

Flavour Physics in 2019

Fundamentals, Challenges, and Opportunities

Danny van Dyk

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Technische Universität München

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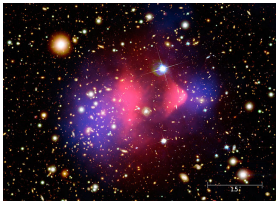
Fundamentals

Flavour at the Heart of Particle Physics?

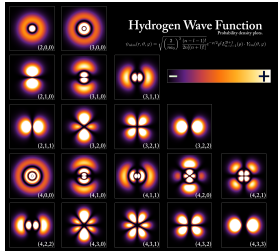
Our job as particle physicists is to determine:

- ▶ what are the elementary building blocks of the universe?
- ▶ what are the fundamental interactions between them?

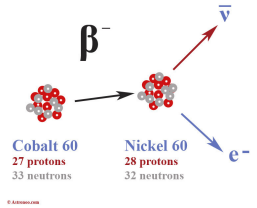
The goal is a uniform description of the building blocks and interactions, which can describe phenomena on very different length and time scales.



$10^{+22} m$



$10^{-10} m$



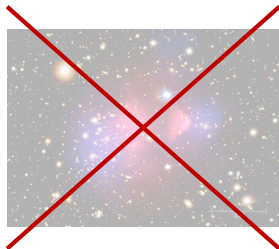
$10^{-19} m$

Flavour at the Heart of Particle Physics?

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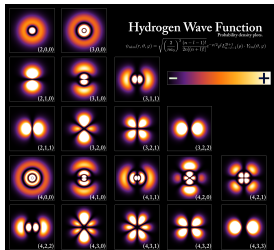
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- ▶ what are the fundamental interactions between them?

To date our best theory is the **Standard Model (SM)** of particle physics.



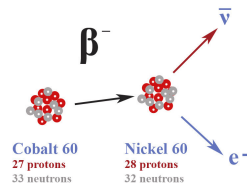
$10^{+22} m$

QED

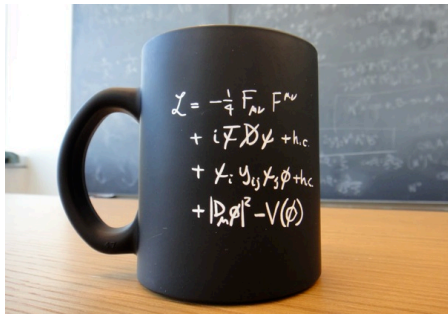


$10^{-10} m$

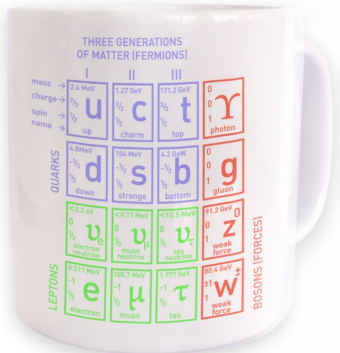
QCD,
weak interaction



$10^{-19} m$

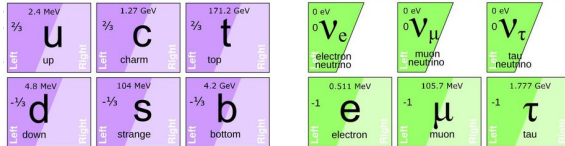


[CERN]



[somethinggeeky.com]

- ▶ matter fields are described in the SM in 3 generations
- ▶ the gauge interactions treat the generations universally
- ▶ the Higgs field **does not!**



$$-\mathcal{L} \supset H\bar{Q}Y_d d + H^c\bar{Q}Y_u u + H\bar{L}Y_e e + \text{h.c.}$$

symmetry between and within generations broken by the **Yukawa** terms

- ▶ we can observe quark masses as eigenvalues of $v/\sqrt{2}Y_{u,d}$
- ▶ we can observe misalignment between u -type and d -type quark fields through changes in the quark family (**horizontal jumps**)

The misalignment implies that “flavour” cannot be a conserved quantum number.

Quark Mixing Matrix

The quark Yukawa matrices can be simultaneously diagonalized through bi-unitary transformations

$$\text{diag}(m_u, m_c, m_t) = v/\sqrt{2} L_u Y_u R_u^\dagger$$

$$\text{diag}(m_d, m_s, m_b) = v/\sqrt{2} L_d Y_d R_d^\dagger$$

L_u, L_d unitary matrices applied to the left-handed fields

R_u, R_d unitary matrices applied to the right-handed fields

Generally: misalignment between L_u and L_d

- ▶ mass terms can be diagonalized simultaneously
- ▶ charged-current terms will be non-diagonal, misaligned by

$$(V_{\text{CKM}})_{ij} = (L_u \cdot L_d^\dagger)_{ij}$$

the **Cabbibo-Kobayashi-Maskawa (CKM) matrix**.

the CKM matrix

- ▶ is a 3×3 complex-valued unitary matrix
- ▶ is unitary by construction (since $L_{u,d}$ are unitary)

as a quark mixing matrix

- ▶ can be redefined such that 5 relative phases are absorbed into the quark field definitions
- ▶ has one one global (unobservable) phase

quark mixing in the SM

- ▶ can be therefore described in terms of only 4 parameters: 3 angles and 1 complex phase

[Kobayashi, Maskawa 1973]

$$\begin{pmatrix}
 V_{ud} & V_{us} & V_{ub} \\
 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\
 V_{cd} & V_{cs} & V_{cb} \\
 -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\
 V_{td} & V_{ts} & V_{tb} \\
 A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1
 \end{pmatrix}$$

expansion in $\lambda \simeq 0.2$ up to $\mathcal{O}(\lambda^3)$

How much information do we require to describe the SM?

gauge sector

- ▶ QED gauge coupling $\alpha_e(m_\mu)$,
- ▶ weak mixing angle via M_W/M_Z
- ▶ QCD gauge coupling $\alpha_s(m_Z)$ 3

Higgs sector

- ▶ Higgs vacuum expectation value v
- ▶ Higgs quartic coupling λ_4 2

flavour sector

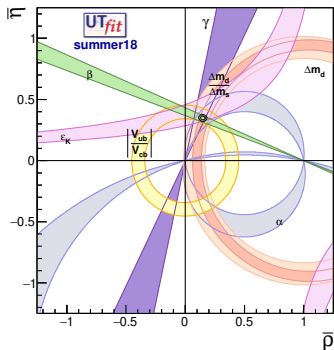
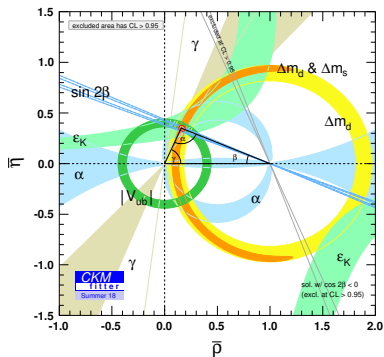
- ▶ charged lepton masses m_ℓ 3
- ▶ quark masses m_q 6
- ▶ CKM parameters λ, A, ρ, η 4

13 / 18 parameters are describing flavour in the SM

CKM Metrology

only four CKM parameters enter all flavour changing processes

- ▶ overconstrain the parameters in a **global fit** and check:
 - ▶ is the CKM matrix unitary?
 - ▶ is the CKM matrix the only source of CP violation?



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obtain constraints from ...

- ▶ tree-level (semi)leptonic decays
- ▶ loop-level neutral-current decays
- ▶ tree-level hadronic decays
- ▶ loop-level meson-mixing

$$c \rightarrow sW^* (\rightarrow \ell^+ \nu)$$

$$b \rightarrow s \{ \ell^+ \ell^-, \gamma \}$$

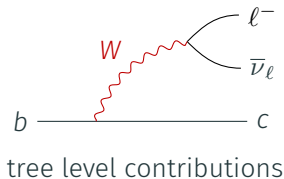
$$b \rightarrow cW^* (\rightarrow s\bar{u})$$

$$(\bar{b}d) \leftrightarrow (\bar{d}b)$$

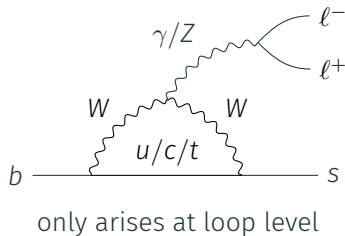
Semileptonic b Decays

2.4 MeV $\frac{2}{3}$ u Left up Right	1.27 GeV $\frac{2}{3}$ c Left charm Right	171.2 GeV $\frac{2}{3}$ t Left top Right
4.8 MeV $-\frac{1}{3}$ d Left down Right	104 MeV $-\frac{1}{3}$ s Left strange Right	4.2 GeV $-\frac{1}{3}$ b Left bottom Right

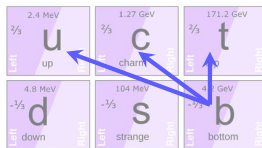
► w/ change of el. charge



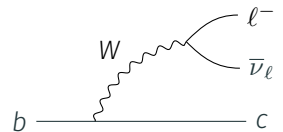
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Semileptonic b Decays

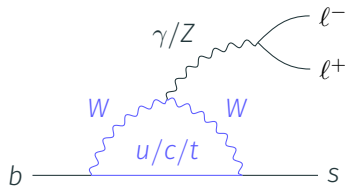


► w/ change of el. charge



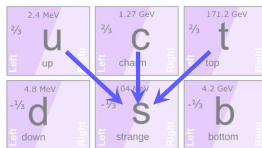
tree level contributions

► w/o change of el. charge

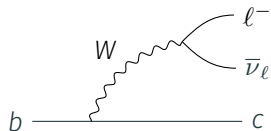


only arises at loop level

Semileptonic b Decays

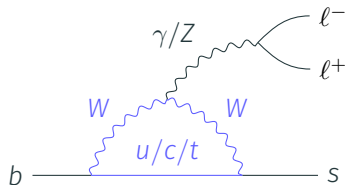


► w/ change of el. charge



tree level contributions

► w/o change of el. charge

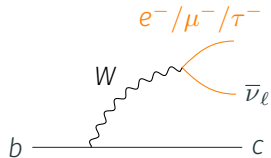


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Semileptonic b Decays

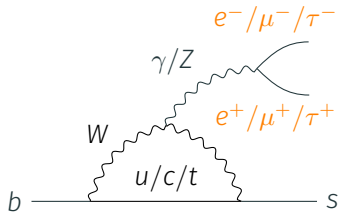
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► w/ change of el. charge



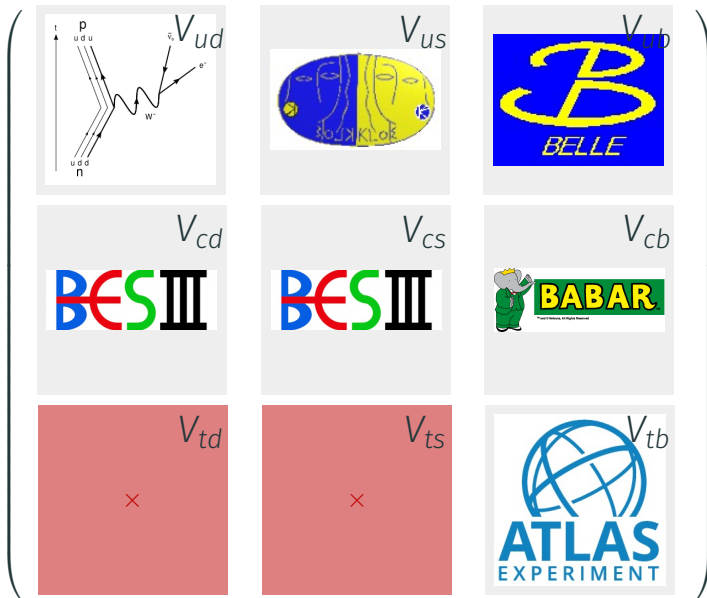
tree level contributions

► w/o change of el. charge



only arises at loop level

in both cases: **lepton-flavour-universal gauge couplings!**



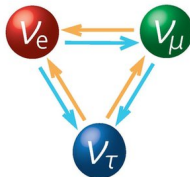
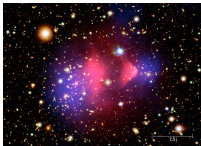
marked with \times : there are no direct measurements of s.l. decays available

So what is Flavour Physics then?

Flavour physics is ...

- ▶ ...determining the (flavour-related) majority of SM parameters with the best possible precision
- ▶ ...testing the assumptions inherent to the SM, e.g.:
 - ▶ the CKM matrix is the only source of flavour change
 - ▶ the CKM matrix is the only source of CP violation
 - ▶ there are exactly three generations of matter fields
- ▶ ...probing indirectly for New Physics by quantifying the deviations between SM predictions and measurements

since we know New Physics must hide somewhere



Challenges

Why “Anomalies”?

Theory predictions of B -meson decays have been and continue to be tested through several experiments

BaBar at the Stanford Linear Accelerator Center, USA

e^+e^- collisions

stopped in 2008

Belle (II) at KEK in Tsukuba, Japan

e^+e^- collisions

Belle data taking stopped in 2010

Belle II data taking (really) from 2019

ATLAS,CMS,LHCb at the Large Hadron Collider, CERN

pp collisions

LHC “run 2” ended in 2018

LHC “run 3” planned for 2021

⋮ ⋮ ⋮

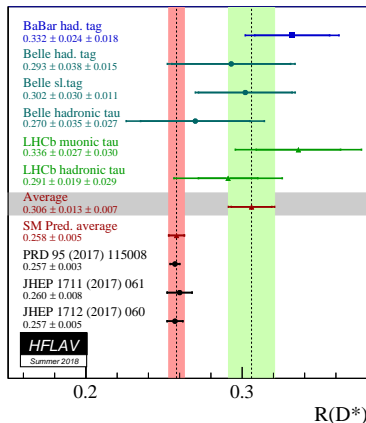
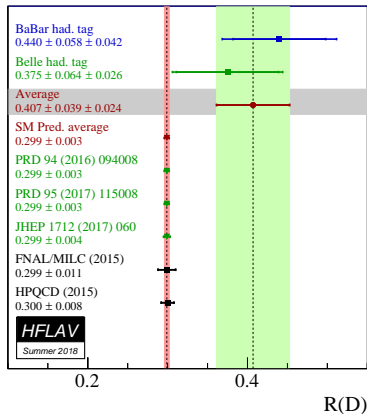
LHC “run 5” planned until 2035

Anomalies are those measurements that deviate from the SM theory predictions by more than 2 but less than 5 standard deviations.

Tests of **Lepton-Flavour-Universality** in $b \rightarrow c l \bar{\nu}$

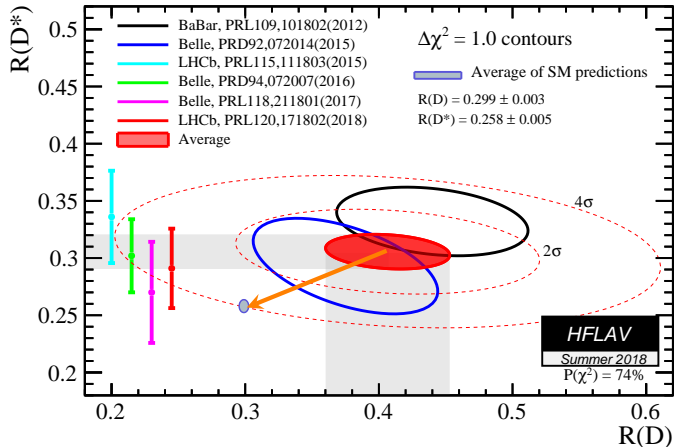
[HFLAV Sommer 2018]

$$\frac{b \rightarrow c \tau \bar{\nu}}{b \rightarrow c \mu \bar{\nu}} \rightarrow R(X) \equiv \frac{\Gamma(\bar{B} \rightarrow X \tau \bar{\nu})}{\Gamma(\bar{B} \rightarrow X \mu \bar{\nu})} \quad X = D, D^*$$



Tests of Lepton-Flavour-Universality in $b \rightarrow c\ell\bar{\nu}$

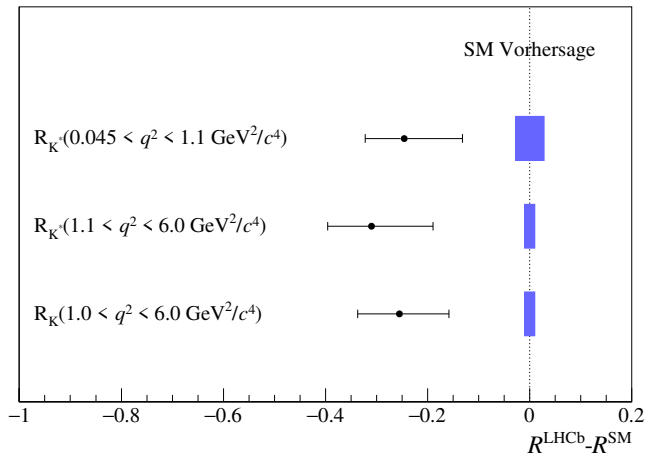
[HFLAV Summer 2018]



► combined significance: 3.62σ or $1 : 6370$

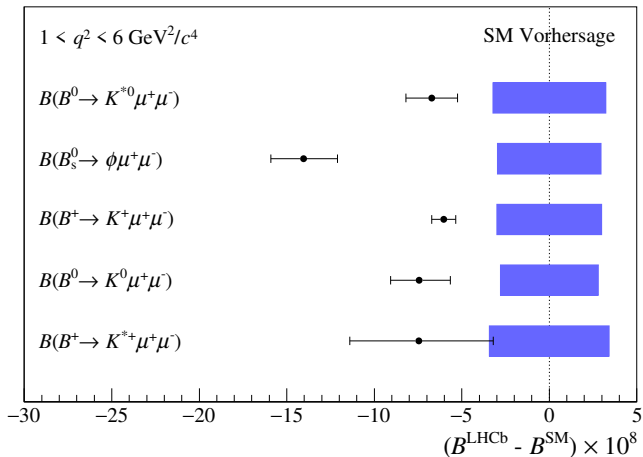
Tests of **Lepton-Flavour-Universalität** ($R(X)$), and decay rates, and angular distributions (P'_5) in $b \rightarrow sl^+\ell^-$

[Albrecht, Langenbruch Physik-Journal 2018]



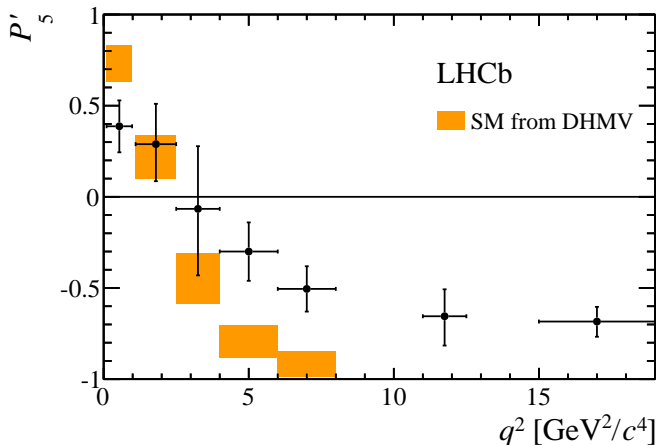
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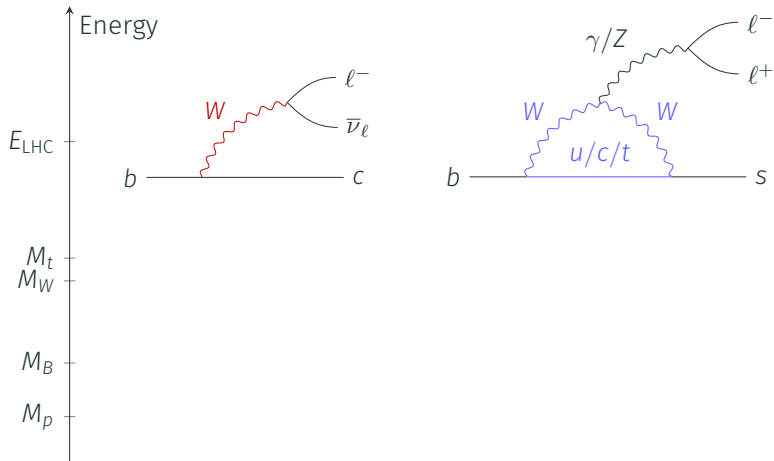
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► combined significance? → need weak effective theory (WET)

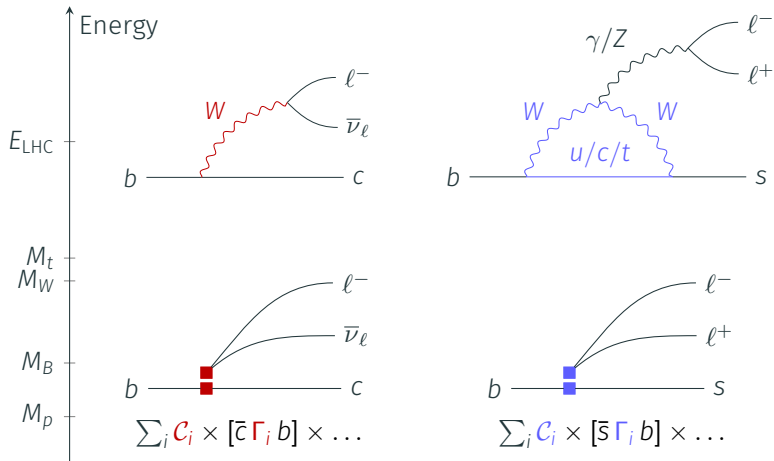
Effective (Quantum) Field Theories

- ▶ widely used tool of theoretical physics



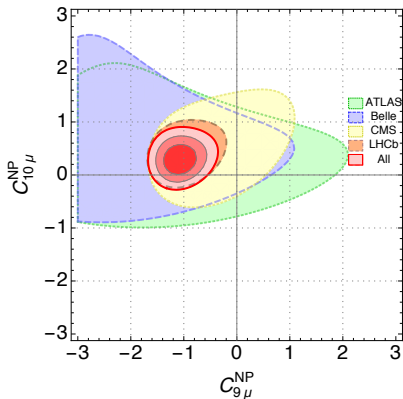
Effective (Quantum) Field Theories

- ▶ widely used tool of theoretical physics
- ▶ replaces dynamical degrees of freedom (here: t, W, Z) by coefficients \mathcal{C}_i and static structures in local operators (here: Γ_i)



- ▶ coefficients \mathcal{C} (over-)constrained through data
- ▶ w/o details: 2 coefficients $\mathcal{C}_{9,10}$ numerically dominant
- ▶ strong tension between best-fit point and SM prediction

[e.g. Descotes-Genon, Hofer, Matias, Virto 2015]



$$\mathcal{C}_{9,10}^{\text{NP}} \equiv \mathcal{C}_{9,10} - \mathcal{C}_{9,10}^{\text{SM}}$$

- ▶ consistent picture across all observ.:
 - ▶ decay rates
 - ▶ angular distributions
 - ▶ LFU ratios

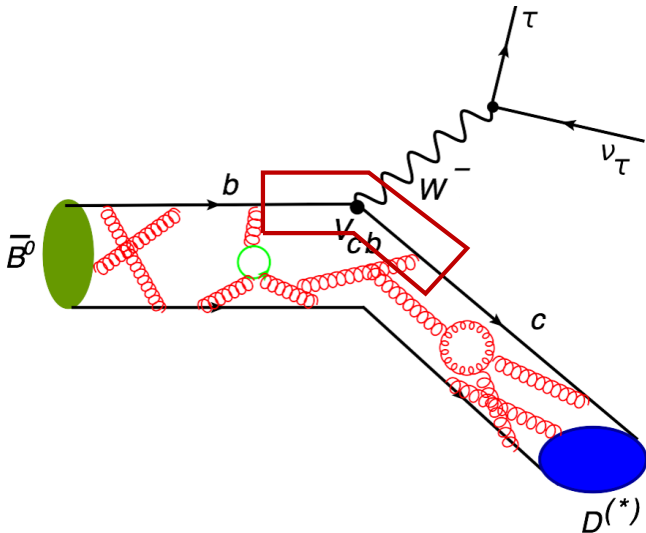
- ▶ significance $\sim 5\sigma$ *****
 - ▶ no single 5σ measurement!
 - ** form factors
 - *** non-local matrix elements

- ▶ tasks

E 5σ measurement! (z.B.: R_K)

T remove ***** caveats!

Matrix elements of local operators $\bar{c} \Gamma b$ (and $\bar{s} \Gamma b$) parametrised through **form factors**

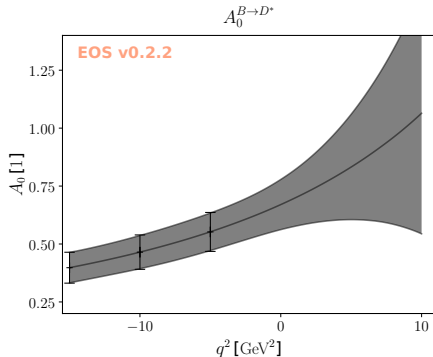
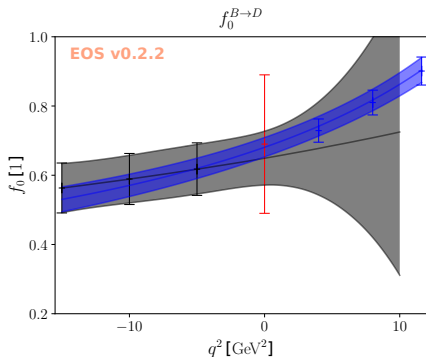


Matrix elements of local operators $\bar{c} \Gamma b$ (and $\bar{s} \Gamma b$) parametrised through **form factors**

- ▶ functions of momentum transfer (typical notation: q^2)
- ▶ 3 independent functions in e.g. $B \rightarrow D$ or $B \rightarrow K$
- ▶ 7 independent functions in e.g. $B \rightarrow D^*$ or $B \rightarrow K^*$
- ▶ low-energy QCD effects complicate direct calculation; requires numerical simulation (lattice QCD)

alternative: determination through light-cone sum rules with B -meson wave functions

- ▶ relatively recent method [Khodjamirian et al 2005,2006,2008]
[Feldmann et al. 2008]
- ▶ complementary to lattice QCD results
- ▶ recent simultaneous analysis of all form factors in $B \rightarrow P$ and $B \rightarrow V$ transitions, expanding on previous works [Gubernari,Kokulu,DvD 2018]



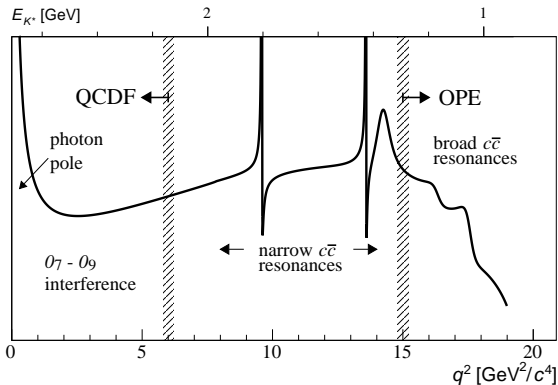
[Gubernari, Kokulu, DvD 2018]

- ▶ allows R_D and R_{D^*} predictions independent of $\bar{B} \rightarrow D^{(*)} \mu \bar{\nu}$ data
- ▶ compatible with $B \rightarrow D$ Lattice predictions
- ▶ must still be tested in framework of “heavy quark expansion”

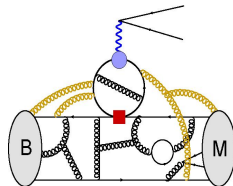
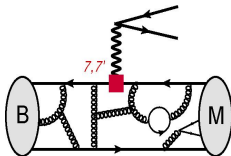
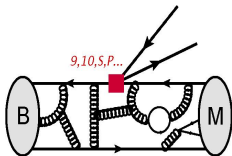
[Bordone, Jung, DvD w.i.p.]

- ▶ of relevance indep. of anomalies, since permits precision determinations of SM parameters (i.e.: $|V_{cb}|$, $|V_{ub}|$)

$B \rightarrow K^* \mu^+ \mu^-$ landscape:



[sketch from Blake, Gershon, Hiller 2015]



$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ (C_9 \mp C_{10}) \mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[C_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

▶ non-local: $\mathcal{H}_\lambda(q^2) = i\mathcal{P}_\mu^\lambda \int d^4x e^{iq \cdot x} \langle \bar{M}_\lambda(k) | T \{ \mathcal{J}_{em}^\mu(x), C_i \mathcal{O}_i(0) \} | \bar{B}(q+k) \rangle$

below J/ψ : light-cone OPE

[Khodjamirian, Mannel, Pivovarov, Wang 2010]

- ▶ leading (local) terms as in QCD factorisation
- ▶ next-to-leading terms from light-cone sum rules

[Beneke, Feldmann, Seidel 2001&2004]

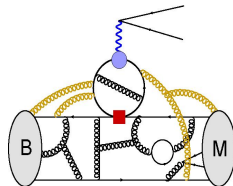
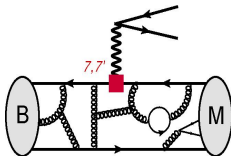
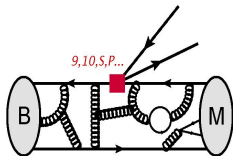
[Khodjamirian, Mannel, Pivovarov, Wang 2010]

above the $\psi(2S)$: local OPE in $1/m_b$ and $1/\sqrt{q^2}$; $q^2 = m^2(\mu\mu)$

- ▶ OPE in HQET operators
- ▶ OPE in QCD operators

[Grinstein, Pirjol 2004]

[Beylich, Buchalla, Feldmann 2011]



$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ (C_9 \mp C_{10}) \mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[C_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

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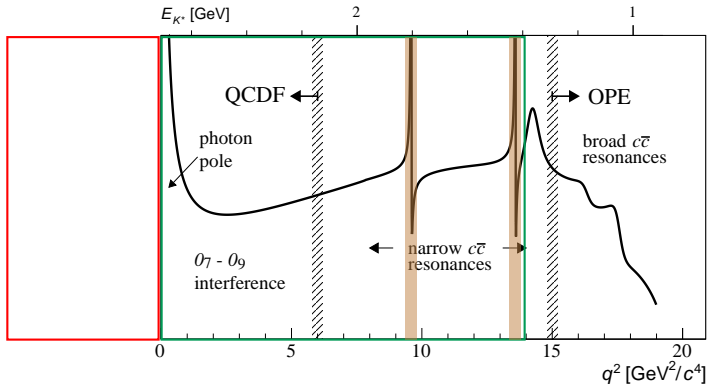
▶ OPE in HQET operators

[Grinstein, Pirjol 2004]

▶ OPE in QCD operators

[Beylich, Buchalla, Feldmann 2011]

Strategy



[sketch from Blake, Gershon, Hiller 2015]

- ▶ **calculate** non-local matrix elements at $q^2 < 0$
- ▶ **extrapolate** to $q^2 > 0$ via some type of analytic continuation
- ▶ **constrain** two narrow resonances at $q^2 > 0$ from data on $B \rightarrow \psi_n K^*$

Non-local Effects in $B \rightarrow K^* \mu^+ \mu^-$

- ▶ new idea to parametrize non-local effects based on analyticity properties

[Bobeth, Chrzaszcz, van Dyk, Virto 2017]

- ▶ includes theory predictions and experimental data **simultaneously**

some details for actual parametrisation

- ▶ parametrize the ratios $\mathcal{H}_\lambda(q^2)/\mathcal{F}_\lambda(q^2)$
- ▶ poles for subthreshold resonances $J/\psi, \psi(2S)$
- ▶ rest: power series in parameter z , which has correct analyticity properties
- ▶ the poles should not modify the asymptotic behaviour at $|q^2| \rightarrow \infty$

$$\mathcal{H}_\lambda(z) = \frac{1 - z z_{J/\psi}^*}{z - z_{J/\psi}} \frac{1 - z z_{\psi(2S)}^*}{z - z_{\psi(2S)}} \hat{\mathcal{H}}_\lambda(z)$$

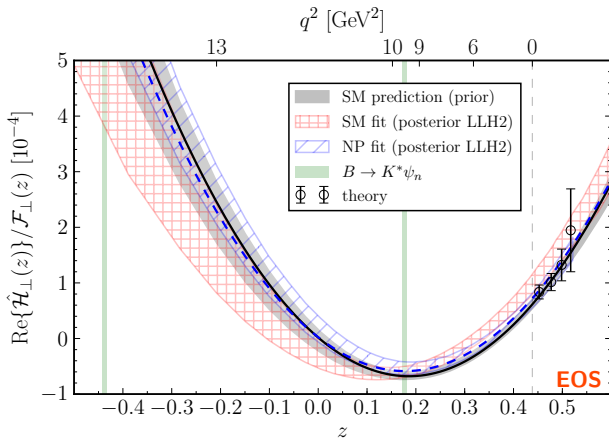
$$\hat{\mathcal{H}}_\lambda(z) = \left[\sum_{k=0}^K \alpha_k^{(\lambda)} z^k \right] \mathcal{F}_\lambda(z)$$

Non-local Effects in $B \rightarrow K^* \mu^+ \mu^-$

- ▶ new idea to parametrize non-local effects based on analyticity properties

[Bobeth, Chrzaszcz, van Dyk, Virto 2017]

- ▶ includes theory predictions and experimental data **simultaneously**



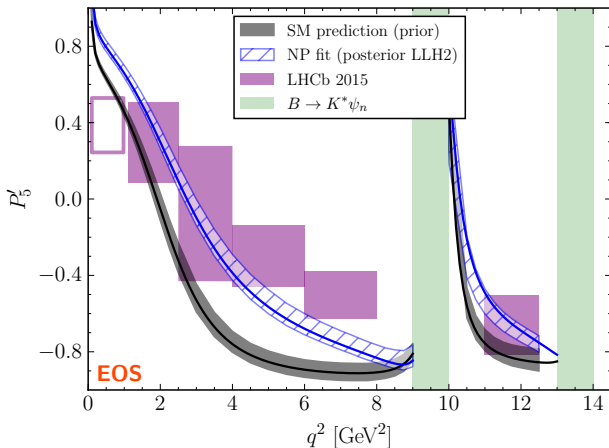
- ▶ parametrisation does not provide enough freedom in the SM fit in order to deviate substantially from the prior

Non-local Effects in $B \rightarrow K^* \mu^+ \mu^-$

- ▶ new idea to parametrize non-local effects based on analyticity properties

[Bobeth, Chrzaszcz, van Dyk, Virto 2017]

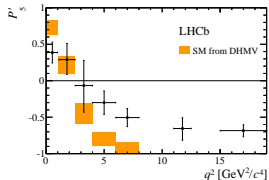
- ▶ includes theory predictions and experimental data **simultaneously**



- ▶ parametrisation does not provide enough freedom in the SM fit in order to deviate substantially from the prior

So ...?

So is this New Physics then?



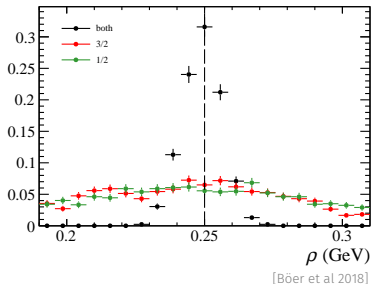
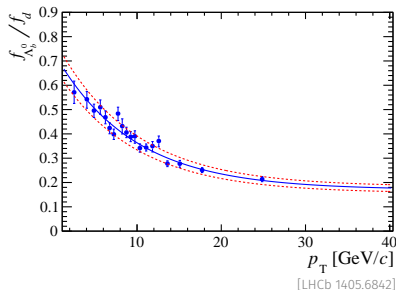
Jury's still out!

- ▶ no single 5σ deviation
- ▶ so far no theorist confident enough about old physics to claim a New Physics discovery!
- ▶ challenging situation indeed!

Opportunities

Opportunity: The LHC as a Λ_b^0 Factory

(bdu) with $J^P = 1/2^+$ and $I = 0$



- ▶ no anomalies (yet?) with b baryons
- ▶ independent “laboratory” to check theory
- ▶ QCD factorisation harder to prove than for B mesons
- ▶ heavy-quark expansion seems to work much better than for mesons

[e.g. Feldman, Yip 2011]

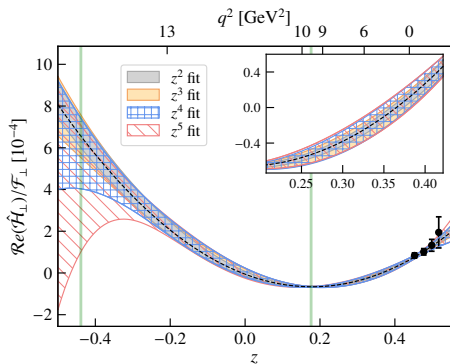
$[\Lambda_b \rightarrow \Lambda_c$: Bernlochner, Ligeti, Robinson, Sutcliffe 2018]

$[\Lambda_b \rightarrow \Lambda_c^*$: Böer, Bordone, Graverini, Owen, Rotondo, DvD 2018]

Opportunity: Future Combined Analyses (Theory + Experiment)

Sensitivity to non-local parameters in $B \rightarrow K^* \mu^+ \mu^-$ from an unbinned fit

[Chrzaszcz, Mauri, Serra, Silva Coutinho, van Dyk 2018]



- ▶ cover various benchmark points for the non-local matrix elements
 - ▶ includes approaches with polynomials up to order z^5
 - ▶ uses central theory inputs from pheno analysis

[Bobeth, Chrzaszcz, van Dyk, Virto 2017]

- ▶ find sensitivity to z^3 and higher coefficients
- ▶ constrain non-local matrix elements well for region $1.1 \text{ GeV}^2 \leq q^2 \leq 9 \text{ GeV}^2$

Opportunity: Belle 2 complementary to LHCb

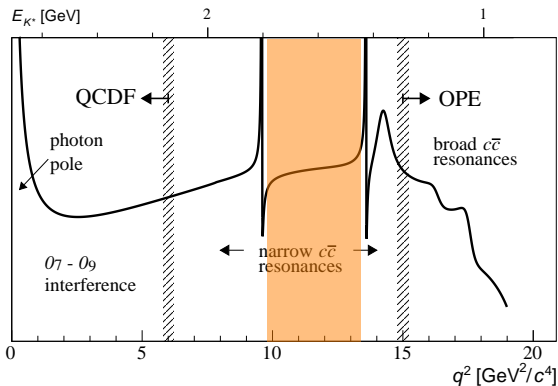
common lore:

[Belle 2 Physics Book]

- ▶ inclusive $B \rightarrow X_S \mu^+ \mu^-$ golden channel for Belle 2
- ▶ for exclusive $B \rightarrow K^* \mu^+ \mu^-$ LHCb **expected** to have better precision throughout

measurements in between and close to $J/\psi, \psi'$ resonances difficult for LHCb!

- ▶ opportunity for Belle 2 to contribute substantially



Opportunity: Shameless Self-advertisement



- ▶ **C++14** software framework for flavour observables
 - ▶ includes Python bindings
- ▶ use cases
 - ▶ theory predictions
 - ▶ Bayesian parameter inference
 - ▶ Monte-Carlo simulation of signal events
- ▶ open source / binary packages for Debian & Ubuntu

[EOS homepage](#)

[GitHub repository](#)

- ▶ available also through the Belle 2 Externals

Summary

Summary

- ▶ flavour physics is a corner stone of particle physics (2019 or not!)
- ▶ accurate and precise theory predictions within the SM and beyond are important to further our understanding of physics at shortest distances
 - ▶ essentiell for the interpretation of the present anomalies
 - ▶ even if anomalie are a (statistical) fluke: theory predictions govern precision in the determination of SM parameters
- ▶ flavour physics thrives through close collaboration between theory and experiment
 - ▶ LHC runs 3 through 5 and Bell 2 will pile up a huge treasure of data
 - ▶ will require concerted efforts to fully exploit the data!

Backup Slides
