

Impact of Vacuum Stability on the Phenomenology of SUSY Models

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and JHEP 08 (2016) 126 and PLB 752 (2016) 7-12

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The Standard Model (In)Stability

$$V_{\text{SM}} = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

- large field values: $V \sim \lambda (H^\dagger H)^2$
- RGE: $\lambda \rightarrow \lambda(Q)$, where $Q \sim H$
- $\lambda \rightarrow 0$ around $Q \sim 10^{10}$ GeV, new minimum beyond M_{Planck}

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The MSSM: less simple

$$V_{\text{MSSM}} = V_F + V_{\text{soft}} + V_D$$

with (only 3rd generation squarks and Higgses)

$$\begin{aligned} V_{\text{soft}} = & m_{H_d}^2 |h_d|^2 + m_{H_u}^2 |h_u|^2 - (B_\mu h_d \cdot h_u + \text{h. c.}) \\ & + \tilde{t}_L^* \tilde{m}_Q^2 \tilde{t}_L + \tilde{t}_R^* \tilde{m}_t^2 \tilde{t}_R + \tilde{b}_L^* \tilde{m}_Q^2 \tilde{b}_L + \tilde{b}_R^* \tilde{m}_b \tilde{b}_R \\ & + (A_t h_u \tilde{t}_L^* \tilde{t}_R + A_b h_d \tilde{b}_L^* \tilde{b}_R + \text{h. c.}) \end{aligned}$$

A multi-scalar theory

- 2 Higgs doublets
 - 2×6 scalar quarks, $6 + 3$ scalar leptons
 - 12 colored and $18 + 2$ charged directions
 - charged Higgs directions “safe” [Casas et al. 1996]
 - SM Higgs potential: $SO(4)$ symmetry
-
- large couplings to Higgs doublets (y_t and y_b comparably large)
 - large stop contribution (X_t, A_t) to light Higgs mass needed
 - SUSY threshold corrections for m_b influence y_b

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impossible!

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An analytic solution?

impossible!

only approximative

$$\begin{aligned}
 V_{\tilde{q},h} = & \tilde{t}_L^* (\tilde{m}_L^2 + |y_t h_2|^2) \tilde{t}_L + \tilde{t}_R^* (\tilde{m}_t^2 + |y_t h_2|^2) \tilde{t}_R \\
 & + \tilde{b}_L^* (\tilde{m}_L^2 + |y_b h_1|^2) \tilde{b}_L + \tilde{b}_R^* (\tilde{m}_b^2 + |y_b h_1|^2) \tilde{b}_R \\
 & - [\tilde{t}_L^* (\mu^* y_t h_1^* - A_t h_2) \tilde{t}_R + \text{h.c.}] \\
 & - [\tilde{b}_L^* (\mu^* y_b h_2^* - A_b h_1) \tilde{b}_R + \text{h.c.}] \\
 & + |y_t|^2 |\tilde{t}_L|^2 |\tilde{t}_R|^2 + |y_b|^2 |\tilde{b}_L|^2 |\tilde{b}_R|^2 \\
 & + \frac{g_1^2}{8} \left(|h_2|^2 - |h_1|^2 + \frac{1}{3} |\tilde{b}_L|^2 + \frac{2}{3} |\tilde{b}_R|^2 + \frac{1}{3} |\tilde{t}_L|^2 - \frac{4}{3} |\tilde{t}_R|^2 \right)^2 \\
 & + \frac{g_2^2}{8} (|h_2|^2 - |h_1|^2 + |\tilde{b}_L|^2 - |\tilde{t}_L|^2)^2 \\
 & + \frac{g_3^2}{8} (|\tilde{t}_L|^2 - |\tilde{t}_R|^2 + |\tilde{b}_L|^2 - |\tilde{b}_R|^2)^2 \\
 & + (m_{h_2}^2 + |\mu|^2) |h_2|^2 + (m_{h_1}^2 + |\mu|^2) |h_1|^2 - 2 \text{Re}(B_\mu h_1 h_2).
 \end{aligned}$$

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 V_{\tilde{q},h} = & \tilde{t}_L^* (\tilde{m}_L^2 + |y_t h_2|^2) \tilde{t}_L + \tilde{t}_R^* (\tilde{m}_t^2 + |y_t h_2|^2) \tilde{t}_R \\
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 & - [\tilde{t}_L^* (\mu^* y_t h_1^* - A_t h_2) \tilde{t}_R + \text{h.c.}] \\
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 & + |y_t|^2 |\tilde{t}_L|^2 |\tilde{t}_R|^2 + |y_b|^2 |\tilde{b}_L|^2 |\tilde{b}_R|^2 \\
 & + \frac{g_1^2}{8} \left(|h_2|^2 - |h_1|^2 + \frac{1}{3} |\tilde{b}_L|^2 + \frac{2}{3} |\tilde{b}_R|^2 + \frac{1}{3} |\tilde{t}_L|^2 - \frac{4}{3} |\tilde{t}_R|^2 \right)^2 \\
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$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, \quad |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|$$

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 V_{\tilde{q},h} = & \tilde{t}^* (\tilde{m}_L^2 + |y_t h_2|^2) \tilde{t} + \tilde{t}^* (\tilde{m}_t^2 + |y_t h_2|^2) \tilde{t} \\
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$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|; |\tilde{b}| = |h_1| = |\phi_1|, |\tilde{t}| = |h_2| = |\phi_2|$$

The tree-level scalar potential

$$\begin{aligned} V_{\tilde{q},h} = & \phi_2^* (\tilde{m}_L^2 + |y_t \phi_2|^2) \phi_2 + \phi_2^* (\tilde{m}_t^2 + |y_t \phi_2|^2) \phi_2 \\ & + \phi_1^* (\tilde{m}_L^2 + |y_b \phi_1|^2) \phi_1 + \phi_1^* (\tilde{m}_b^2 + |y_b \phi_1|^2) \phi_1 \\ & - [\phi_2^* (\mu^* y_t \phi_1^* - A_t \phi_2) \phi_2 + \text{h.c.}] \\ & - [\phi_1^* (\mu^* y_b \phi_2^* - A_b \phi_1) \phi_1 + \text{h.c.}] \\ & + |y_t|^2 |\phi_2|^2 |\phi_2|^2 + |y_b|^2 |\phi_1|^2 |\phi_1|^2 \\ & + (m_{h_2}^2 + |\mu|^2) |\phi_2|^2 + (m_{h_1}^2 + |\mu|^2) |\phi_1|^2 - 2 \text{Re}(B_\mu \phi_1 \phi_2). \end{aligned}$$

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$$- [\phi_2^* (-A_t \phi_2) \phi_2 + \text{h.c.}]$$

$$+ |y_t|^2 |\phi_2|^2 |\phi_2|^2$$

$$+ (m_{h_2}^2 + |\mu|^2) |\phi_2|^2$$

$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|; \quad \cancel{|\tilde{b}| = |h_1| = |\phi_1|}, |\tilde{t}| = |h_2| = |\phi_2|$$

Minimize the potential

$$V(\phi) = m^2 \phi^2 - A\phi^3 + \lambda\phi^4,$$

with $m^2 = m_{h_2}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_t^2$, $A = -A_t$ and $\lambda = 3y_t^2$.

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Answer:

$$\phi_0 = 0, \quad \phi_{\pm} = \frac{3A \pm \sqrt{9A^2 - 32\lambda m^2}}{8\lambda}.$$

Condition to be safe from non-standard (i.e. non-trivial) minima:

$$V(\phi_{\pm}) > 0 \quad \Leftrightarrow \quad m^2 > \frac{A^2}{4\lambda}$$

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Well-known constraints

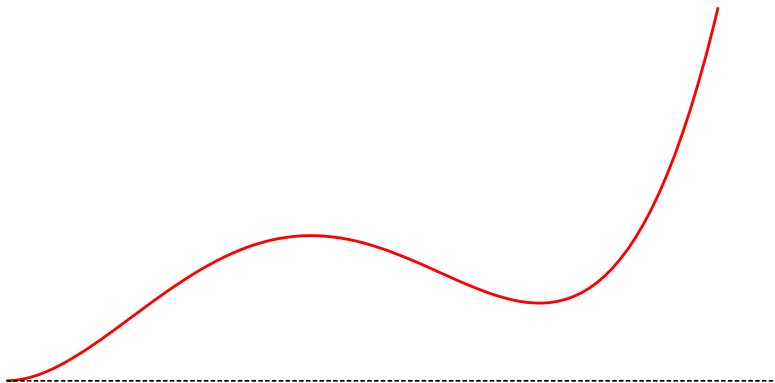
[Gunion, Haber, Sher '88]

$$|A_t|^2 < 3y_t^2 (m_{h_2}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_t^2)$$

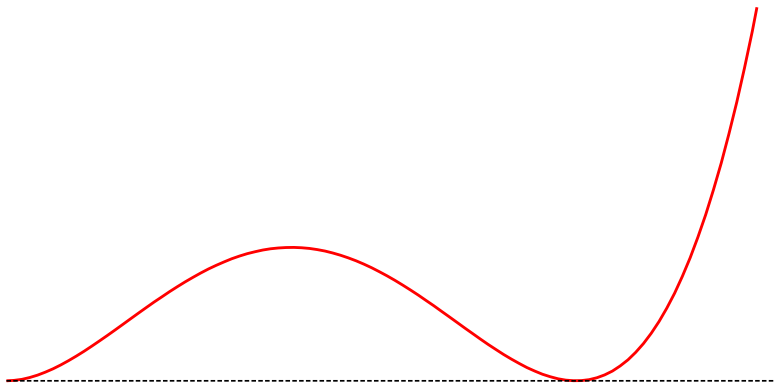
$$|A_b|^2 < 3y_b^2 (m_{h_1}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_b^2)$$

for the limiting cases $|\tilde{t}_L| = |\tilde{t}_R| = |h_2|$ and $|\tilde{b}_L| = |\tilde{b}_R| = |h_1|$!

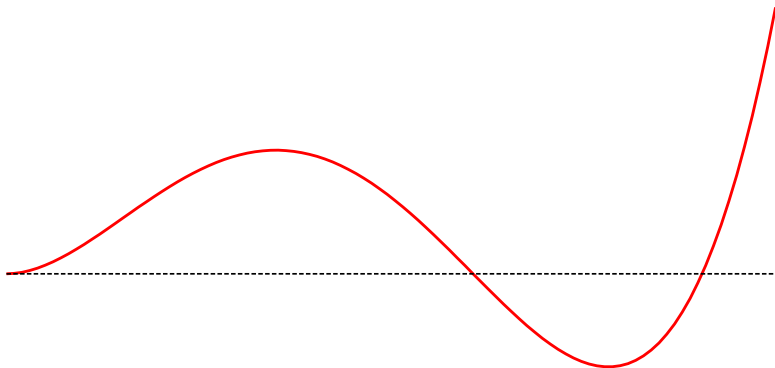
$$A^2 < 4\lambda m^2$$



$$A^2 = 4\lambda m^2$$



$$A^2 > 4\lambda m^2$$



A simple view of a complicated object

$$h_2 = \phi, \quad |\tilde{t}| = \alpha|\phi|, \quad h_1 = \eta\phi, \quad |\tilde{b}| = \beta|\phi|$$

$$\begin{aligned} V_\phi = & (m_{h_2}^2 + \eta^2 m_{h_1}^2 + (1 + \eta^2)\mu^2 - 2B_\mu\eta \\ & + (\alpha^2 + \beta^2)\tilde{m}_L^2 + \alpha^2\tilde{m}_t^2 + \beta^2\tilde{m}_b^2)\phi^2 \\ & - 2(\alpha^2(\mu y_t \eta - A_t) + \beta^2(\mu y_t - \eta A_b))\phi^3 + (\alpha^2 y_t^2 + \beta^4 y_b^2)\phi^4 \\ & + \left(\frac{g_1^2 + g_2^2}{8}(1 - \eta^2 + \beta^2 - \alpha^2)^2 + 2\alpha^2 y_t^2 + 2\beta^2 y_b^2 \right)\phi^4 \\ \equiv & M^2(\eta, \alpha, \beta)\phi^2 - \mathcal{A}(\eta, \alpha, \beta)\phi^3 + \lambda(\eta, \alpha, \beta)\phi^4, \end{aligned}$$

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with

$$\begin{aligned} M^2 = & m_{h_2}^2 + \eta^2 m_{h_1}^2 + (1 + \eta^2)\mu^2 - 2B_\mu\eta \\ & + (\alpha^2 + \beta^2)\tilde{m}_L^2 + \alpha^2\tilde{m}_t^2 + \beta^2\tilde{m}_b^2, \\ \mathcal{A} = & 2\alpha^2\eta\mu y_t - 2\alpha^2 A_t + 2\beta^2\mu y_b - 2\eta\beta^2 A_b, \\ \lambda = & \frac{g_1^2 + g_2^2}{8}(1 - \eta^2 + \beta^2 - \alpha^2)^2 \\ & + (2 + \alpha^2)\alpha^2 y_t^2 + (2\eta^2 + \beta^2)\beta^2 y_b^2. \end{aligned}$$

[Gunion, Haber, Sher '88; Casas, Lleyda, Muñoz '96]

The same but different ("A-parameter bounds")

$$A^2 < 4\lambda M^2$$

$$\downarrow$$

$$4 \min_{\{\eta, \alpha, \beta\}} \lambda(\eta, \alpha, \beta) M^2(\eta, \alpha, \beta) > \max_{\{\eta, \alpha, \beta\}} (A(\eta, \alpha, \beta))^2$$

$$h_u = \tilde{b}, h_d^0 = 0$$

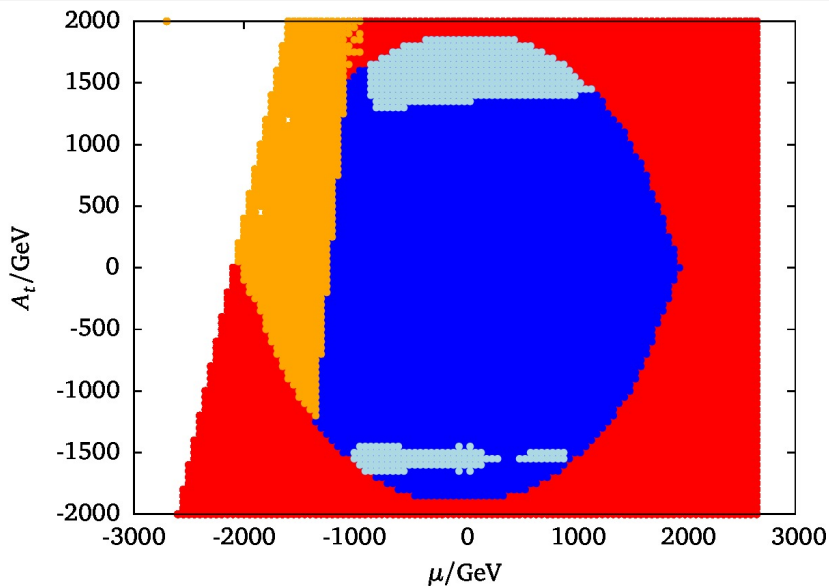
[WGH'15]

$$m_{H_u}^2 + \mu^2 + \tilde{m}_Q^2 + \tilde{m}_b^2 > \frac{(\mu y_b)^2}{y_b^2 + (g_1^2 + g_2^2)/2}$$

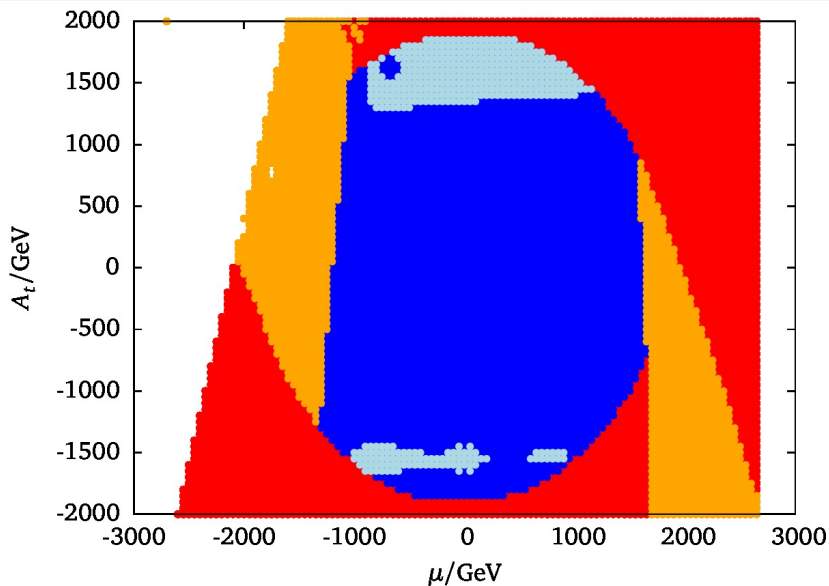
$$|h_d|^2 = |h_u|^2 + |\tilde{b}|^2, \tilde{b} = \alpha h_u$$

[WGH'15]

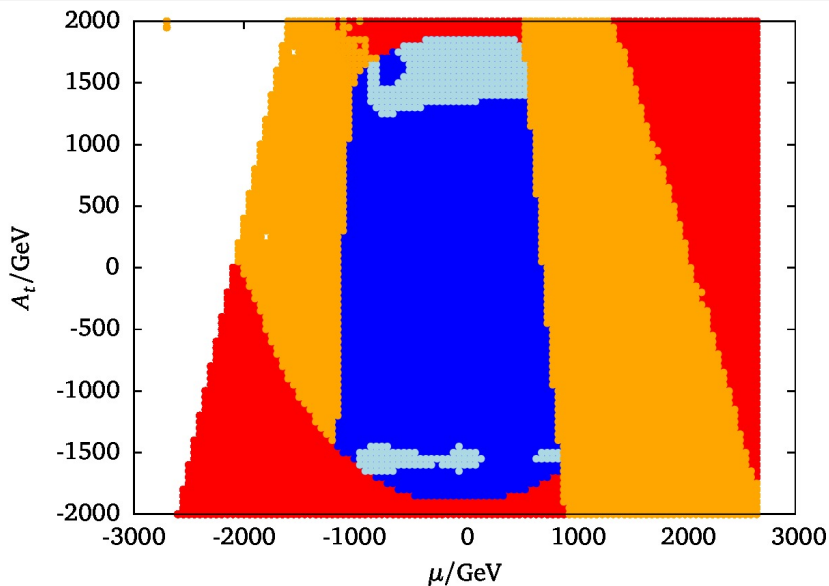
$$m_{11}^2(1 + \alpha^2) + m_{22}^2 \pm 2m_{12}^2 \sqrt{1 + \alpha^2} + \alpha^2(\tilde{m}_Q^2 + \tilde{m}_b^2) > \frac{4\mu^2 \alpha^2}{2 + 3\alpha^2}$$



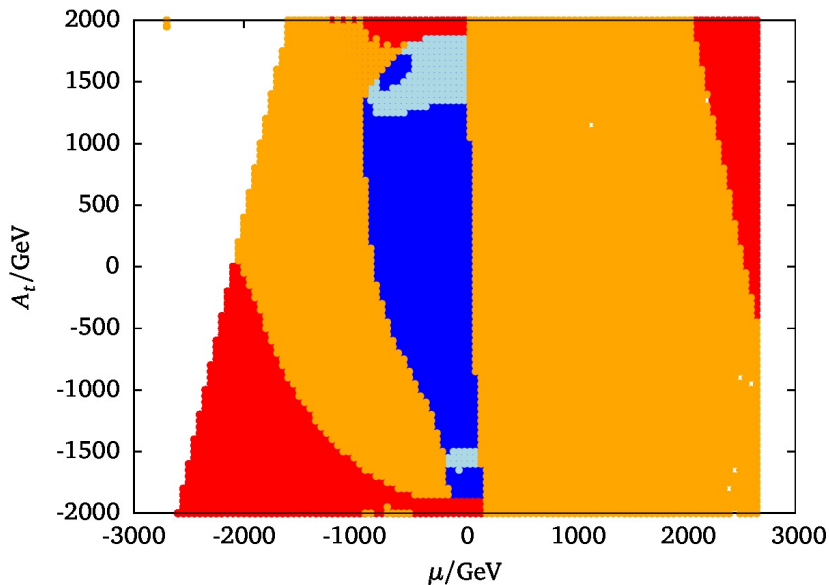
[WGH 2016]



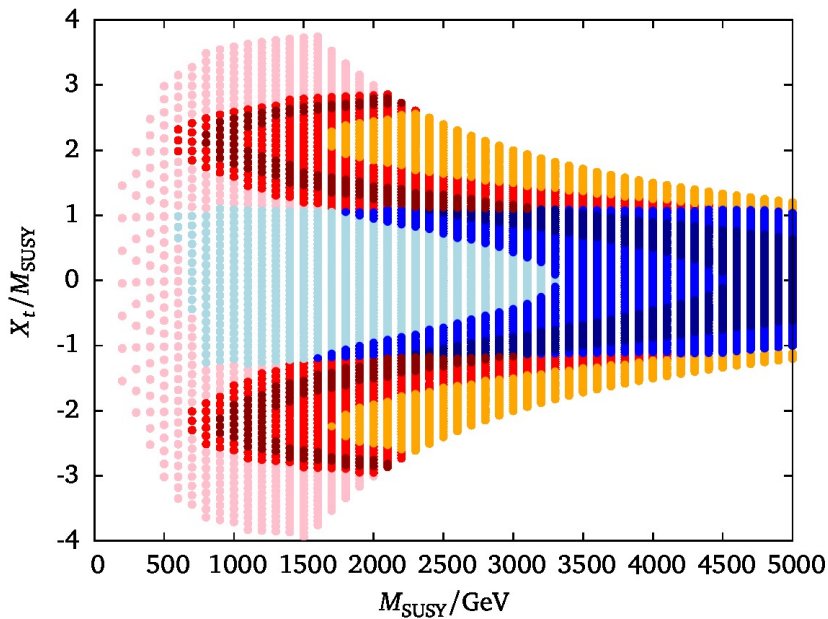
[WGH 2016]

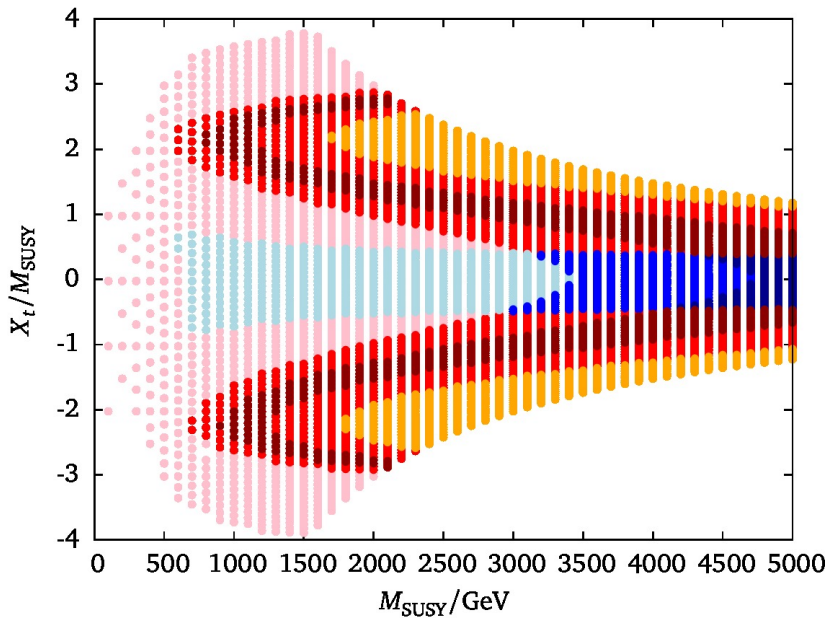


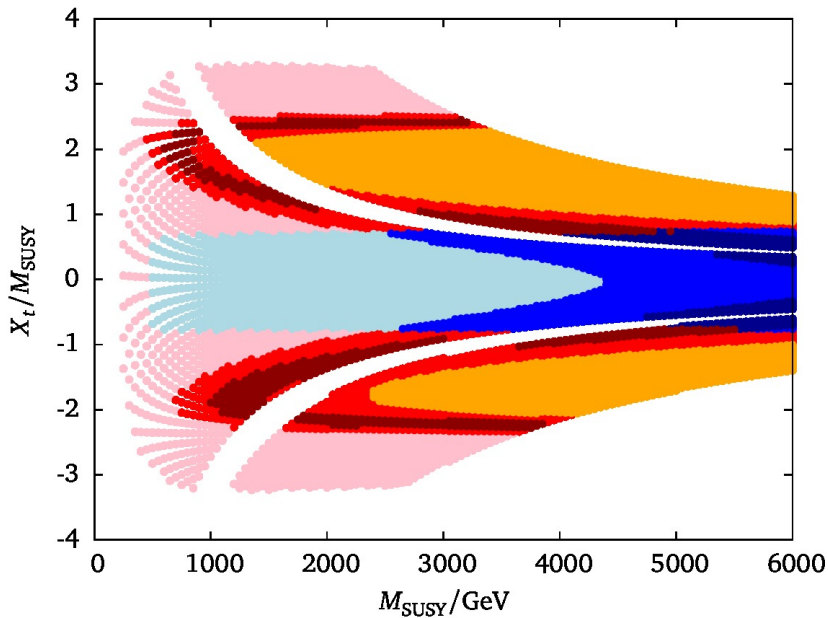
[WGH 2016]

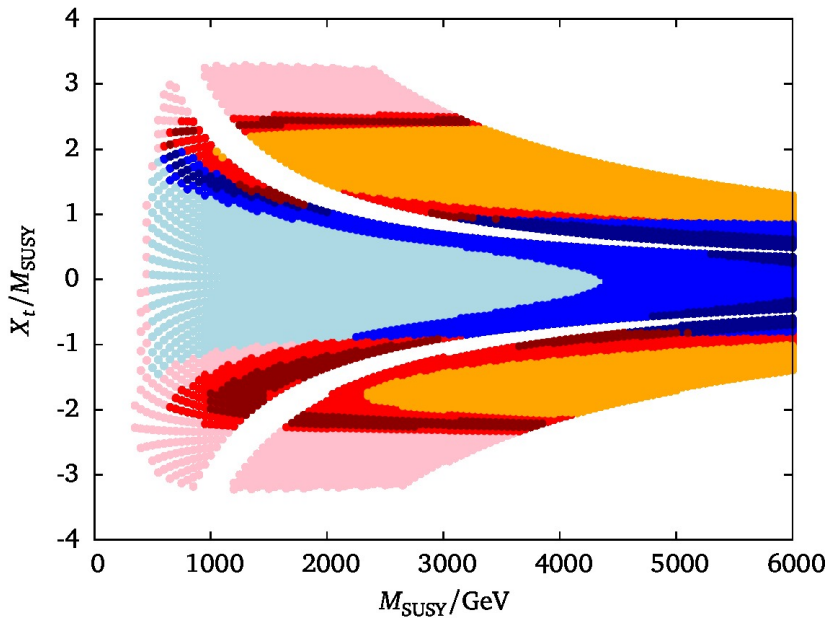


[WGH 2016]









A comment on metastability and quantum tunneling

Cosmological stability

bounce action

$$B \gtrsim 400$$

↔ life-time longer than age of the universe

Decay probability (per unit volume)

$$\frac{\Gamma}{V} = Ae^{-B/\hbar}$$

[Coleman '77]

Death and doom

- value of B crucially depends on field space path
- multifield spaces: reduction to single field space (!)
- *independent* of SUSY parameter choice

A general n scalar potential

$$V(\vec{\phi}) = \lambda_{abcd} \phi_a \phi_b \phi_c \phi_d + A_{abc} \phi_a \phi_b \phi_c + m_{ab}^2 \phi_a \phi_b + t_a \phi_a + c$$

- includes up to 3^n stationary points
- initial vacuum at $\vec{\phi} = \vec{\phi}_v$

$$\left. \frac{\partial V}{\partial \phi_a} \right|_{\vec{\phi} = \vec{\phi}_v} = 0$$

Expanding around the vacuum

- $\vec{\phi} = \vec{\phi}_v + \vec{\varphi}$, with $\vec{\varphi} = (\varphi_1, \dots, \varphi_n)^T$

$$V(\vec{\varphi}) = \lambda'_{abcd} \varphi_a \varphi_b \varphi_c \varphi_d + A'_{abc} \varphi_a \varphi_b \varphi_c + m_{ab}'^2 \varphi_a \varphi_b$$

- rewrite $\vec{\varphi} \rightarrow \varphi \hat{\varphi}$ with unit vector $\hat{\varphi}$, $\varphi = \sqrt{\varphi_1^2 + \dots + \varphi_n^2}$

$$V(\varphi, \hat{\varphi}) = \lambda(\hat{\varphi}) \varphi^4 - A(\hat{\varphi}) \varphi^3 + m^2(\hat{\varphi}) \varphi^2$$

A general n scalar potential

$$V(\vec{\phi}) = \lambda_{abcd} \phi_a \phi_b \phi_c \phi_d + A_{abc} \phi_a \phi_b \phi_c + m_{ab}^2 \phi_a \phi_b + t_a \phi_a + c$$

- includes up to 3^n stationary points
- initial vacuum at $\vec{\phi} = \vec{\phi}_v$

$$\left. \frac{\partial V}{\partial \phi_a} \right|_{\vec{\phi}=\vec{\phi}_v} = 0$$

A semi-analytic approximation

A quartic potential:

$$V(\phi) = \lambda \phi^4 - A^2 \phi^3 + m^2 \phi^2$$

$$B = \frac{\pi^2}{3\lambda} (2 - \delta)^{-3} (13.832 \delta - 10.819 \delta^2 + 2.0765 \delta^3)$$

with

$$\delta = \frac{8\lambda m^2}{A^2}$$

[Adams 1993]

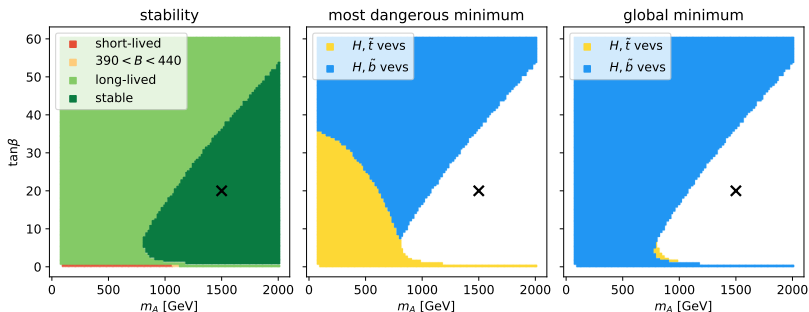
Benchmark scenario M_h^{125}

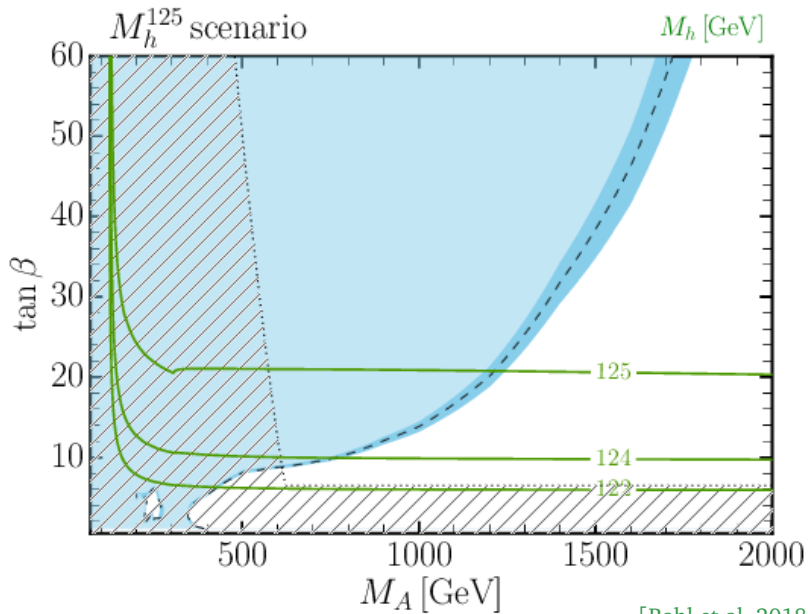
[Bahl et al. 2018]

$$m_{Q_3} = m_{U_3} = m_{D_3} = 1.5 \text{ TeV}, \quad m_{L_3} = m_{E_3} = 2 \text{ TeV}, \quad \mu = 1 \text{ TeV},$$

$$X_t = A_t - \frac{\mu}{\tan \beta} = 2.8 \text{ TeV}, \quad A_b = A_\tau = A_t,$$

$$M_1 = M_2 = 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV}$$





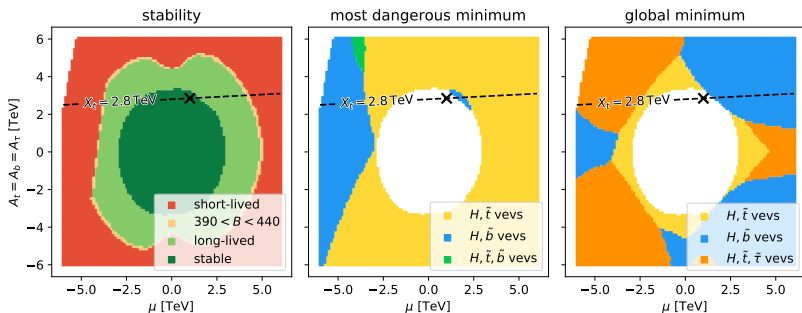
[Bahl et al. 2018]

An absolutely stable and experimentally allowed point

$$\tan\beta = 20$$

$$m_A = 1500 \text{ GeV}$$

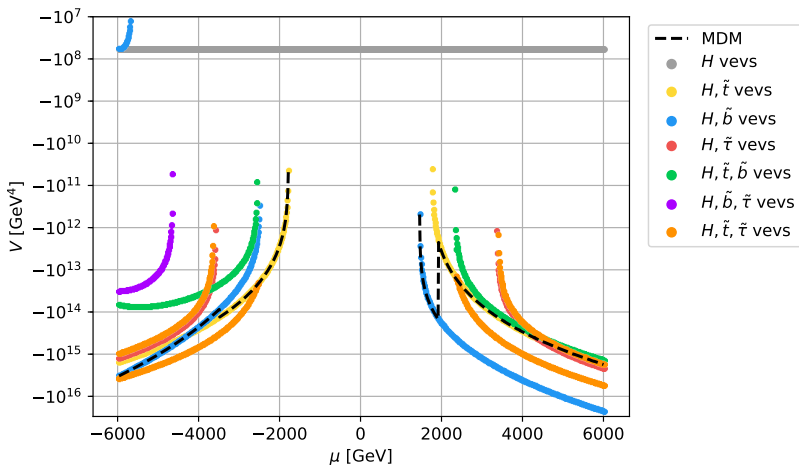
$$A \equiv A_t = A_b = A_\tau$$



- tachyonic sbottom masses upper left corner
- caveat: still limited numbers of fields included!

How is the “most dangerous minimum” (MDM) defined?

Go along the dashed line with $X_t = 2.8 \text{ TeV}$



A vast set of constraints

$$A_t^2 + 3\mu^2 < (m_{\tilde{t}_R}^2 + m_{\tilde{t}_L}^2) \cdot \begin{cases} 3 & \text{stable,} \\ 7.5 & \text{long-lived.} \end{cases}$$

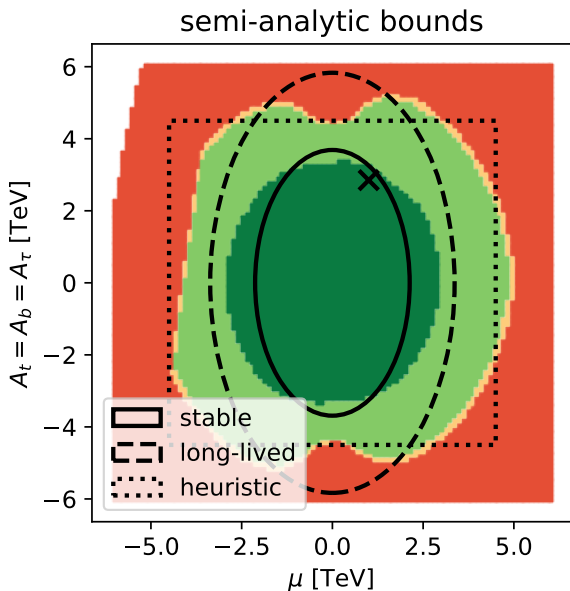
[Casas, Lleyda, Muñoz 1996, Kusenko, Langacker, Segre 1996]

Furthermore, a “heuristic” bound of

$$\frac{\max(A_{\tilde{t}, \tilde{b}}, \mu)}{\min(m_{Q_3, U_3})} \lesssim 3$$

exists. [Bechtle, Haber, Heinemeyer, Stefaniak, Stål, Weiglein, Zeune 2016]

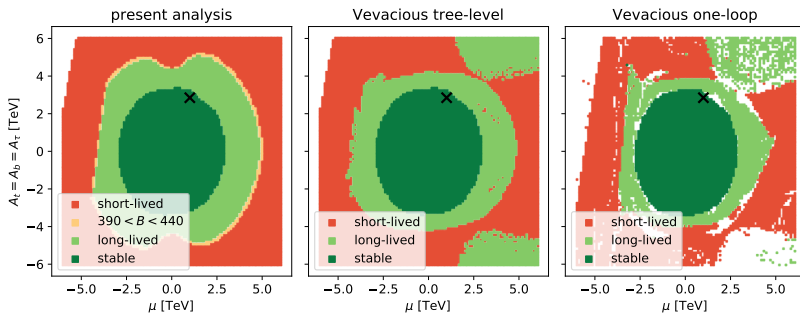
- vacuum tunneling weakens the “traditional” constraint
- metastable vacuum vs. absolute minimum
- quick and dirty versus sophisticated and precise (i. e. slow)
 \hookrightarrow needs numerical evaluation!



- inclusion of one-loop effective potential
- thermal corrections
- quantum tunneling by CosmoTransitions [Wainwright 2011]

Vevacious

[Camargo-Molina, O'Leary, Prood, Staub 2013]



[WGH, Weiglein, Wittbrodt 2019]

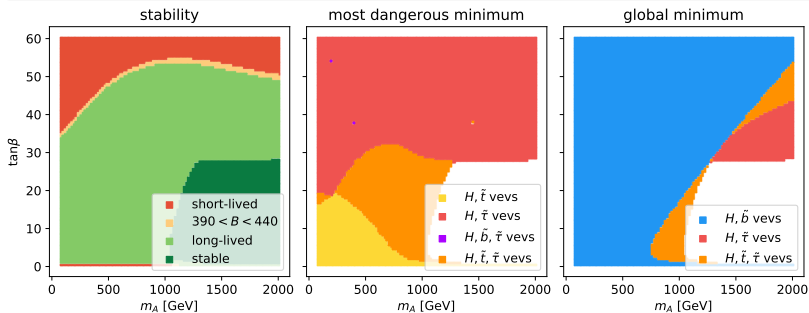
Benchmark scenario $M_h^{125}(\tilde{\tau})$: light stau

[Bahl et al. 2018]

$$m_{Q_3} = m_{U_3} = m_{D_3} = 1.5 \text{ TeV}, \quad m_{L_3} = m_{E_3} = 350 \text{ GeV}, \quad \mu = 1 \text{ TeV},$$

$$X_t = A_t - \frac{\mu}{\tan\beta} = 2.8 \text{ TeV}, \quad A_b = A_t, \quad A_\tau = 800 \text{ GeV},$$

$$M_1 = M_2 = 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV}$$

large $\tan\beta$ values: short-lived

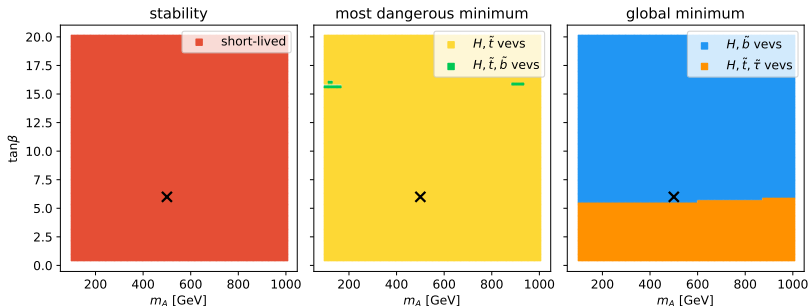
Benchmark scenario M_h^{125} (alignment)

[Bahl et al. 2018]

$$m_{Q_3} = m_{U_3} = m_{D_3} = 2.5 \text{ TeV}, \quad m_{L_3} = m_{E_3} = 2 \text{ TeV},$$

$$\mu = 7.5 \text{ TeV}, \quad A_t = A_b = A_\tau = 6.25 \text{ TeV},$$

$$M_1 = 500 \text{ GeV}, \quad M_2 = 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV}$$



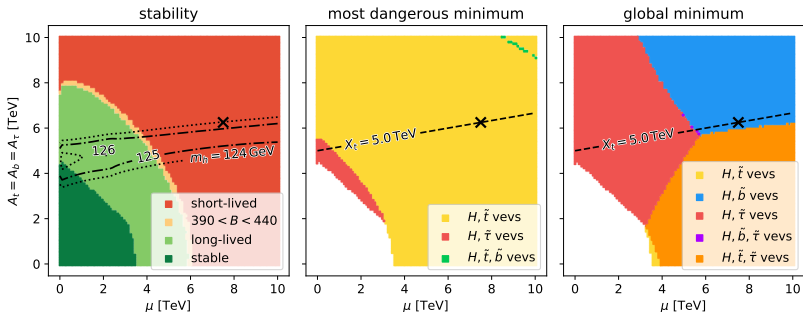
Benchmark scenario M_h^{125} (alignment)

[Bahl et al. 2018]

$$m_{Q_3} = m_{U_3} = m_{D_3} = 2.5 \text{ TeV}, \quad m_{L_3} = m_{E_3} = 2 \text{ TeV},$$

$$\mu = 7.5 \text{ TeV}, \quad A_t = A_b = A_\tau = 6.25 \text{ TeV},$$

$$M_1 = 500 \text{ GeV}, \quad M_2 = 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV}$$



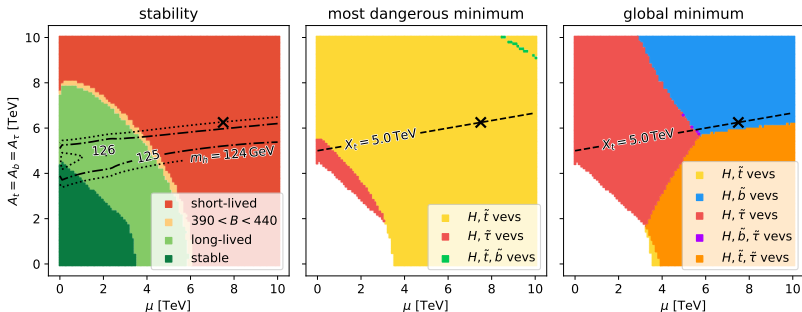
Benchmark scenario M_h^{125} (alignment)

[Bahl et al. 2018]

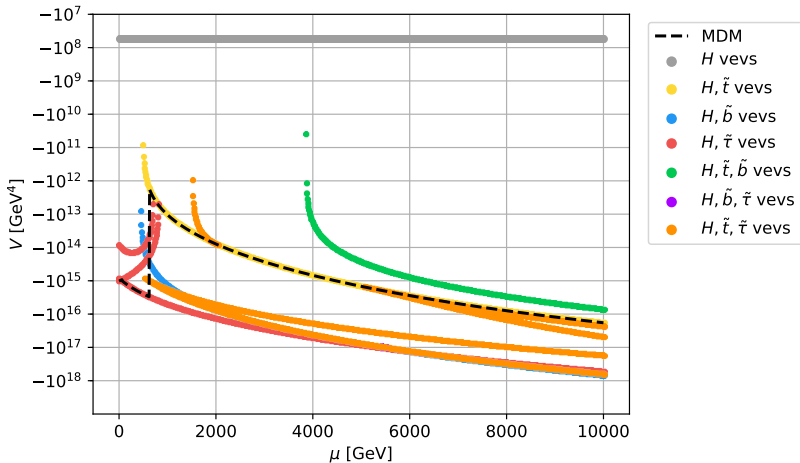
$$m_{Q_3} = m_{U_3} = m_{D_3} = 2.5 \text{ TeV}, \quad m_{L_3} = m_{E_3} = 2 \text{ TeV},$$

$$\mu = 7.5 \text{ TeV}, \quad A_t = A_b = A_\tau = 6.25 \text{ TeV},$$

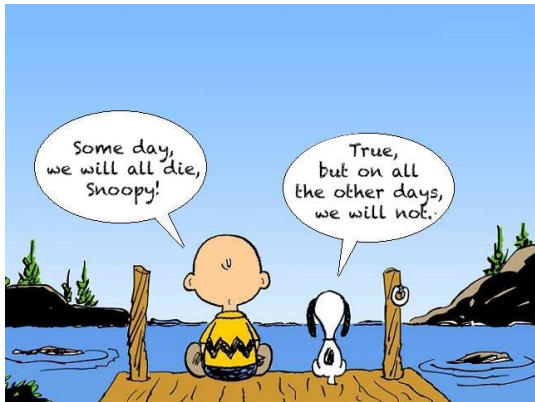
$$M_1 = 500 \text{ GeV}, \quad M_2 = 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV}$$



$$\mu = 7.5 \text{ TeV} \rightarrow 4 \text{ TeV} \quad \text{and} \quad A = 6.25 \text{ TeV} \rightarrow 5 \text{ TeV}$$



- severe constraints from vacuum (meta)stability in SUSY
- fast and numerically stable approach
- global minimum not the “most dangerous” one
- tree-level analysis sufficient (in comparison with 1-loop)



Backup

Slides

False vacuum decay

The bounce solution

[Coleman 1977]

$$\frac{d^2 \phi}{d\rho^2} + \frac{3}{\phi} \frac{d\phi}{d\rho} = \frac{\partial U}{\partial \phi}$$

with boundary conditions

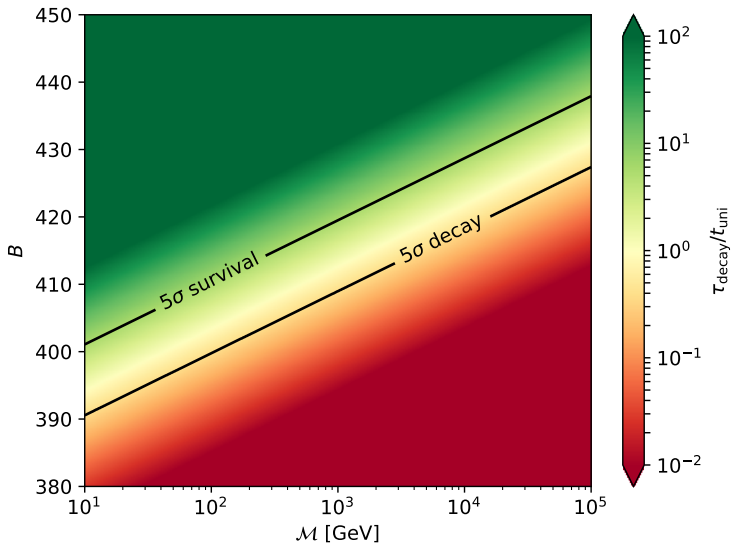
$$\phi(\infty) = \phi_v, \quad \left. \frac{d\phi}{d\rho} \right|_{\rho=0} = 0$$

U is the euclidean scalar potential, ρ is a spacetime variable and ϕ_v is the location of the metastable minimum.

The bounce action B is the stationary point of the euclidean action given by the integral

$$B = 2\pi^2 \int_0^\infty \rho^3 d\rho \left[\frac{1}{2} \left(\frac{d}{d\rho} \phi_B(\rho) \right)^2 + U(\phi_B(\rho)) \right]$$

The Bounce variability



Reduction to a single real scalar field

$$\phi \rightarrow \frac{1}{\sqrt{2}} \operatorname{Re}(\phi) + \frac{i}{\sqrt{2}} \operatorname{Im}(\phi)$$

- φ is canonically normalised after expanding $\vec{\varphi} = \varphi \hat{\varphi}$
- EW vacuum is given by

$$\operatorname{Re}(h_u^0) = v \sin \beta, \quad \operatorname{Re}(h_d^0) = v \cos \beta$$

where $v = \sqrt{v_u^2 + v_d^2} \approx 246 \text{ GeV}$ is the SM Higgs vev

Unfeasible to vary all real scalar degrees of freedom simultaneously \hookrightarrow selection of fields

$$\begin{aligned} & \{ \operatorname{Re}(h_u^0), \operatorname{Re}(h_d^0), \operatorname{Re}(\tilde{t}_L), \operatorname{Re}(\tilde{t}_R), \operatorname{Re}(\tilde{b}_L), \operatorname{Re}(\tilde{b}_R) \} \\ & \{ \operatorname{Re}(h_u^0), \operatorname{Re}(h_d^0), \operatorname{Re}(\tilde{t}_L), \operatorname{Re}(\tilde{t}_R), \operatorname{Re}(\tilde{\tau}_L), \operatorname{Re}(\tilde{\tau}_R) \} \\ & \{ \operatorname{Re}(h_u^0), \operatorname{Re}(h_d^0), \operatorname{Re}(\tilde{b}_L), \operatorname{Re}(\tilde{b}_R), \operatorname{Re}(\tilde{\tau}_L), \operatorname{Re}(\tilde{\tau}_R) \} \end{aligned}$$