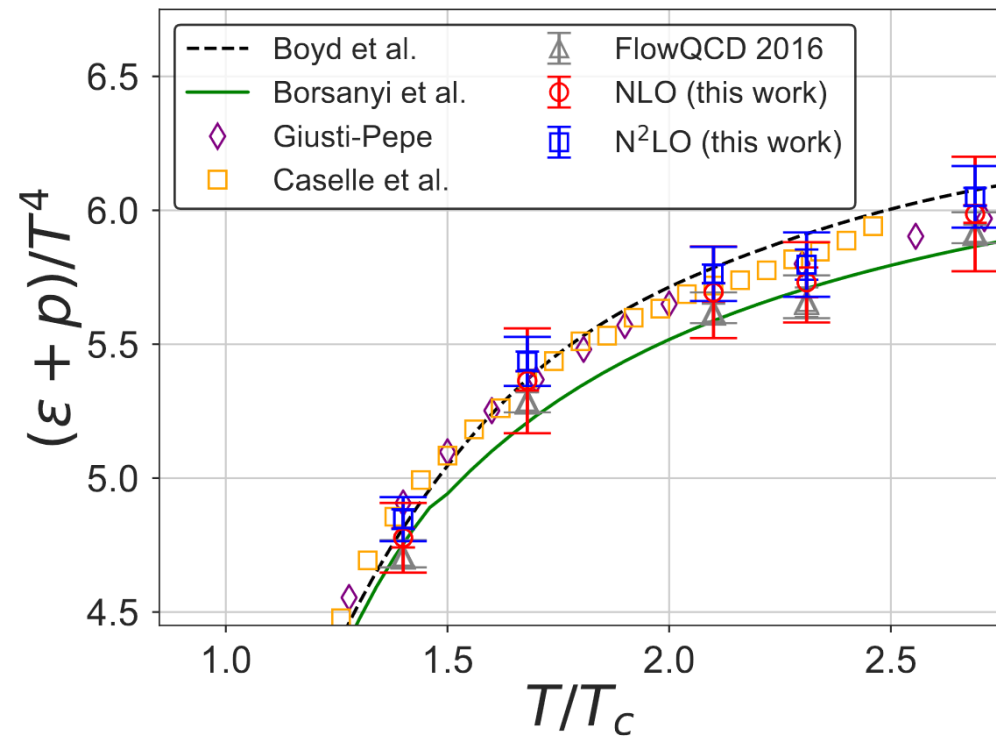
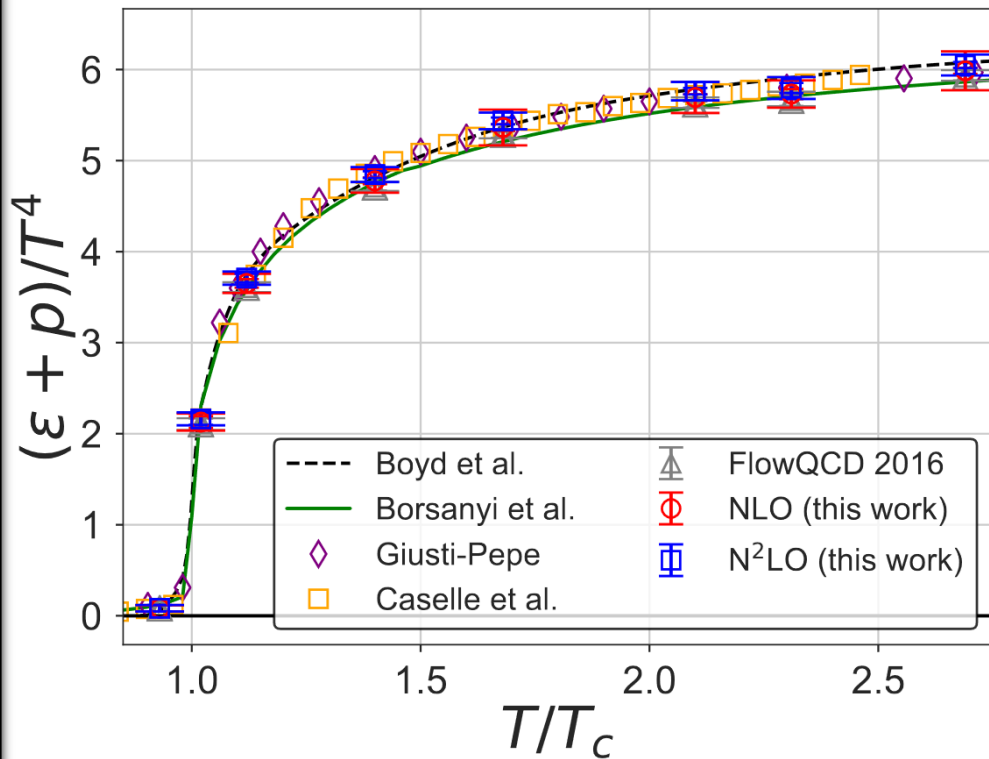
Lattice: *No renormalization required!*

[Suzuki 2013]

[Harlander, Kluth, Lange 2018]

Application

2-loop PT + lattice



Entropy density: $\varepsilon + p = -\frac{4}{3} \left\langle T_{00}(x) - \frac{1}{4} T_{\mu\mu}(x) \right\rangle$

[Iritani, Kitazawa, Suzuki, Takaura 2019]



Other Applications?

Other Applications!

- ❖ Strong coupling
- ❖ Renormalization of EFTs
- ❖ Neutron electric dipole moment
- ❖ Hadronic vacuum polarization
- ❖ Meson mixing and lifetimes
- ❖ QCD static force

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Applications

Strong coupling

[Lüscher 2010; Lüscher 2014]

[Harlander, Neumann 2016]

$$\langle G_{\mu\nu}^a G_{\mu\nu}^a \rangle = 0 \longrightarrow E(t) = \frac{1}{4} \mathcal{G}_{\mu\nu}^a(t) \mathcal{G}_{\mu\nu}^a(t)$$

Applications

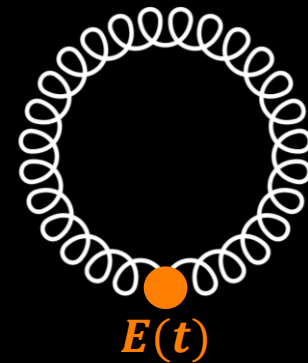
Strong coupling

[Lüscher 2010; Lüscher 2014]

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$$\langle t^2 E(t) \rangle = \frac{3\alpha_s}{4\pi} \frac{N_A}{8} + \mathcal{O}(\alpha_s^2)$$



A Feynman diagram showing a circular loop of gluons (represented by curly lines). An external gluon line, also represented by a curly line, is attached to the bottom of the loop at a vertex marked with an orange dot. Below this vertex is the label $E(t)$.

$$\sim \int \frac{d^D p}{(4\pi)^2} e^{-tp^2} \neq 0$$

Applications

Strong coupling

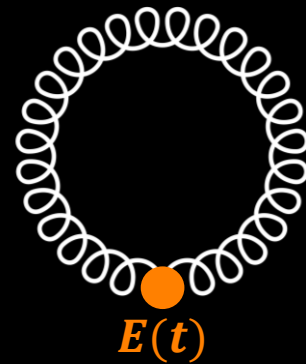
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$$\alpha_s = \frac{32\pi}{3N_A} \langle t^2 E(t) \rangle + \dots$$

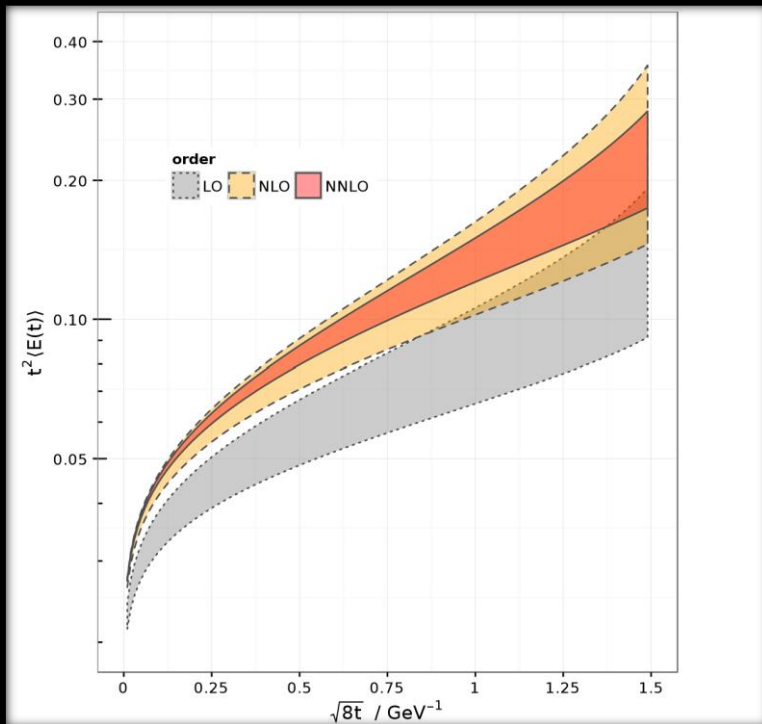
Lattice



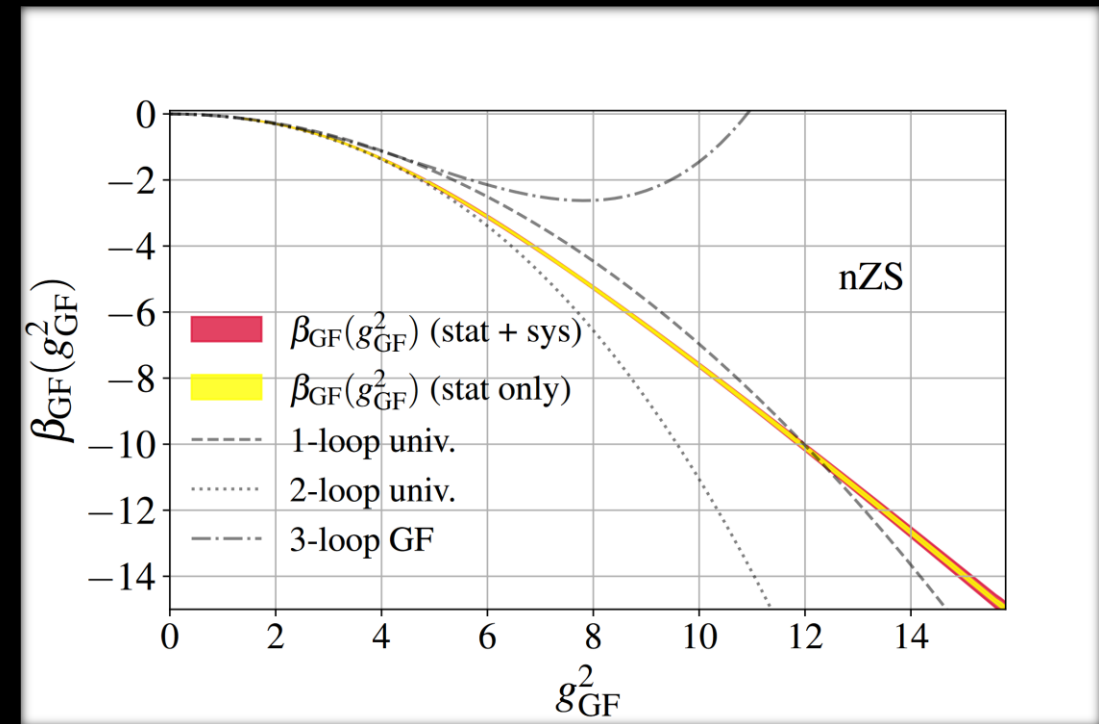
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Applications

Strong coupling



$N_f = 3, \quad 3/\sqrt{8t} < \mu < 1.15/\sqrt{8t}$
[Harlander, Neumann 2016]



$N_f = 0$
[Hasenfratz, Peterson, van Sickle, Witzel 2023]

Applications

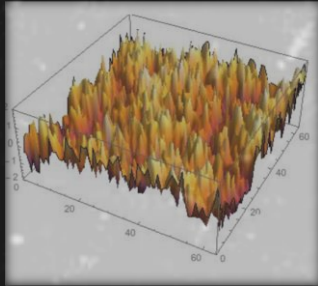
Renormalizing EFTs

Let's recall our approach to the EMT ...

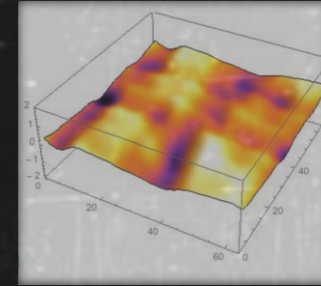
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Perturbative

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Applications

Renormalizing EFTs

Let's have a closer look: $\tilde{\mathcal{O}}(t, x) = \zeta(t)\mathcal{O}(x)$

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$$\tilde{\mathcal{O}}(t) = \zeta^{\textcircled{B}}(t, \epsilon) \mathcal{O}^{\textcircled{B}}(\epsilon) = \underbrace{\zeta^B(t, \epsilon) Z^{-1}(\epsilon)}_{\zeta^R(t) \stackrel{!}{=} \text{fin.}} \mathcal{O}^R \Rightarrow \boxed{Z(\epsilon)}$$

Applications

Renormalizing EFTs

Let's have a closer look:

$$\tilde{\mathcal{O}}(t) = \zeta^B(t, \epsilon) \mathcal{O}^B(\epsilon)$$

?

Applications

Renormalizing EFTs

Let's have a closer look: $\tilde{\mathcal{O}}(t) = \zeta^B(t, \epsilon) \mathcal{O}^B(\epsilon)$

Method of Projectors

[Gorishny, Larin, Tkachov 1983; Gorishny, Larin 1987]

$$P[X] \sim \Pi \left(\frac{\partial}{\partial p_i^\nu}, \frac{\partial}{\partial m} \right) \langle A_\mu^a(p_1) \psi(p_2) \dots | X | 0 \rangle \Big|_{p_i=m=0}$$

Only tree-level contributions at $t=0$

Applications

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Only tree-level contributions at $t=0$

$$P_i[\mathcal{O}_j] = \delta_{ij} \quad \Rightarrow \quad \zeta_{ij}^B(t, \epsilon) = P_j[\tilde{\mathcal{O}}_i(t)]$$

Higher orders for $t > 0$

Applications

Renormalizing EFTs

$$\mathcal{O} \rightarrow \tilde{\mathcal{O}}(t) \rightarrow \zeta^B(t, \epsilon) \rightarrow Z(\epsilon)$$

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Renormalizing EFTs

$$\mathcal{O} \rightarrow \tilde{\mathcal{O}}(t) \rightarrow \zeta^B(t, \epsilon) \rightarrow Z(\epsilon)$$

No need to separate UV and IR divergences!

$$0 = \frac{1}{\epsilon_{UV}} + \frac{1}{\epsilon_{IR}} \sim \int d^D k \frac{1}{(k^2)^n}$$

flow

$$\int d^D k \frac{e^{-tk^2}}{(k^2)^n} \sim \frac{1}{\epsilon_{IR}} = -\frac{1}{\epsilon_{UV}}$$

$Z(\epsilon)$ cancels IR divergences

Applications

Renormalizing EFTs

Flowed standard model?

- ❖ Known: $SU(N)$ + fermions
- ❖ $SU(3) \otimes SU(2)$
- ❖ $U(1) \longrightarrow$ linear flow equation
- ❖ Higgs
- ❖ All the flowed renormalization constants

$$B_{\mu}^B = B_{\mu}, \quad \chi^B = Z_{\chi}^{1/2} \chi$$

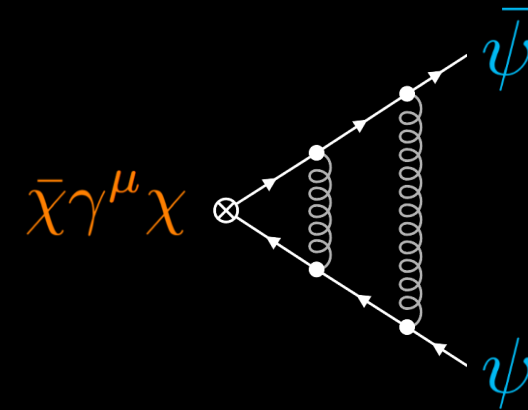
[Lüscher 2011; Lüscher 2013]

Applications

Renormalizing EFTs

How to find Z_χ ?

- ❖ Vector currents $\bar{\psi}\gamma^\mu\psi$ finite
see also: [JB, Harlander, Kohnen, Lange 2023]
- ❖ No operator renormalization

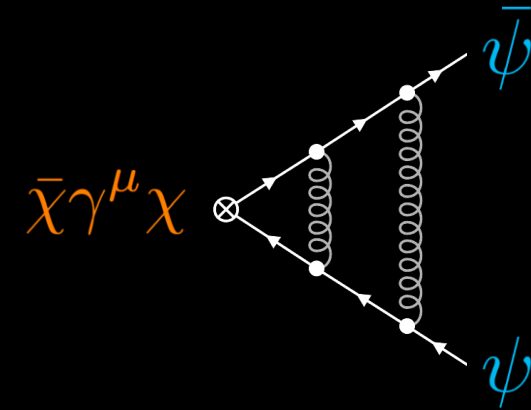


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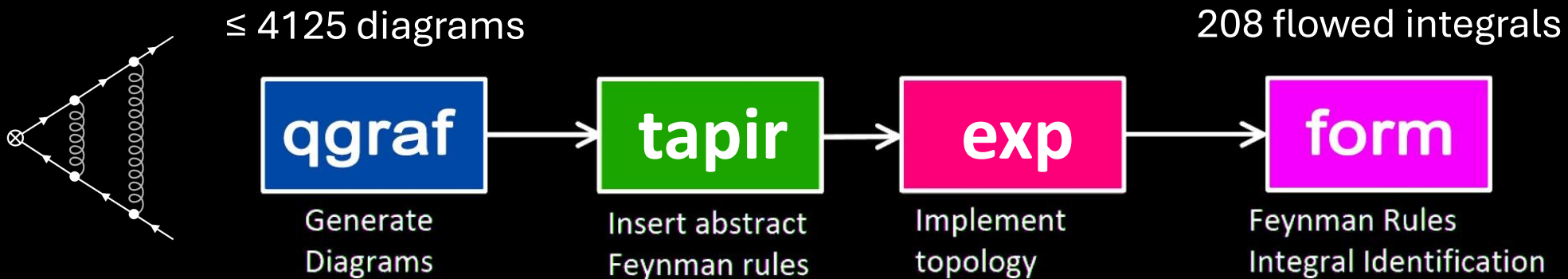
- ❖ Vector currents $\bar{\psi}\gamma^\mu\psi$ finite
see also: [JB, Harlander, Kohnen, Lange 2023]
- ❖ No operator renormalization
- ❖ Get Z_χ from SFTX



$$Z_\chi P_{\bar{\psi}\gamma^\mu\psi}[\bar{\chi}\gamma^\mu\chi] \stackrel{!}{=} \text{fin.}$$

Applications

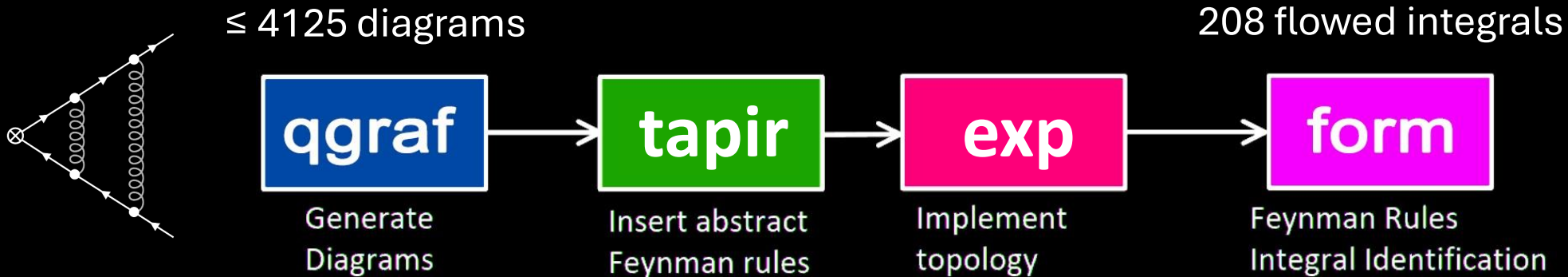
Renormalizing EFTs



Applications

Renormalizing EFTs

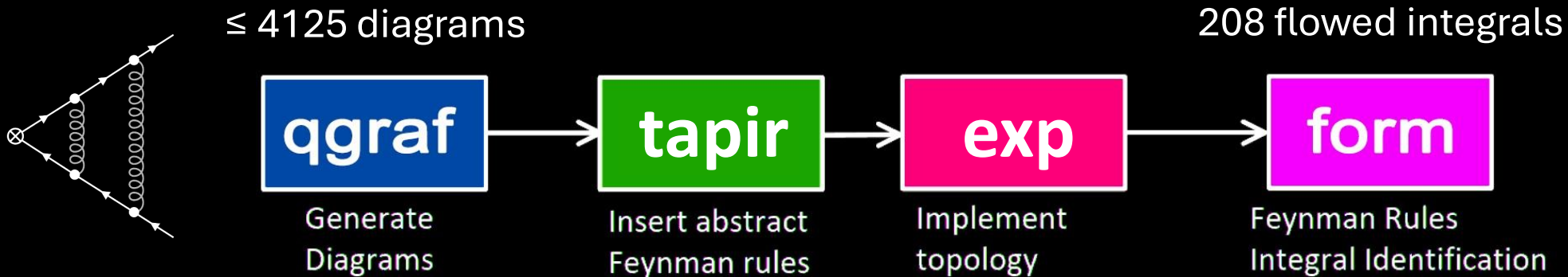
$$\int_{p,k} \int_{[0,1]^3} d^3u u^c \frac{e^{-[P_1(u)p^2 + P_2(u)k^2 + P_3(u)(p+k)^2]t}}{(p^2)^{a_1} (k^2)^{a_2} ((p+k)^2)^{a_3}}$$



Applications

Renormalizing EFTs

$$\int_{p,k} \int_{[0,1]^3} d^3u u^c \frac{e^{-[P_1(u)p^2 + P_2(u)k^2 + P_3(u)(p+k)^2]t}}{(p^2)^{a_1} (k^2)^{a_2} ((p+k)^2)^{a_3}}$$



Generate IBP relations

[Tkachov 1981; Chetyrkin, Tkachov 1981;
Artz, Harlander, Lange, Neumann, Prausa 2019]

Reduction with **kira+FireFly**

[Maierhöfer, Usovitsch, Uwer 2018; Klappert *et al.* 2021;
Klappert, Lange 2020; Klappert, Klein, Lange 2021]

→ 6 master integrals

Applications

Renormalizing EFTs

Results (unbroken phase, vanishing Yukawa couplings)

- ❖ Z_χ through NNLO for all fermions & Higgs
- ✓ Flowed currents independent of R_ξ gauge
- ✓ Agreement with SM & flowed QCD results
- ✓ Necessary ingredients for SMEFT renormalization

More Applications!

This talk:

- ❖ Energy-momentum tensor
- ❖ Strong coupling
- ❖ Renormalization of EFTs
- ❖ Neutron electric dipole moment
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- ❖ Meson mixing and lifetimes
- ❖ QCD static force

...

[Suzuki 2013; Harlander, Kluth, Lange 2018; ...]

[Lüscher 2010, 2014; Harlander, Neumann 2016; ...]

[*in preparation*]

[Mereghetti, Monahan, Rizik, Shindler, Stoffer 2022; ...]
[JB, Harlander, Rizik, Shindler 2022]

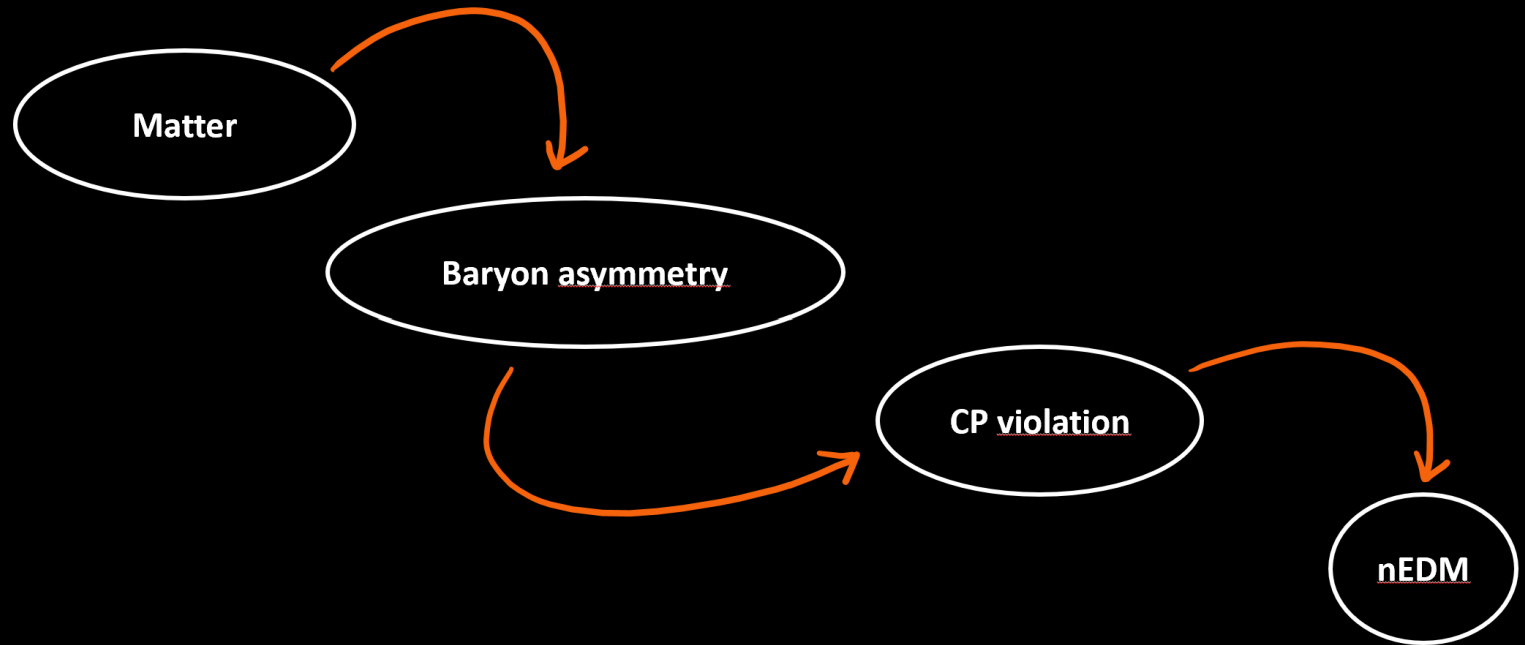
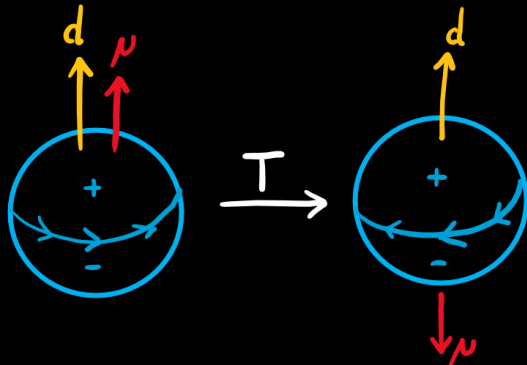
[Harlander, Lange, Neumann 2022]

[Black *et al.* 2024]

[Brambilla *et al.* 2022]

Appendix

nEDM



Experimental upper bound:

$$d_n = (0.0 \pm 1.1_{\text{stat}} \pm 0.2_{\text{sys}}) \times 10^{-26} e \text{ cm} \quad [\text{Abel et al. 2020}]$$

Future experiments estimate:

$$d_n \sim 10^{-27} e \text{ cm} \quad [\text{nEDM collaboration 2021, TUCAN collaboration 2022}]$$

Standard model (CKM):

$$d_n \sim 10^{-32} e \text{ cm} \quad [\text{Seng 2015, Gavela et al. 1994}]$$

Appendix

nEDM

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SU}(N)} + \sum_{n \geq 5} \frac{C_n}{\Lambda^n} \mathcal{O}^{(n)}$$

$$\mathcal{O}_{\text{CE}} = \bar{\psi} \sigma_{\mu\nu} \gamma_5 G^{\mu\nu} \psi$$

↓ γ_5 ⚡ DimReg

$$\mathcal{O}_{\text{CM}} = \bar{\psi} \sigma_{\mu\nu} G^{\mu\nu} \psi$$

Let's start with

Appendix

nEDM

Flowed operator: $\tilde{\mathcal{O}}_{\text{CM}}(t, x) \equiv \mathcal{O}_{\text{CM}}(x) \Big|_{A_\mu \rightarrow B_\mu(t) \ \psi \rightarrow \chi(t)}$

Recall SFTX: $\tilde{\mathcal{O}}_{\text{CM}}(t, x) = C_{\text{CM}}(t)\mathcal{O}_{\text{CM}}(x) + C_{\text{S}}(t)\mathcal{O}_{\text{S}}(x) + \dots$

Results:

$$\begin{aligned} C_{\text{CM}}(t) &= 1 + a_s(-4.023 + 0.166l_{\mu t}) \\ &\quad + a_s^2(-11.611 - 10.147l_{\mu t} + 0.229l_{\mu t}^2) + \mathcal{O}(a_s^3) \\ C_{\text{S}}(t)/i &= -2a_s + a_s^2(6.136 + 3.167) + \mathcal{O}(a_s^3) \end{aligned}$$

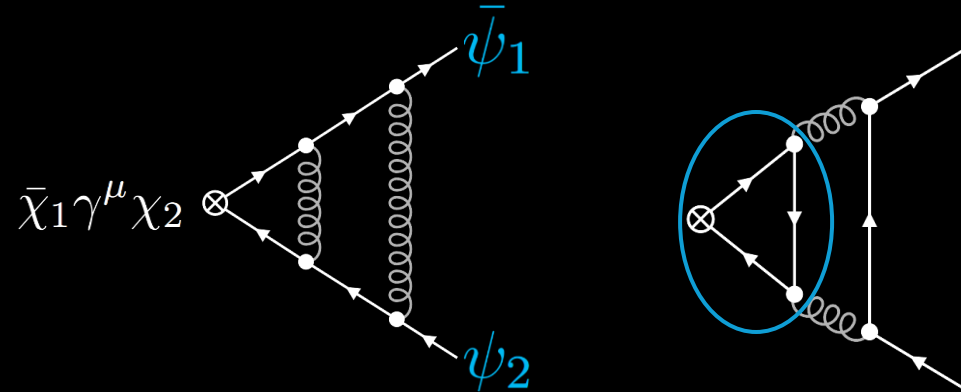
$$a_s = \frac{\alpha_s}{\pi}, \quad l_{\mu t} = \log(8\pi\mu^2 t)$$

[Mereghetti, Monahan, Rizik, Shindler, Stoffer 2022]
[JB, Harlander, Rizik, Shindler 2022]

Appendix

Flowed Standard Model

$$\tilde{Z}_f P_{\tilde{f}_1 \gamma^\mu f_2} [\tilde{f}_1 \gamma^\mu \tilde{f}_2] = \zeta_f \stackrel{!}{=} \text{fin.}$$



Definition of flowed fields

$$\tilde{f}(t, x)|_{t=0} = f(x)$$

$$\tilde{\phi}(t, x)|_{t=0} = \phi(x)$$

$$\mathcal{G}_\mu^a(t, x)|_{t=0} = G_\mu^a(x)$$

$$\mathcal{W}_\mu^a(t, x)|_{t=0} = W_\mu^a(x)$$

$$\mathcal{B}_\mu(t, x)|_{t=0} = B_\mu(x)$$

Flow equations

$$\partial_t \tilde{f} = \left[\mathcal{D}^2 - \kappa (\partial_\mu \mathcal{D}_\mu - \partial^2) \right] \tilde{f},$$

$$\partial_t \mathcal{G}_\mu^a = \mathcal{D}_\nu^{ab} \mathcal{G}_{\nu\mu}^b - \kappa \mathcal{D}_\mu^{ab} \partial_\nu \mathcal{G}_\nu^b,$$

$$\partial_t \mathcal{W}_\mu^a = \mathcal{D}_\nu^{ab} \mathcal{W}_{\nu\mu}^b - \kappa \mathcal{D}_\mu^{ab} \partial_\nu \mathcal{W}_\nu^b,$$

$$\begin{aligned} \partial_t \mathcal{B}_\mu &= \partial_\nu \mathcal{B}^{\nu\mu} + \kappa \partial_\mu \partial_\nu \mathcal{B}_\nu \\ &= \left[\partial^2 \delta_{\mu\nu} + (\kappa - 1) \partial_\mu \partial_\nu \right] \mathcal{B}_\nu \end{aligned}$$

→ Linear

Appendix

QCD flow equations

[Lüscher 2010; Lüscher 2013; Narayanan, Neuberger 2006]
 [Lüscher, Weisz 2011]

Flowed QCD: $\mathcal{L}_\chi = \int_0^\infty dt \bar{\lambda} (\partial_t \chi - \mathcal{D}_\mu \mathcal{D}_\mu \chi + \kappa \partial_\mu \mathcal{G}_\mu \chi)$

Flowed gauge-fixing

Lagrange multiplier field
 EOM = Flow Equation

Flowed ghost

$$\mathcal{L}_{\text{fQCD}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}_\chi + \mathcal{L}_B + \mathcal{L}_{\text{fl-gh}}$$

Flowed gauge-fixing

Flowed gluon: $\int_0^\infty dt L_\mu^a (\partial_t \mathcal{G}_\mu^a - \mathcal{D}_\nu^{ab} \mathcal{G}_{\nu\mu}^b + \alpha_0 \mathcal{D}_\mu^{ab} \partial_\nu \mathcal{G}_\nu^b)$, $\mathcal{G}_\mu^a(t, x)|_{t=0} = G_\mu^a(x)$

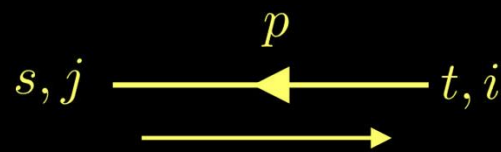
Appendix

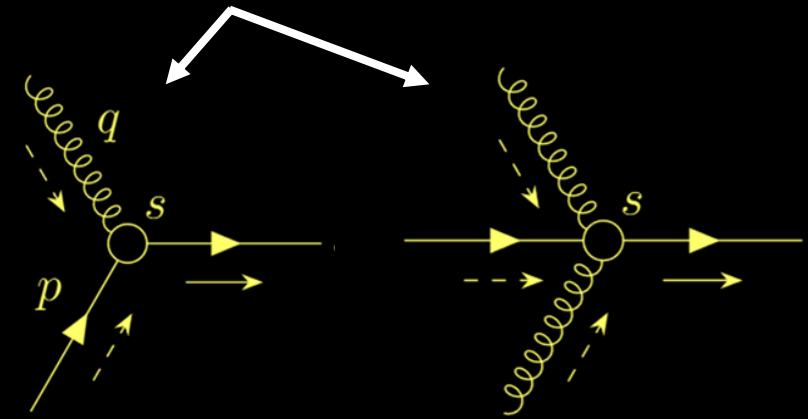
Feynman rules

[Lüscher 2010; Lüscher 2013; Narayanan, Neuberger 2006]
 [Lüscher, Weisz 2011]

Flowed QCD: $\mathcal{L}_\chi = \int_0^\infty dt \bar{\lambda} (\partial_t \chi - \mathcal{D}_\mu \mathcal{D}_\mu \chi + \kappa \partial_\mu \mathcal{G}_\mu \chi)$

$s, j \xrightarrow{p} t, i = \delta_{ij} \theta(t-s) e^{-(t-s)p^2}$





$\langle \bar{\chi}(t) \chi(s) \rangle = e^{-(t+s)p^2} \langle \bar{\psi} \psi \rangle \longrightarrow s, j \xrightarrow{p} t, i = \delta_{ij} \frac{-i\not{p} + m_0}{p^2 + m_0^2} e^{-(t+s)p^2}$

