

Project C2b

Exclusive non-leptonic and rare b -quark decays

Guido Bell, Thorsten Feldmann and Tobias Huber

CRC Annual Meeting, 6.-8. October 2020



Project C2b: People and Work Areas

<i>Name</i>	<i>Position</i>	<i>financed by</i>	<i>Period</i>	<i>contributes to</i>	
Marzia Bordone	postdoc	C2b	Oct 19 - Sep 20	WA 4	WA 5
Nico Gubernari	postdoc	C2b	Oct 20 –	WA 2	WA 5
Gilberto T.-X.	postdoc	C2b	Oct 19 –	WA 1	WA 4
Nicolas Seitz	PhD	C2a	Jul 19 –	WA 2	
Oscar Cata	Akad. Rat	U Siegen		WA 5	
Rusa Mandal	postdoc	A.v.H. fellow		WA 5	

WA 1: NNLO QCD corrections.

WA 2: QED corrections.

WA 3: Power corrections (and factorization).

WA 4: Phenomenology of nonleptonic decays.

WA 5: Phenomenology of rare semileptonic and radiative decays.

Publications and Preprints:

- 1** **G. Bell**, M. Beneke, **T. Huber** and X. Q. Li, “Two-loop non-leptonic penguin amplitude in QCD factorization,” JHEP **04** (2020), 055. QFET → WA1
- 2** A. Lenz and **G. Tetlalmatzi-Xolocotzi**, “Model-independent bounds on NP effects in non-leptonic tree-level decays of B-mesons,” JHEP **07** (2020), 177. WA4
- 3** **T. Huber**, J. Virto and K. Vos, “3-Body non-leptonic heavy-to-heavy B decays at NNLO in QCD,” arXiv:2007.08881 [hep-ph]. WA4
- 4** **M. Bordone**, N. Gubernari, **T. Huber**, M. Jung and D. van Dyk, “A puzzle in $\bar{B}_{(s)}^0 \rightarrow D_{(s)}^+ \{\pi^-, K^-\}$ decays and extraction of the f_s/f_d fragmentation fraction,” arXiv:2007.10338 [hep-ph]. WA4
- 5** **O. Catà** and **T. Mannel**, “Linking lepton number violation with B anomalies,” arXiv:1903.01799 [hep-ph] QFET → WA5
- 6** **M. Bordone**, **O. Catà** and **T. Feldmann**, “Effective theory approach to NP with flavour: general framework and a leptoquark example,” JHEP **2001** (2020) 067. QFET → WA5
- 7** **M. Bordone**, **O. Catà**, **T. Feldmann** and **Rusa Mandal**, “Constraining flavour patterns of scalar leptoquarks in the EFT ,” arXiv:2010.03297 [hep-ph] WA5

1 Results from WA5 (phenomenology of rare decays)

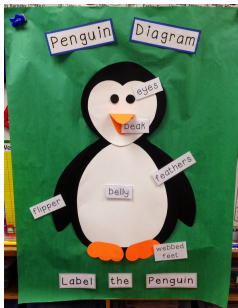
- Bordone/Cata/Feldmann P3H-19-032
- Bordone/Cata/Feldmann/Mandal P3H-20-046

2 Results from WA4 (phenomenology of nonleptonic decays)

- Huber/Virto/Vos P3H-20-037

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- some technical aspects relevant for WA1 – WA3
have also been discussed at SCET workshop of the CRC, January 2020

WA 5



LQ symbol	$[SU(3), SU(2), U(1)]$	Spin	F
S_3	$(\bar{3}, 3, 1/3)$	0	-2
\tilde{S}_1	$(\bar{3}, 1, 4/3)$	0	-2
S_1	$(\bar{3}, 1, 1/3)$	0	-2
\tilde{S}_1	$(\bar{3}, 1, -2/3)$	0	-2
R_2	$(3, 2, 7/6)$	0	0
\tilde{R}_2	$(3, 2, 1/6)$	0	0
U_3	$(3, 3, 2/3)$	1	0
\tilde{U}_1	$(3, 1, 5/3)$	1	0
U_1	$(3, 1, 2/3)$	1	0
\tilde{U}_1	$(3, 1, -1/3)$	1	0
V_2	$(\bar{3}, 2, 5/6)$	1	-2
\tilde{V}_2	$(\bar{3}, 2, -1/6)$	1	-2

[leptoquark table from arXiv:1903.04977]

“Effective Theory Approach to New Physics with Flavour: General Framework and a Leptoquark Example”

[Bordone, Cata, Feldmann 2019]

SMEFT is not built as an EFT of flavour, . . .

- NP operators in SM-EFT introduce various flavour coefficients.
→ too many free parameters for generic pheno analyses
- $\mathcal{O}(1)$ couplings would require NP scales above LHC reach.
→ global-fit scan would have to include “non-natural” regions of parameter-space
- minimal flavour violation cannot explain current “ B -anomalies”
→ needs extension/generalization of MFV

“Effective Theory Approach to New Physics with Flavour: General Framework and a Leptoquark Example”

[Bordone, Cata, Feldmann 2019]

SMEFT is not built as an EFT of flavour, . . . but it can account for flavour patterns:

- extending MFV with new “flavour spurions”
→ NP explanations for B -anomalies favour leptoquark scenarios
- providing a power-counting scheme
→ simplest solution: flavour-dependent Froggatt-Nielsen charges,
for quark and lepton couplings

[Feldmann, Mannel 2006]

⇒ Rather general and systematic framework,
that contains and extends MFV.

⇒ Spurion scenarios and FN charges to be fixed by Phenomenology.

- consider, as an example, a four-fermion operator in SMEFT: (→ LFU violation)

$$\frac{1}{\Lambda^2} [C_{\ell q}]^{ij\alpha\beta} (\bar{Q}_i \gamma_\mu Q_j) (\bar{L}_\alpha \gamma^\mu L_\beta)$$

- flavour tensor $[C_{\ell q}]^{ij\alpha\beta}$ introduces 3^4 free real parameters:

generic EFT: $\mathcal{O}(1)$

MFV:

FN charges:

LQ+FN charges:

- SM Yukawa matrices Y_U, Y_D
- small expansion parameter $\lambda \sim \sin \theta_C \sim 0.2$
- family-dependent charges b'_Q etc. of a hypothetical $U(1)$ symmetry
- new “leptoquark spurion” $\Delta_{QL}^{i\alpha}$

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MFV: $(\#\delta^{ij} + \#(Y_U Y_U^\dagger)^{ij} + \#(Y_D Y_D^\dagger)^{ij} + \dots) (\delta^{\alpha\beta} + \dots)$

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LQ+FN charges:	$\sim (\Delta_{QL})^{i\beta} (\Delta_{QL}^\dagger)^{\alpha j} + \dots \sim \lambda b_Q^j - b_L^\beta \lambda b_L^\alpha - b_Q^j $

- SM Yukawa matrices Y_U, Y_D
- small expansion parameter $\lambda \sim \sin \theta_C \sim 0.2$
- family-dependent charges b_Q^j etc. of a hypothetical $U(1)$ symmetry
- new “leptoquark spurion” $\Delta_{QL}^{i\alpha}$

Features:

- independent of particular *dynamical* realization in the UV ✓
- exponents in the power-counting fulfill **triangle inequalities**, like e.g.

$$\left| b_Q^j - b_L^\beta \right| + \left| b_L^\alpha - b_Q^j \right| \geq \left| b_Q^j - b_Q^j + b_L^\alpha - b_L^\beta \right|$$

→ power-counting not violated by higher-orders in spurion expansion ✓

- allows to play with different scenarios in SMEFT, depending on:
 - which fundamental spurions to consider (Δ_{QL}, \dots),
 - which values for the FN charges to take (see below),
 - + possibly additional simplifying assumptions (see below).

→ **test different scenarios** against **NP signals in flavour observables !**

- | | |
|---|---|
| • $b \rightarrow sll$ processes (R_K, R_K^*, P_5') | • $B \rightarrow K^{(*)} \nu \bar{\nu}$ |
| • $b \rightarrow cl\nu$ processes (R_D, R_{D^*}) | • $K^+ \rightarrow \pi^+ \nu \bar{\nu}$, |
| • B -meson mixing ($\Delta M_d, \Delta M_s$) | • LFV kaon decays ($K_L \rightarrow \mu e$) |
| • W universality | • LFV B -meson decays |
| • Z decays ($Z \rightarrow \nu \bar{\nu}, Z \rightarrow \ell^+ \ell^-$) | • charged LFV ($\mu \rightarrow e \gamma, \tau \rightarrow \mu \gamma$) |

Determining the FN charges:

- operators with (Dirac-) fermions only depend on **charge differences**
- left-handed quark charges from CKM angles (canonical) ✓

$$b_Q^1 \equiv 3, \quad b_Q^2 \equiv 2, \quad b_Q^3 \equiv 0$$

- L-R differences from Yukawa couplings ✓

$$\begin{aligned} |b_Q^1 - b_U^1| &\simeq 8, & |b_Q^2 - b_U^2| &\simeq 4, & b_U^3 &= 0, \\ |b_Q^1 - b_D^1| &\simeq 7, & |b_Q^2 - b_D^2| &\simeq 5, & |b_Q^3 - b_D^3| &\simeq 3 \end{aligned}$$

and

$$|b_L^1 - b_E^1| \simeq 9, \quad |b_L^2 - b_E^2| \simeq 5, \quad |b_L^3 - b_E^3| \simeq 3$$

⇒ 3 unconstrained charges plus a 2^8 -fold discrete ambiguity (!)

- Further information from particular LQ scenario plus additional assumptions

Scenario with U_1 vector leptoquark: (hypercharge $y = 2/3$)

- involves two additional spurions:

$$\text{left-handed: } \Delta_{QL}^{i\alpha} \sim \lambda |b_Q^i - b_L^{\alpha}|, \quad \text{right-handed: } \Delta_{DE}^{i\alpha} \sim \lambda |b_D^i - b_E^{\alpha}|$$

- (approx.) tree-level relations for 4-fermion operators
relates $SU(2)$ -singlet and -triplet couplings

→ suppress e.g. rare $K \rightarrow \pi \nu \bar{\nu}$ decays

- use various precision and/or rare flavour observables

$$Z \rightarrow \nu \bar{\nu}, b \rightarrow s \mu^+ \mu^-, b \rightarrow c \tau^- \nu, \\ \bar{B}_d \rightarrow \tau^- \mu^+, \bar{B}_s \rightarrow \tau^\pm \mu^\mp, K_L \rightarrow \mu^\pm e^\mp$$

to constrain **order of magnitude** of spurion entries

⇒ 24 possible combinations of FN charges

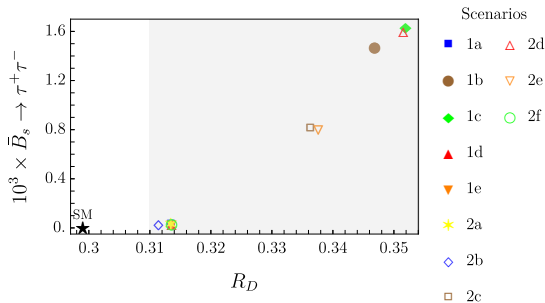
- fit of overall coefficients to experimental data
(assuming real numbers in down-quark basis, and no fine-tuning/texture zeros)

⇒ 11 phenomenologically viable scenarios for FN charges

Conclusions from U_1 vector leptoquark scenario:

- Simultaneous explanation of W universality tests and $R_D^{(*)}$ provides the strongest tensions in the fit.
- Interesting predictions for $B \rightarrow \tau^+ \tau^-$ and $B \rightarrow \tau^\pm \mu^\mp$.

Example: Correlations between $B_s \rightarrow \tau^+ \tau^-$ and R_D for different scenarios



“Constraining flavour patterns of scalar leptoquarks in the EFT”

[Bordone, Cata, Feldmann, Mandal 2020 (today)]

Motivation:

- models with scalar leptoquarks are renormalizable (independent of UV completion)
→ consistent computation of **loop-induced decays**
- sensitive to decays like $B \rightarrow K^{(*)} \nu \nu$ and $K \rightarrow \pi \nu \nu$
(which are naturally suppressed in U_1 scenario)
- EFT with FN charges → **more sophisticated analyses** compared to existing fits
(without ad-hoc flavour structures and including all relevant observables)

"Which scalar leptoquark-inspired spurions can explain the B anomalies
without upsetting existing low-energy constraints?"

Scenario with two scalar leptoquarks S_3 and S_1 (hypercharge $y = 1/3$)

spurions: $S_{QL}^{i\alpha} = c_L^{i\alpha} \lambda^{|b_Q^i - b_L^\alpha|}$, $\tilde{S}_{QL}^{i\alpha} = \tilde{c}_L^{i\alpha} \lambda^{|b_Q^i - b_L^\alpha|}$, $\tilde{S}_{UE}^{i\alpha} = c_R^{i\alpha} \lambda^{|b_U^i - b_E^\alpha|}$

- essentially two different solutions for viable FN charge assignments

Nominal scenario (A):

$$S_{QL}^{(A)} \sim \begin{pmatrix} \lambda^8 & \lambda^4 & \lambda^2 \\ \lambda^7 & \lambda^3 & \lambda \\ \lambda^5 & \lambda & -\lambda \end{pmatrix}, \quad \tilde{S}_{QL}^{(A)} \sim \begin{pmatrix} \lambda^8 & \lambda^4 & -\lambda^2 \\ \lambda^7 & \lambda^3 & \lambda \\ \lambda^5 & -\lambda & \lambda \end{pmatrix}, \quad \tilde{S}_{UE}^{(A)} \sim \begin{pmatrix} \lambda^{25} & \lambda^{17} & \lambda^{13} \\ \lambda^{12} & \lambda^4 & -\lambda^0 \\ \lambda^{14} & \lambda^6 & \lambda^2 \end{pmatrix}$$

corresponding to FN charges:

$$b_Q = (3, 2, 0), \quad b_U = (11, -2, 0), \quad b_L = (-5, -1, 1), \quad b_E = (-14, -6, -2)$$

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corresponding to FN charges:

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Some phenomenologically interesting predictions:

Mode	\mathcal{B}_{SM}	\mathcal{B}_{Exp}	scalar LQs with FN
$\tau \rightarrow \mu\phi$	0	$< 8.4 \times 10^{-8}$	$[0.58, 1.25] \times 10^{-10}$
$B_s \rightarrow \tau^\pm \mu^\mp$	0	$< 4.2 \times 10^{-5}$	$[1.21, 2.60] \times 10^{-6}$
$B^+ \rightarrow K^+ \tau^+ \tau^-$	$(1.60 \pm 0.12) \times 10^{-7}$	$< 2.2 \times 10^{-3}$	$[7.9, 13.3] \times \mathcal{B}_{\text{SM}}$
$\bar{B}_s \rightarrow \tau^+ \tau^-$	$(7.30 \pm 0.49) \times 10^{-7}$	$< 6.8 \times 10^{-3}$	$[7.75, 13.1] \times \mathcal{B}_{\text{SM}}$
$B^+ \rightarrow K^+ \tau^+ \mu^-$	0	$< 3.9 \times 10^{-5}$	$(1.8 \pm 0.7) \times 10^{-6}$

$\tau \rightarrow \mu\gamma$: dominated by top-quark loops,

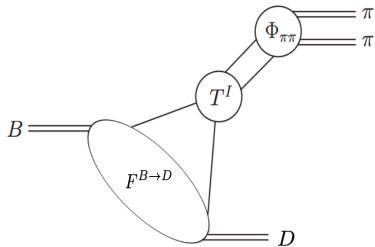
$$\mathcal{B}(\tau \rightarrow \mu\gamma) \sim g_1^2 g_R^2 (\tilde{c}_L^{32})^2 (c_R^{33})^2 \times 8.67 \cdot 10^{-7}.$$

Same combination of spurion entries also appears in $R_{D^{(*)}}$

\Rightarrow Tension between $b \rightarrow c \tau^- \bar{\nu}$ and $\tau \rightarrow \mu\gamma$

$(g-2)_\mu$: can only be explained by bending the power-counting for c_R^{32} and/or accepting a tension with $\tau \rightarrow \mu\gamma$.

WA 4



“3-body non-leptonic heavy-to-heavy B decays at NNLO in QCD”

[Huber/Virto/Vos 2020]

QCD Factorization for two-body $\bar{B}_d \rightarrow D^+ \pi^-$ decays well established

$$\langle D^+ \pi^- | \mathcal{Q}_i | \bar{B}_d \rangle = \sum_j F_j^{B \rightarrow D}(m_L^2) \int_0^1 du T_{ij}(u) \phi_\pi(u) + \mathcal{O}(\Lambda_{\text{QCD}}/m_b)$$

- Only vertex terms of colour-allowed tree-amplitude contribute (no colour-suppressed tree-amplitude, no penguin amplitudes)
- Hard spectator scattering and weak annihilation power-suppressed
- hadronic information in terms of form factors $F_j^{B \rightarrow D}$ at $q^2 = m_\pi^2$ and light-cone distribution amplitude ϕ_π for pion

Idea:

Establish factorization for $\bar{B}^0 \rightarrow D^+ K^- \pi^0$ and $\bar{B}^0 \rightarrow D^+ \pi^- \pi^0$ decays
in phase-space region of small invariant mass, $k^2 = M_{K\pi, \pi\pi}^2 \ll m_B^2$

- Hard kernels $T_{ij}(u)$ identical to two-body decay and known to NNLO [Huber, Kränkl, Li'16]
- new hadronic information in di-meson distribution amplitudes, e.g.

$$\Phi_{\pi\pi}(u) = 6u(1-u) \sum_{n=0}^{\infty} \alpha_n^{\pi\pi} C_n^{3/2}(2u-1)$$

where Gegenbauer coefficients $\alpha_n^{\pi\pi}$ now depend on $(k^2, \cos \theta)$

- first NNLO study of non-leptonic 3-body decays
- numerical impact of higher-order corrections
- estimate of finite-width effects in quasi-2-body kinematics

Gegenbauer coefficients can be expanded in **partial waves**:

- in case of $\bar{B}^0 \rightarrow D^+ \pi^- \pi^0$ case:

$$\alpha_n^{\pi\pi}(k^2, \theta_\pi) = \sum_{\ell=1,3,\dots}^{n+1} B_{n\ell}^{\pi\pi}(k^2) P_\ell(\cos \theta_\pi) \quad (n \text{ even})$$

with normalization for $n = 0$ fixed by (time-like) form factors, (from data)

$$B_{01}^{\pi\pi}(k^2) \propto f_+^{\pi\pi}(k^2)$$

- in case of $\bar{B}^0 \rightarrow D^+ K^- \pi^0$:

$$\alpha_n^{K\pi}(k^2, \theta_\pi) = \sum_{\ell=0}^{n+1} B_{n\ell}^{K\pi}(k^2) P_\ell(\cos \theta_\pi) \quad (\text{all } n)$$

with normalization for $n = 0$ fixed by two form factors (from data)

$$B_{00}^{K\pi}(k^2) \propto f_0^{K\pi}(k^2), \quad B_{01}^{K\pi}(k^2) \propto f_+^{K\pi}(k^2)$$

The set of functions $B_{n\ell}^L(k^2)$ determines the k^2 spectrum of each partial wave.

Perturbative expansion of the decay amplitude:

$$\mathcal{A} \propto (V_{ux}^* V_{cb}) F_0^{B \rightarrow D} \sum_{n \geq \max(\ell-1, 0)} B_{nl}^{\pi/K} \mathcal{G}_n$$

- process-dependent information in short-distance functions

$$\mathcal{G}_n \equiv C_1(\mu) \mathcal{V}_{1n}(\mu) + C_2(\mu) \mathcal{V}_{2n}(\mu)$$

- The functions \mathcal{V}_{in} are known to NNLO in QCD

[TH/Kränkl/Li 2016]

$$\mathcal{G}_0(\mu_b) = 1.034_{\text{LO}} + (0.026 + i0.020)_{\text{NLO}} + (0.013 + i0.027)_{\text{NNLO}},$$

$$\mathcal{G}_1(\mu_b) = (-0.013 + i0.030)_{\text{NLO}} + (-0.044 + i0.018)_{\text{NNLO}},$$

$$\mathcal{G}_2(\mu_b) = (0.0023 - i0.0017)_{\text{NLO}} + (0.0017 - i0.0054)_{\text{NNLO}}.$$

- NNLO contributions large relative to NLO (\leftarrow colour factors!)
- NNLO essential for contributions of higher terms ($n \geq 1$) in Gegenbauer expansion

Application: Probing higher-order QCD effects,
higher Gegenbauer moments, and higher partial waves:

- Define ratios of angular-integrated decay widths

$$\mathcal{R}_{MM}[z_1, z_2; z'_1, z'_2](k^2) \equiv \frac{\int_{z_1}^{z_2} dz \frac{d\Gamma(\bar{B} \rightarrow D^+ M^- \pi^0)}{dk^2 dz}}{\int_{z'_1}^{z'_2} dz \frac{d\Gamma(\bar{B} \rightarrow D^+ M^- \pi^0)}{dk^2 dz}}$$

- dependence on form factors, CKM elements cancel.

These ratios are being analyzed at the Belle-2 experiment.

Modelling the functions $B_{n\ell}^L(k^2)$ from sum over individual resonances R :

- Yields for S - and P -wave coefficients

$$B_{00}^{K\pi}(s) = \sum_{R_0} \frac{m_{R_0} f_{R_0} g_{R_0 K\pi} e^{i\varphi_{R_0}}}{\sqrt{2}[m_{R_0}^2 - s - i\sqrt{s}\Gamma_{R_0}(s)]}$$

$$B_{01}^{\pi\pi}(s) = \sum_R \frac{m_R f_R g_{R\pi\pi} e^{i\varphi_R}}{\sqrt{2}[m_R^2 - s - i\sqrt{s}\Gamma_R(s)]}$$

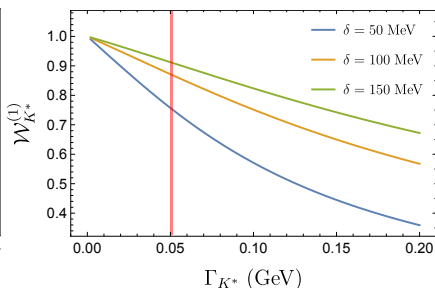
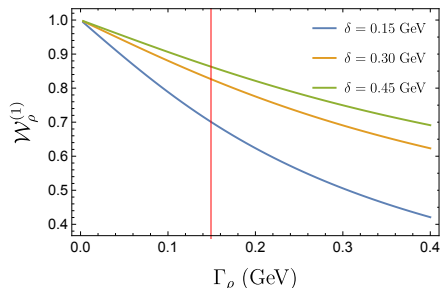
$$B_{01}^{K\pi}(s) = \frac{\sqrt{\lambda_{K\pi}(s)}}{s} \sum_R \frac{m_R f_R g_{RK\pi} e^{i\varphi_R}}{\sqrt{2}[m_R^2 - s - i\sqrt{s}\Gamma_R(s)]}$$

- Satisfies narrow-width limit in case of stable vector resonance (ρ or K^*)

Application: leading corrections to narrow-width approximation:

- Define resonance width as a function of bin size 2δ :

$$\Gamma_{[R]}(\delta) \equiv \int_{(m_R-\delta)^2}^{(m_R+\delta)^2} ds \frac{d\Gamma(\bar{B} \rightarrow D^+ M^- \pi^0)}{ds} = \sum_{\ell} \Gamma_{[R]}^{(\ell)}, \quad \mathcal{W}_R^{(\ell)} \equiv \Gamma_{[R]}^{(\ell)} / \Gamma_{[R],\text{NWL}}^{(\ell)}$$



Typically, finite-width corrections are of the order 20%.



Ongoing projects:

- NNLO corrections to penguin coefficient a_6 (**TH, GTX**)
- mixed QCD/QED corrections in B -decays (**TF, TH, Gubernari, Seitz**)
- factorization of endpoint divergencies (**GB, Böer, TF**)
- QCD factorization and flavour symmetries (**TH, GTX**)
- rare decays of Λ_b baryons (**TF, Gubernari**)

→ WA1

→ WA2

→ WA3

→ WA4

→ WA5